

Cofactor Definition:

$$C_{ij} = (-1)^{i+j} \det$$

$$\begin{array}{|c|c|} \hline a_{11} \dots & a_{1m} \\ \hline \vdots & \vdots \\ \hline a_{i1} \dots & a_{ij} \dots a_{im} \\ \hline \vdots & \vdots \\ \hline a_{m1} \dots & a_{mm} \\ \hline \end{array} M_{ij}$$

$$A_{ik} C_{ij} = \delta_{kj} |A|$$

$$A_{ki} C_{ji} = \delta_{kj} |A|$$

Inverse of a $n \times n$ matrix:

Consult Exercise 3.2.24 on p. 113.

Assume determinant $|A| \neq 0$. Then,

$$(A^{-1})_{ij} = \frac{c_{ji}}{|A|}$$

where c_{ji} is the j th cofactor of $|A|$.

Proof:

$$A_{ki} c_{ji} = \delta_{kj} |A|$$

$$A_{ki} (A^{-1})_{ij} \Rightarrow A A^{-1} = 1$$

$$(A^{-1})_{ij} A_{jk} = \frac{c_{ji} A_{jk}}{|A|} = \delta_{ik}$$

$$\Rightarrow A^{-1} A = 1$$

Note: $|A^{-1} A| = |A^{-1}| |A| \quad (3, 49)$

$$|1| = 1 \Rightarrow |A^{-1}| = \frac{1}{|A|}$$

Orthogonal Matrices: (p. 193)

Application of matrices to vectors:

$$\vec{x}' = A \vec{x}$$

$$\vec{x}' = \begin{pmatrix} x'_1 \\ \vdots \\ x'_n \end{pmatrix}, \quad \vec{x} = \begin{pmatrix} x_1 \\ \vdots \\ x_n \end{pmatrix}$$

$$x'_i = a_{ij} x_j$$

Einstein
convention.

We are interested in "rotations", which leave the Euclidean norm $|\vec{x}| = \sqrt{\vec{x}^2}$ invariant, i.e., $\vec{x}'^2 = \vec{x}^2$.

$$\begin{aligned} x'_i x'_i &= a_{ij} x_j a_{ik} x_k = x_j x_k a_{ij} a_{ik} \\ &= x_j x_k \delta_{jk} \quad \text{for all } x_j, x_k \end{aligned}$$

$$\Rightarrow \underbrace{a_{ij} a_{ik}} = \delta_{jk} \quad \left. \begin{array}{l} \text{Mutually} \\ \text{orthonormal} \\ \text{vectors.} \end{array} \right\}$$

M (3)

Orthogonality condition on matrix A.

Transpose Matrix definition:

$$\tilde{A} = (\tilde{a}_{ij}) \text{ with } \tilde{a}_{ij} = a_{ji}.$$

Orthogonality relation: $\tilde{A} A = 1,$

$$\tilde{A} = A^{-1}.$$

Therefore also) $A \tilde{A} = 1$

$$a_{ij} \tilde{a}_{jk} = a_{ij} a_{kj} = \delta_{ik}$$

Determinants:

$$1 = |A \tilde{A}| = |A| |\tilde{A}|$$

$$= |A|^2$$

$$\Rightarrow \underline{|A| = \pm 1.}$$

$\frac{1}{2} n(n+1)$ independent equations.

$\frac{1}{2} n(n-1)$ free parameters left.

M (4)

2D: $\frac{1}{2} n (n-1) = 1$ free parameter

$$\sum_{i=1}^2 a_{ij} a_{ik} = a_{1j} a_{1k} + a_{2j} a_{2k} = \delta_{jk}$$

$$a_{11} a_{11} + a_{21} a_{21} = 1$$

$$a_{11} a_{12} + a_{21} a_{22} = 0$$

$$a_{12} a_{11} + a_{22} a_{21} = 0$$

$$a_{12} a_{12} + a_{22} a_{22} = 1$$

Solutions: $a_{11} = \cos \phi$, $a_{21} = -\sin \phi$
(+)

$$\cos \phi a_{12} - \sin \phi a_{22} = 0$$

$$a_{12} \cos \phi - a_{22} \sin \phi = 0$$

$$a_{12} = + \sin \phi, \quad a_{22} = + \cos \phi$$

(-)

$$A = \begin{pmatrix} \cos \phi & \sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix}$$

↑
Angle

Complex possibility:

$$a_{11} = \cosh \xi, \quad a_{21} = i \sinh \xi$$

$$\cosh \xi \quad a_{12} - i \sinh \xi \quad a_{22} = 0$$

$$a_{12} = i \sinh \xi, \quad a_{22} = \cosh \xi$$

$$A = \begin{pmatrix} \cosh \xi & -i \sinh \xi \\ i \sinh \xi & \cosh \xi \end{pmatrix}$$

↑
Rapidity

$S_1 \Rightarrow S$
↓

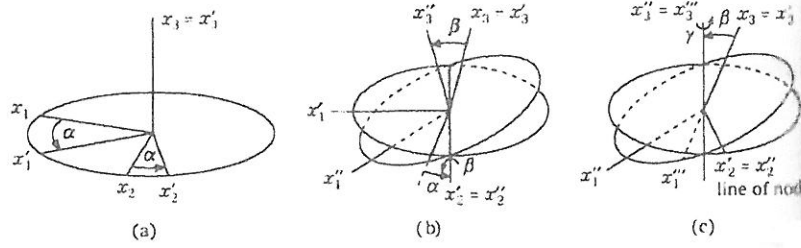
Special
Relativity
Supplement

3D: Show and explain Euler Angl

book p. 200/201.

Figure 3.7

- (a) Rotation About x_3
Through Angle α ,
(b) Rotation About x'_2
Through Angle β , and
(c) Rotation About x''_3
Through Angle γ



3.3 Orthogonal Matrices

201

The three matrices describing these rotations are

$$R_z(\alpha) = \begin{pmatrix} \cos \alpha & \sin \alpha & 0 \\ -\sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{pmatrix}, \quad (3.98)$$

exactly like Eq. (3.92), but different sign for β in

$$R_y(\beta) = \begin{pmatrix} \cos \beta & 0 & -\sin \beta \\ 0 & 1 & 0 \\ \sin \beta & 0 & \cos \beta \end{pmatrix} \quad (3.99)$$

and

$$R_z(\gamma) = \begin{pmatrix} \cos \gamma & \sin \gamma & 0 \\ -\sin \gamma & \cos \gamma & 0 \\ 0 & 0 & 1 \end{pmatrix}. \quad (3.100)$$

The total rotation is described by the triple matrix product:

$$A(\alpha, \beta, \gamma) = R_z(\gamma)R_y(\beta)R_z(\alpha). \quad (3.101)$$

Note the order: $R_z(\alpha)$ operates first, then $R_y(\beta)$, and finally $R_z(\gamma)$. Direct multiplication gives

$$A(\alpha, \beta, \gamma) = \begin{pmatrix} \cos \gamma \cos \beta \cos \alpha - \sin \gamma \sin \alpha & & & \\ -\sin \gamma \cos \beta \cos \alpha - \cos \gamma \sin \alpha & & & \\ \sin \beta \cos \alpha & & & \\ \cos \gamma \cos \beta \sin \alpha + \sin \gamma \cos \alpha & -\cos \gamma \sin \beta & & \\ -\sin \gamma \cos \beta \sin \alpha + \cos \gamma \cos \alpha & \sin \gamma \sin \beta & & \\ \sin \beta \sin \alpha & \cos \beta & & \end{pmatrix}. \quad (3.102)$$

Hermitian and Unitary Matrices

QM: Complex Variables.

$$z = x + iy, \quad i = \sqrt{-1}$$

$$z^* = x - iy \quad \text{complex conjugate.}$$

Matrix: $A^* = (a_{ij}^*)$

Adjoint matrix: $A^\dagger = (a_{ik})^\dagger = (\tilde{a}_{ik}^*) = (a_{ki}^*)$.

Hermitian matrix: $A^\dagger = A$

also called self-adjoint.

If A is real: real symmetric matrix.

Unitary matrix: $U^\dagger = U^{-1}$

\uparrow
U real: $U^\dagger = \tilde{U} = U^{-1}$

orthogonal matrices back.

Leave length of complex vector unchanged, Conservation of Probability in QM.

$$\vec{x} = \begin{pmatrix} x_1 \\ \vdots \\ x_m \end{pmatrix}, \quad \vec{x}^t = (x_1^*, \dots, x_m^*)$$

$$\vec{x}^t \cdot \vec{x} = \vec{x}^t \vec{x} = \sum_i x_i^* x_i = \sum_i |x_i|^2 \geq 0$$

Let $\vec{y} = U \vec{x}$:

$$\begin{aligned} \vec{y}^t \vec{y} &= (U \vec{x})^t U \vec{x} = \vec{x}^t U^t U \vec{x} \\ &= \vec{x}^t U^{-1} U \vec{x} = \vec{x}^t \vec{x} \end{aligned}$$

Transformation of an operator A :

$$\text{let } \vec{y} = A \vec{x}, \quad \vec{x}' = U \vec{x}, \quad \vec{y}' = U \vec{y}.$$

We want:
$$\vec{y}' = A' \vec{x}' \quad (*)$$

Solution:
$$A' = U A U^t$$

$$A' \vec{x}' = U A U^t U \vec{x} = U A \vec{x} = U \vec{y} = \vec{y}'$$

(*) is called Similarity Transformation.

Diagonalization of Matrices by Similarity Transformation.

Example: Moments of Inertia Matrix

$$\vec{L} = \mathbf{I} \vec{\omega}$$

with
$$I_{ij} = \sum_k m_k (\delta_{ij} r_k^2 - x_i x_j)$$

By rotation of axis we can bring

\mathbf{I} into diagonal form:

Principal Axis. In these coordinates

$$\mathbf{I}' = \begin{pmatrix} I'_1 & 0 & 0 \\ 0 & I'_2 & 0 \\ 0 & 0 & I'_3 \end{pmatrix} = \begin{pmatrix} \lambda_1 & 0 & 0 \\ 0 & \lambda_2 & 0 \\ 0 & 0 & \lambda_3 \end{pmatrix}$$

orthonormal
with eigenvectors

$$\hat{v}_i' \hat{v}_j' = \delta_{ij}$$

$$\hat{v}_1' = \begin{pmatrix} 1 \\ 0 \\ 0 \end{pmatrix}, \quad \hat{v}_2' = \begin{pmatrix} 0 \\ 1 \\ 0 \end{pmatrix}, \quad \hat{v}_3' = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

$$\mathbf{I}' \hat{v}_i' = \lambda_i \hat{v}_i'$$

(No summation on r.h.s.!))

Orthogonal

M (8)

Transformation from the original

frame: $I' = R I R^{-1}$, $\hat{V}'_i = R \hat{V}_i$

$$I = \begin{pmatrix} I_{11} & I_{12} & I_{13} \\ I_{21} & I_{22} & I_{23} \\ I_{31} & I_{32} & I_{33} \end{pmatrix}, \quad R^{-1} = R^T$$

$$I' \hat{V}'_i = \lambda_i \hat{V}'_i = \lambda_i R \hat{V}_i$$

$$R I R^{-1} R \hat{V}_i = R I \hat{V}_i$$

$$\Rightarrow I \hat{V}_i = \lambda_i \hat{V}_i$$

Eigenvalues

Eigenvectors

So, we are looking for solutions of the homogeneous linear equation

$$(I - \lambda_i I_3) \hat{V}_i = 0$$

I_3
3x3 unit matrix.

Required:

$$0 = \det$$

$$\begin{vmatrix} I_{11} - \lambda & I_{12} & I_{13} \\ I_{21} - \lambda & I_{22} - \lambda & I_{23} \\ I_{31} & I_{32} & I_{33} - \lambda \end{vmatrix}$$

M (10)

$$\Rightarrow \text{Polynomial } \lambda^3 + a_2 \lambda^2 + a_1 \lambda + a_0 = 0$$

which will have three (possibly degenerate) roots $\lambda_1, \lambda_2, \lambda_3$.

Once the roots are known, the eigenvectors follow from the equations

$$A_i \hat{v}_i = (I - \lambda_i) \hat{v}_i = 0.$$

Example: \rightarrow Classwork 14.

$$\hat{v}_i = \begin{pmatrix} v_{i1} \\ v_{i2} \\ v_{i3} \end{pmatrix},$$

$$\hat{v}_i = R \hat{v}_i$$

$$\delta_{ij} = \hat{v}_i R R \hat{v}_j = \hat{v}_i \hat{v}_j$$

Calculation of the matrix R_0

$$\hat{V}_i = \tilde{R} \hat{V}_i^1 \quad \text{implies}$$

$$\begin{pmatrix} V_1^1 \\ V_2^1 \\ V_3^1 \end{pmatrix} = \begin{pmatrix} \tilde{r}_{11} \\ \tilde{r}_{21} \\ \tilde{r}_{31} \end{pmatrix}, \quad \begin{pmatrix} V_2^1 \\ V_2^2 \\ V_2^3 \end{pmatrix} = \begin{pmatrix} \tilde{r}_{12} \\ \tilde{r}_{22} \\ \tilde{r}_{32} \end{pmatrix}$$

$$\begin{pmatrix} V_3^1 \\ V_3^2 \\ V_3^3 \end{pmatrix} = \begin{pmatrix} \tilde{r}_{13} \\ \tilde{r}_{23} \\ \tilde{r}_{33} \end{pmatrix}$$

$$\tilde{R} = \begin{pmatrix} V_1^1 & V_2^1 & V_3^1 \\ V_1^2 & V_2^2 & V_3^2 \\ V_1^3 & V_2^3 & V_3^3 \end{pmatrix}$$

$$R = \begin{pmatrix} V_1^1 & V_1^2 & V_1^3 \\ V_2^1 & V_2^2 & V_2^3 \\ V_3^1 & V_3^2 & V_3^3 \end{pmatrix}$$

First eigenvector V_i^1

(12)

Normal Modes of Vibration (Molecules)

Double Oscillator \rightarrow ~~Fig 3.10~~ p. 219,

$$\ddot{x}_1 = + \frac{k}{M} (x_2 - x_1) \quad \left| \quad x_j \text{ are displacements from zero position.} \right.$$

$$\ddot{x}_2 = + \frac{k}{m} (x_1 - x_2) + \frac{k}{m} (x_3 - x_2)$$

$$\ddot{x}_3 = + \frac{k}{M} (x_2 - x_3)$$

Ansatz:

$$x_j = a_j e^{i\omega t}, \quad j=1,2,3.$$

$\frac{1}{2}$

$$\ddot{x}_j = -\omega^2 a_j e^{i\omega t} = -\omega^2 x_j$$

Matrix form of eqns:

$$\begin{pmatrix} \frac{k}{M} - \omega^2 & -\frac{k}{M} & 0 \\ -\frac{k}{m} & \frac{2k}{m} - \omega^2 & -\frac{k}{m} \\ 0 & -\frac{k}{M} & \frac{k}{M} - \omega^2 \end{pmatrix} \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = 0$$

Not symmetric! Eigenvectors not orthogonal!

Computer assisted calculation

FORM by J.Vermaseren, version 3.3 (Jun 10 2009) Run at: Wed Apr 17 07:20:58 2013

```
* PHZ 3113 Weber & Arfken, p.219:
* 3x3 determinant for normal modes.
*  $kM=k/M$ ,  $km=k/m$ ,  $la=\omega^2$  eigenvalues wanted.
  Symbols  $kM, km, la$ ;
  Off stat;
  Local det= $(kM-la)*(2*km-la)*(kM-la)+0+0$ 
             $-(kM-la)*(-km)*(-kM)$ 
             $-(-kM)*(-km)*(kM-la)-0$ ;
* Reduction by eigenvalue 1:  $la=0$ .
  Local d1a= $-\det/la$ ;
* Remaining eigenvalues  $la_{2,3}=km+kM-\pm\sqrt{(km+kM)^2-2*kM*km-kM^2}$ .
*  $\sqrt{\quad}=-/+km$ :  $la_2=kM$ ,  $la_3=kM+2*km$ . Check:
  Local Zero= $d1a-(la-kM)*(la-kM-2*km)$ ;
  Bracket  $la$ ;
  Print;
  .sort
```

```
det =
+ la * ( - 2*kM*km - kM^2 )
+ la^2 * ( 2*km + 2*kM )
+ la^3 * ( - 1 );
```

```
d1a =
+ la * ( - 2*km - 2*kM )
+ la^2 * ( 1 )
+ 2*kM*km + kM^2;
```

```
Zero = 0;
```

```
.end
```

```
0.00 sec out of 0.00 sec
```

$\det | \dots | = 0 \Rightarrow$ Eigenvalues

$$\omega^2 = 0, \quad k/m, \quad \frac{k}{m} + 2\frac{k}{m}$$

by computer-assisted calculation.

Eigenvectors: Substitute eigenvalues into eqns.

$$\underline{\underline{\omega^2 = 0:}} \quad \left. \begin{array}{l} 1. \quad X_2 - X_1 = 0 \\ 3. \quad X_2 - X_3 = 0 \end{array} \right\} X_1 = X_2 = X_3$$

Translation.

$$\underline{\underline{\omega^2 = \frac{k}{m}:}}$$

$$1. \quad \cancel{\frac{k}{m} X_1} = \frac{k}{m} X_2 - \cancel{\frac{k}{m} X_1}$$

$$\Rightarrow \underline{\underline{X_2 = 0}}$$

$$2. \quad -\frac{k}{m} X_2 = 0 = \frac{k}{m} X_1 + \frac{k}{m} X_3$$

$$\underline{\underline{X_1 = -X_3}}$$

Center stationary. Asymmetric
 X_1, X_3 motion.

$$\omega^2 = \frac{k}{M} + 2 \frac{k}{m}$$

$$1. \quad -\frac{k}{M} x_1 - 2 \frac{k}{m} x_1 = \frac{k}{M} x_2 - \frac{k}{M} x_1$$

$$x_2 = -2 \frac{M}{m} x_1$$

$$3. \quad -\frac{k}{M} x_3 - 2 \frac{k}{m} x_3 = -2 \frac{k}{m} x_1 - \frac{k}{M} x_2$$

$$x_3 = x_1$$

The outer masses moving together.

The center mass is moving

opposite to them.

Numerical calculations:

Be careful about "ill-conditioned" systems. Rounding error may be multiplied by

$$K_{\text{clipping}} = n |A_{ij}|_{\max} |A^{-1}|_{\max}$$

\nearrow order of matrix \uparrow largest element
 \nwarrow largest element of inverse matrix.

Functions of matrices:

Defined by series expansions. E.g.,

$$\exp(A) = \sum_{n=0}^{\infty} \frac{A^n}{n!}$$

Example: σ_n Pauli matrix.

$$\exp(i\sigma_n \theta) = \sum_{n=0}^{\infty} (i)^n \frac{(\sigma_n \theta)^n}{n!}$$

M (18)

$$(\sigma^k)^n = \begin{cases} \sigma^k & \text{for } n \text{ odd} \\ 1_2 & \text{for } n \text{ even} \end{cases}$$

$$\Rightarrow \exp(i\sigma^k \theta) = \begin{cases} 1_2 \cos \theta \\ \sigma^k \sin \theta \end{cases}$$
