

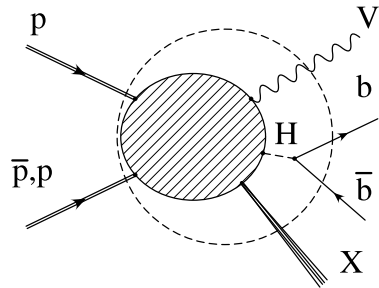
W/Z production with b -jets: a challenging background

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Università di Milano-Bicocca, Luglio 2010

- Motivations: main background to
 - WH/ZH associated production;
 - single-top production;
 - $H/A + b\bar{b}$ and other signals of new physics;
 - $t\bar{t}$ production;
 - several non-standard model signatures.
- New studies:
 - $Wb\bar{b}/Zb\bar{b}$ at NLO, b massive (F. Febres Cordero, L. R., D. Wackerath)
 - $W + 1 b$ -jet, 4FNS and 5FNS merged at NLO (J. Campbell, K. Ellis, F. Febres Cordero, F. Maltoni, L. R., D. Wackerath, S. Willenbrock)
- Outlook

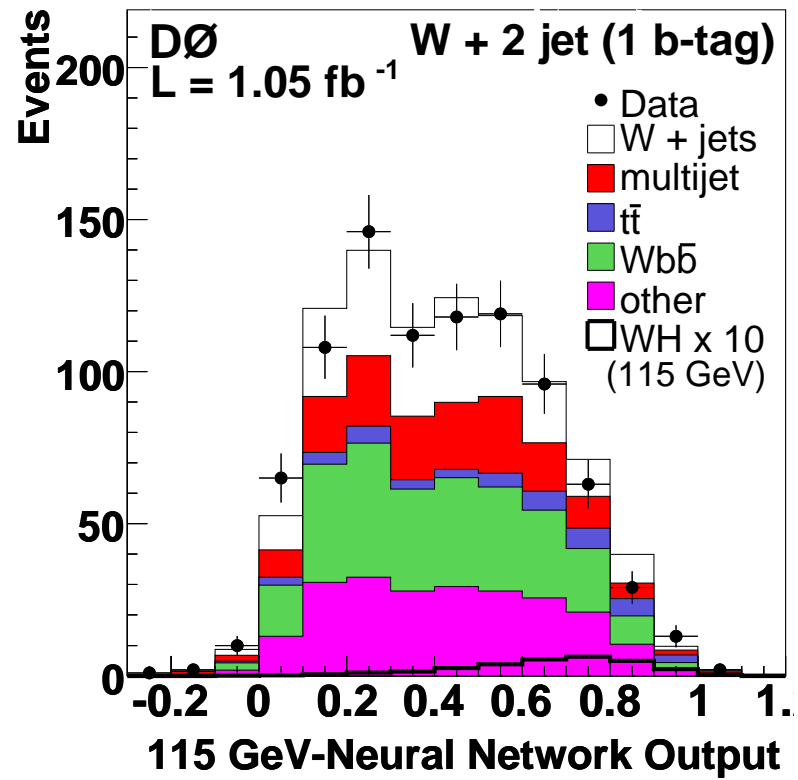
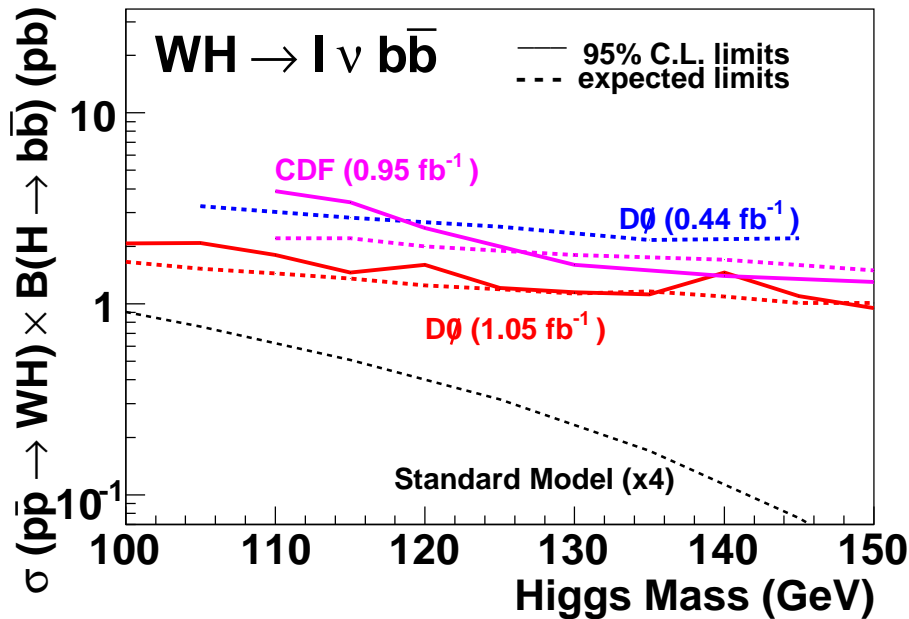
Associated production of SM Higgs with weak vector bosons



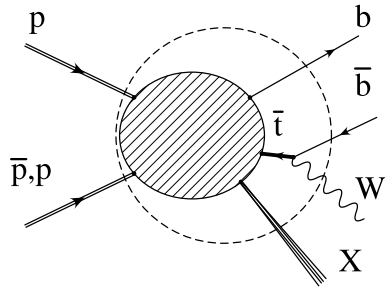
→ NNLO QCD corrections have been calculated for the signal [O'Brien, A.Djouadi and R.Harlander, 2004]

→ $O(\alpha)$ EW corrections have been calculated for the signal [M.L.Ciccolini, S.Dittmaier and M.Kramer, 2003]

→ Results for WH associated production, August 2008



SM Single-Top production



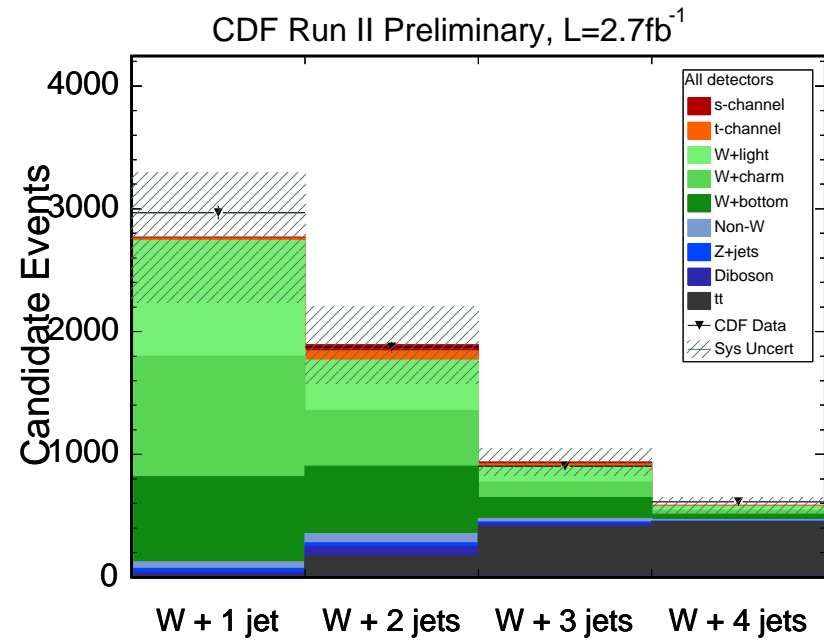
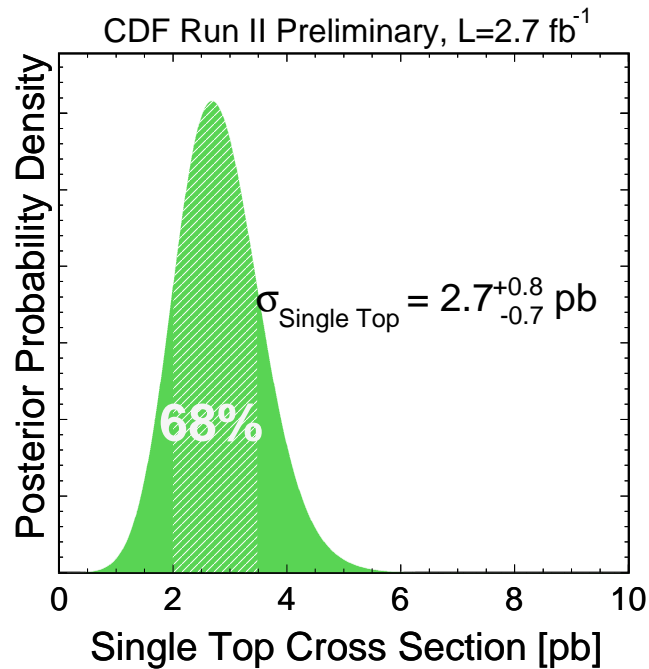
→ **NLO QCD** corrections have been thoroughly studied [T.Stelzer, Z.Sullivan and S.Willenbrock, 1998;

B.W.Harris, E.Laenen, L.Phaf, Z.Sullivan and S.Weinzierl, 2002; ...]

→ **NLO EW** corrections have been calculated for the (SM and MSSM) signal [M.Beccaria, G.Macorini,

F.M.Renard and C.Verzegnassi, 2006]

→ Summer 2008

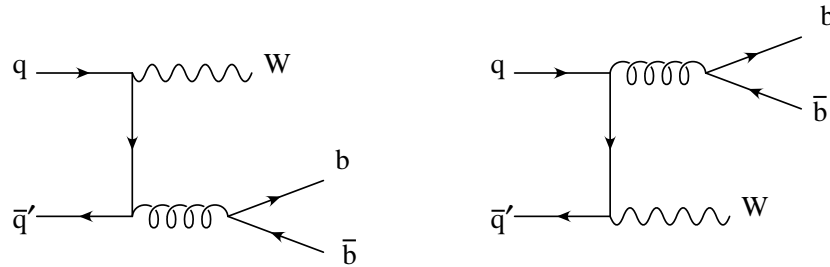


$Wb\bar{b}/Zb\bar{b}$ production at NLO, some history ...

- $V \longrightarrow 4$ partons (1-loop massless amplitudes) (Bern, Dixon, Kosower (97))
- $p\bar{p}, pp \rightarrow Vb\bar{b}$ (at NLO, 4FNS, $m_b = 0$) (Campbell, Ellis (99))
- $p\bar{p}, pp \rightarrow Vb + j$ (at NLO, 5FNS) (Campbell, Ellis, Maltoni, Willenbrock (05,07))
- $p\bar{p}, pp \rightarrow Wb\bar{b}$ (at NLO, 4FNS, $m_b \neq 0$) (Febres Cordero, L.R., Wackerth (06,09))
- $p\bar{p}, pp \rightarrow Zb\bar{b}$ (at NLO, 4FNS, $m_b \neq 0$) (Febres Cordero, L.R., Wackerth (08,09))
- $p\bar{p}, pp \rightarrow W + 1 b\text{-jet}$ (at NLO, 5FNS+4FNS with $m_b \neq 0$) (Campbell, Ellis, Febres Cordero, Maltoni, L.R., Wackerth, Willenbrock (08))

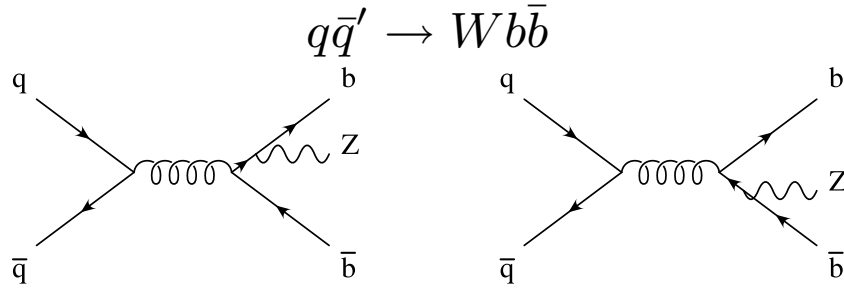
$Wb\bar{b}/Zb\bar{b}$ production with full m_b effects

LO Feynman diagrams:



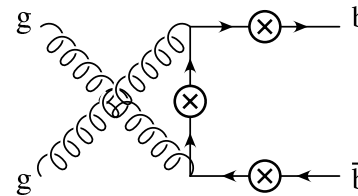
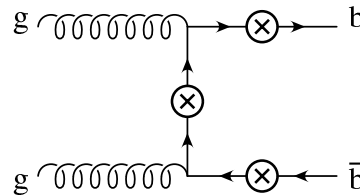
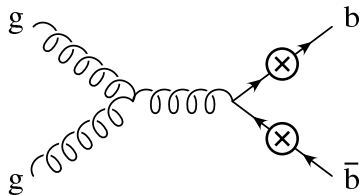
Subprocesses at LO:

- $Wb\bar{b}$: $q\bar{q}' \rightarrow Wb\bar{b}$
- $Zb\bar{b}$: $q\bar{q} \rightarrow Zb\bar{b}$ and $gg \rightarrow Zb\bar{b}$



$q\bar{q}' \rightarrow Wb\bar{b}$

$q\bar{q} \rightarrow Zb\bar{b}$



$gg \rightarrow Zb\bar{b}$

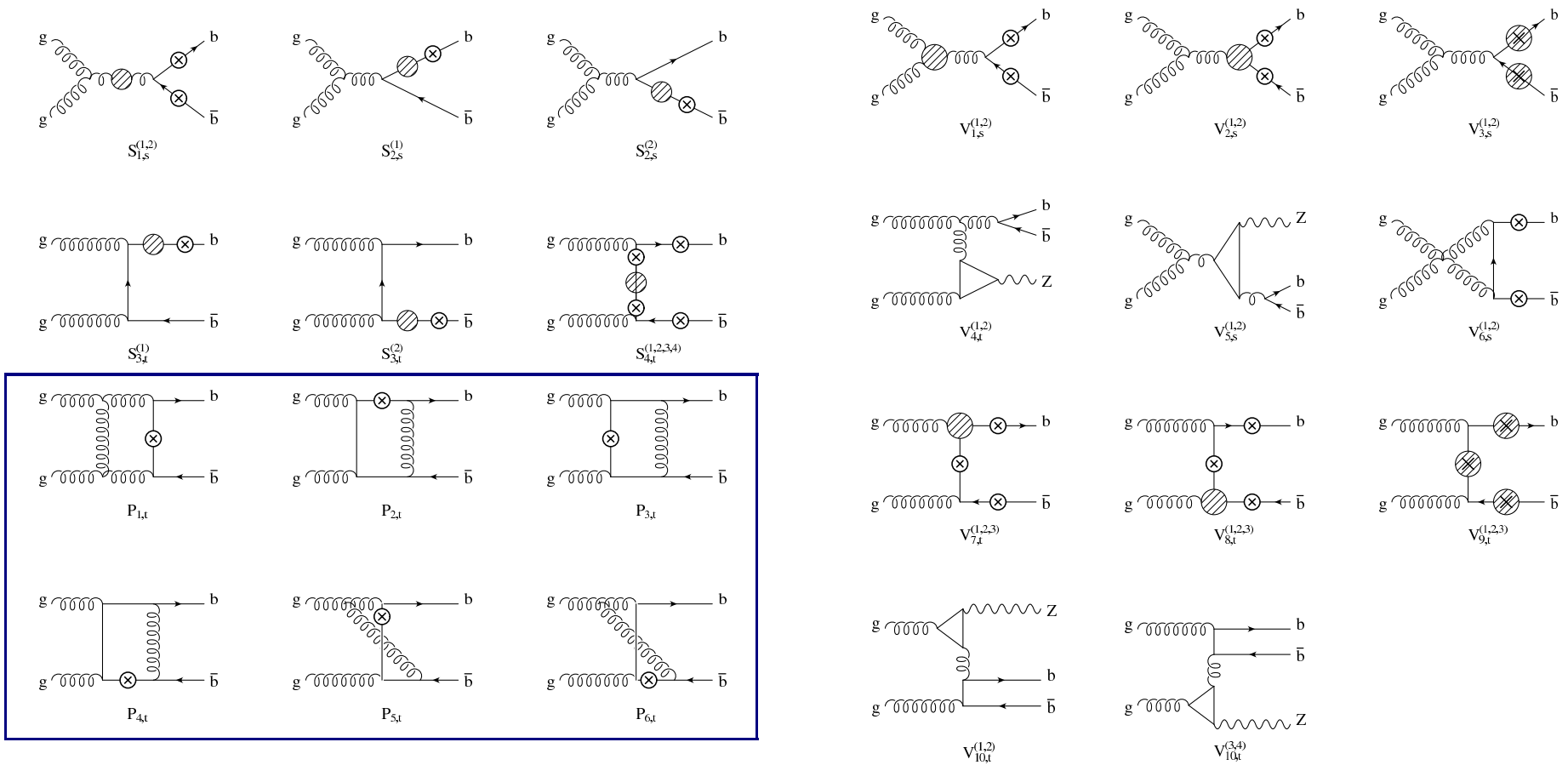
Including $O(\alpha_s)$ corrections

$$\begin{aligned}\hat{\sigma}_{ij}^{\text{NLO}}(x_1, x_2, \mu) &= \alpha_s^2(\mu) \left\{ f_{ij}^{\text{LO}}(x_1, x_2) + \frac{\alpha_s(\mu)}{4\pi} f_{ij}^{\text{NLO}}(x_1, x_2, \mu) \right\} \\ &\equiv \hat{\sigma}_{ij}^{\text{LO}}(x_1, x_2, \mu) + \delta\hat{\sigma}_{ij}^{\text{NLO}}(x_1, x_2, \mu),\end{aligned}$$

$$\delta\hat{\sigma}_{ij}^{\text{NLO}} = \hat{\sigma}_{ij}^{\text{virt}} + \hat{\sigma}_{ij}^{\text{real}}.$$

- **Virtual Corrections:** consist of one-loop diagrams interfered with corresponding LO amplitude
 - $Wb\bar{b}$: one subprocess, $q\bar{q}' \rightarrow Wb\bar{b}$
 - $Zb\bar{b}$: two subprocesses, $q\bar{q} \rightarrow Zb\bar{b}$ and $gg \rightarrow Zb\bar{b}$
- **Real Corrections:** consist of tree level diagrams with one extra parton
 - $Wb\bar{b} + k$: two subprocess, $q\bar{q}' \rightarrow Wbb + g$ and $q(\bar{q})g \rightarrow Wb\bar{b} + q'(\bar{q}')$
 - $Zb\bar{b} + k$: three subprocesses, $q\bar{q} \rightarrow Zb\bar{b} + g$, $gg \rightarrow Zb\bar{b} + g$ and $q(\bar{q})g \rightarrow Zb\bar{b} + q(\bar{q})$

Virtual corrections at a glance: the $gg \rightarrow Zb\bar{b}$ virtual diagrams.



→ Counting: 8 diagrams at LO - ~ 100 at NLO - 12 pentagons

Virtual corrections: calculating $\hat{\sigma}_{ij}^{\text{virt}}$

$$\hat{\sigma}_{ij}^{\text{virt}} = \int d(P S_3) \overline{\sum} |\mathcal{A}_{\text{virt}}(ij \rightarrow W/Z b\bar{b})|^2$$

where:

$$\overline{\sum} |\mathcal{A}_{\text{virt}}(ij \rightarrow W/Z b\bar{b})|^2 = \sum_D \overline{\sum} \left(\mathcal{A}_0 \mathcal{A}_D^\dagger + \mathcal{A}_0^\dagger \mathcal{A}_D \right) = \sum_D \overline{\sum} 2\text{Re} \left(\mathcal{A}_0 \mathcal{A}_D^\dagger \right) .$$

- Use **dimensional regularization** to regularize UV and IR divergencies.
- UV divergencies are canceled by a suitable set of **counterterms**.
- Calculate each diagram as linear combination of **Dirac structures** with coefficients that depend on both **tensor and scalar integrals**.
- Tensor integrals reduced analytically to scalar integrals and organized to avoid spurious divergences due to appearance of inverse power of Gram Determinant.
- IR divergencies will cancel with $\hat{\sigma}_{ij}^{\text{real}}$.

Checking boxes and pentagons using unitarity methods.

The one-loop amplitude can be written as

$$\mathcal{A}_{\text{virt}} = \sum_i d_i I_4^i + \sum_i c_i I_3^i + \sum_i b_i I_2^i + \sum_i a_i I_1^i + R$$

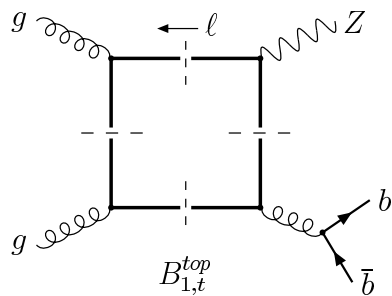
→ tadpoles, bubbles and vertices are easy in FD's language;

→ boxes and pentagons are the real hurdle (tensor integrals up to rank 4)



I_4^i scalar 4-point functions originate from box and pentagon diagrams.

Calculating d_i with unitarity methods is a powerful check!



easy using quadrupole cuts!

Britto, Cachazo, Feng

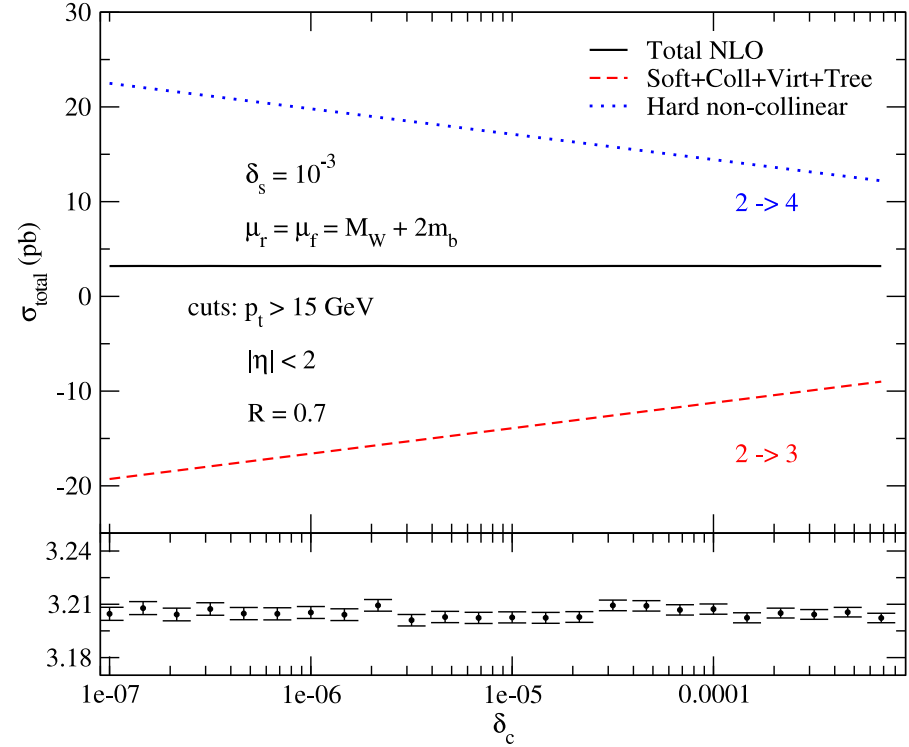
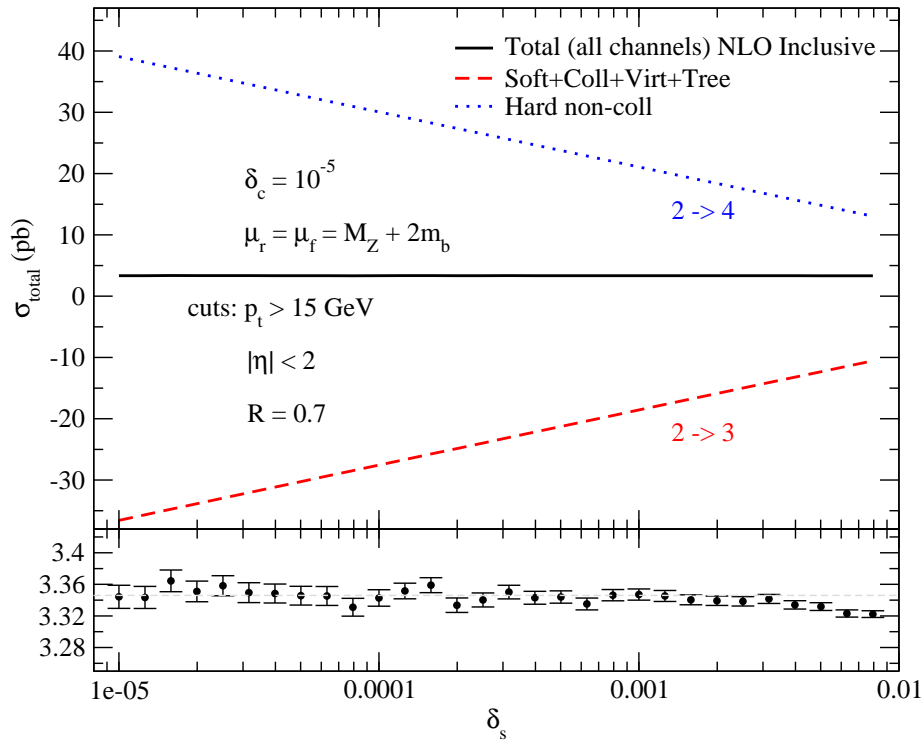
Bern, Dixon, Kosower

Real corrections: calculating $\hat{\sigma}_{ij}^{\text{real}}$

$$\hat{\sigma}_{ij}^{\text{real}} = \int d(PS_4) \overline{\sum} |\mathcal{A}_{\text{real}}(ij \rightarrow W/Z b\bar{b} + k)|^2$$

- IR divergencies associated with the integration over the PS of the extra parton, can be extracted using the so called **Phase Space Slicing (PSS)** method with *two cutoffs*.
- PSS with two cutoffs uses two unphysical parameters, δ_s and δ_c to isolate soft and collinear divergent regions, where IR singularities are extracted analytically.
- Same soft/collinear structure as $Ht\bar{t}/Hb\bar{b}$, tested against one-cutoff PSS and dipole subtraction method.
- Physical quantities are independent of δ_s and δ_c , for small enough values of these parameters.

Check independence of the total cross section of δ_s and δ_c cuts



δ_s -dep. for the $Zb\bar{b}$ total cross section

δ_c -dep. for the $Wb\bar{b}$ total cross section

→ Cross section independent of unphysical cutoffs δ_s and δ_c

→ Results obtained fixing $\delta_s = 10^{-3}$ and $\delta_c = 10^{-5}$

$W/Z + 2 b\text{-jets}$

Febres Cordero, L.R., Wackerath (06-09)

- We use the k_T jet algorithm with $R = 0.7$ and study two cases:
 - **Inclusive Cross Section:** events with two $(b + \bar{b})$ or three $(b + \bar{b} + j)$ jets resolved contribute to the cross section.
 - **Exclusive Cross Section:** only events with two $(b + \bar{b})$ jets resolved contribute to the cross section.

Same convention used by MCFM (used to obtain the results for $m_b = 0$).

- **b -jet kinematical cuts:**
 - **Transverse momentum** of the b -jets: $p_t > p_{t, \min}$ (15-25 GeV) for both b and \bar{b} jets.
 - **Pseudorapidity:** $|\eta| < \eta_{\max}$ (2-2.5) for both b and \bar{b} jets.
- **PDF:** for LO results we use 1-loop evolution of α_s and CTEQ6L1, while for NLO results 2-loop evolution of α_s and CTEQ6M.

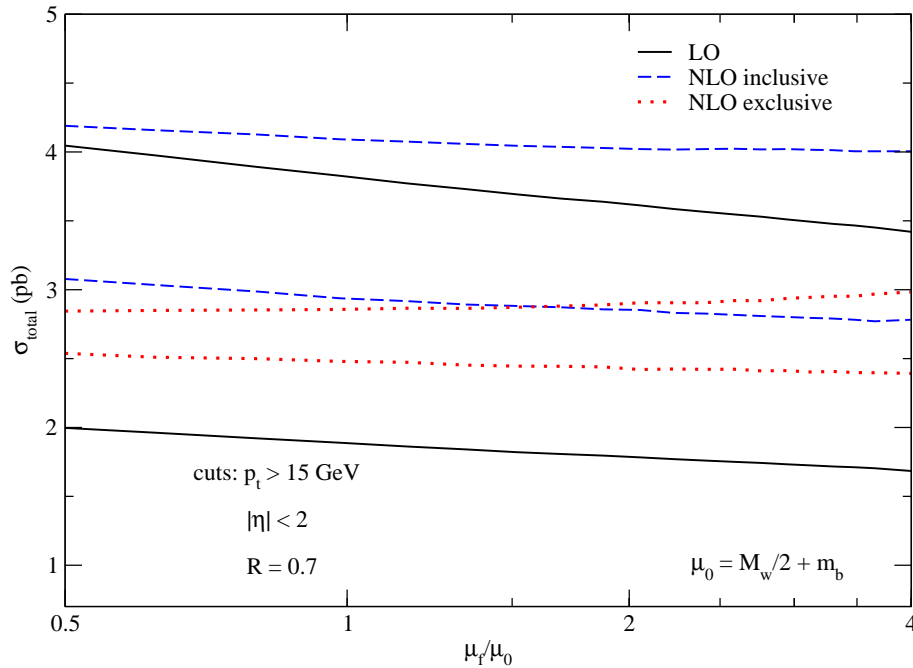
Tevatron: summary of LO and NLO total cross sections

massive and massless calculation, setting $\mu_r = \mu_f = M_V + 2m_b$ ($V = W, Z$).

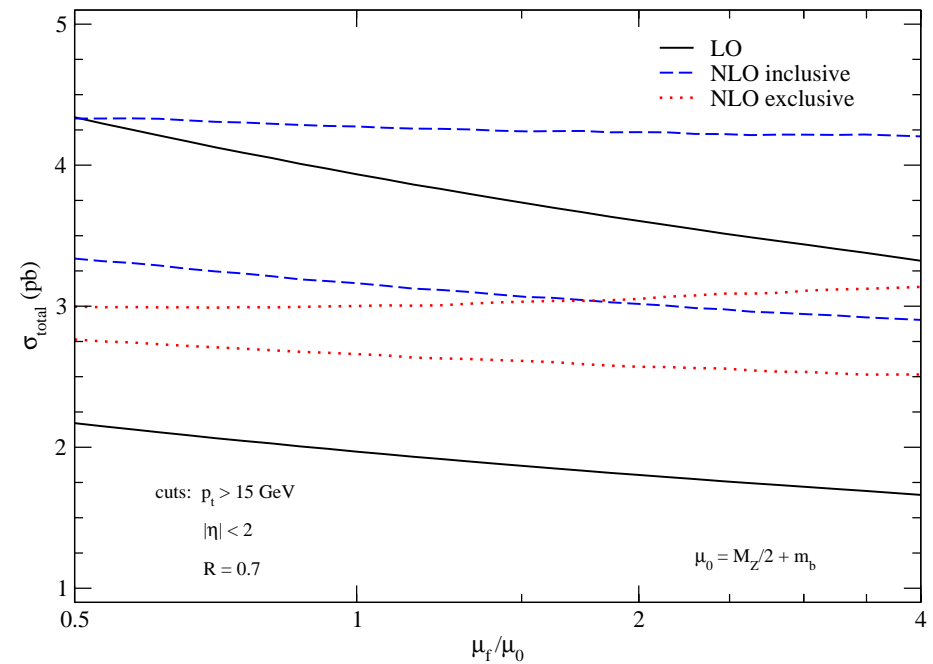
Cross Section, $Wb\bar{b}$	$m_b \neq 0$ (pb) [ratio]	$m_b = 0$ (pb) [ratio]
σ^{LO}	2.20[-]	2.38[-]
σ^{NLO} inclusive	3.20[1.45]	3.45[1.45]
σ^{NLO} exclusive	2.64[1.2]	2.84[1.2]

Cross Section, $Zb\bar{b}$	$m_b \neq 0$ (pb) [ratio]	$m_b = 0$ (pb) [ratio]
σ^{LO}	2.21[-]	2.37[-]
σ^{NLO} inclusive	3.34[1.51]	3.64[1.54]
σ^{NLO} exclusive	2.75[1.24]	3.01[1.27]

Tevatron: scale dependence and theoretical uncertainty at NLO



$Wb\bar{b}$: Tevatron (PRD 74 (2006) 034007)

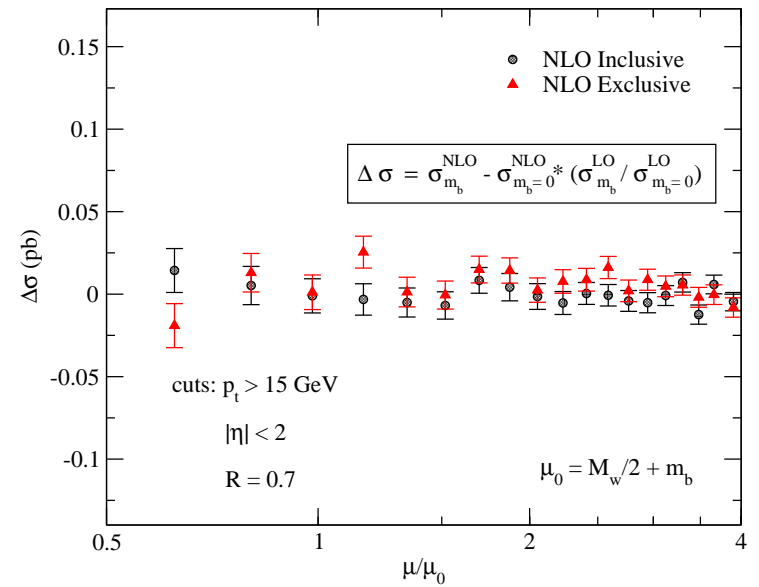
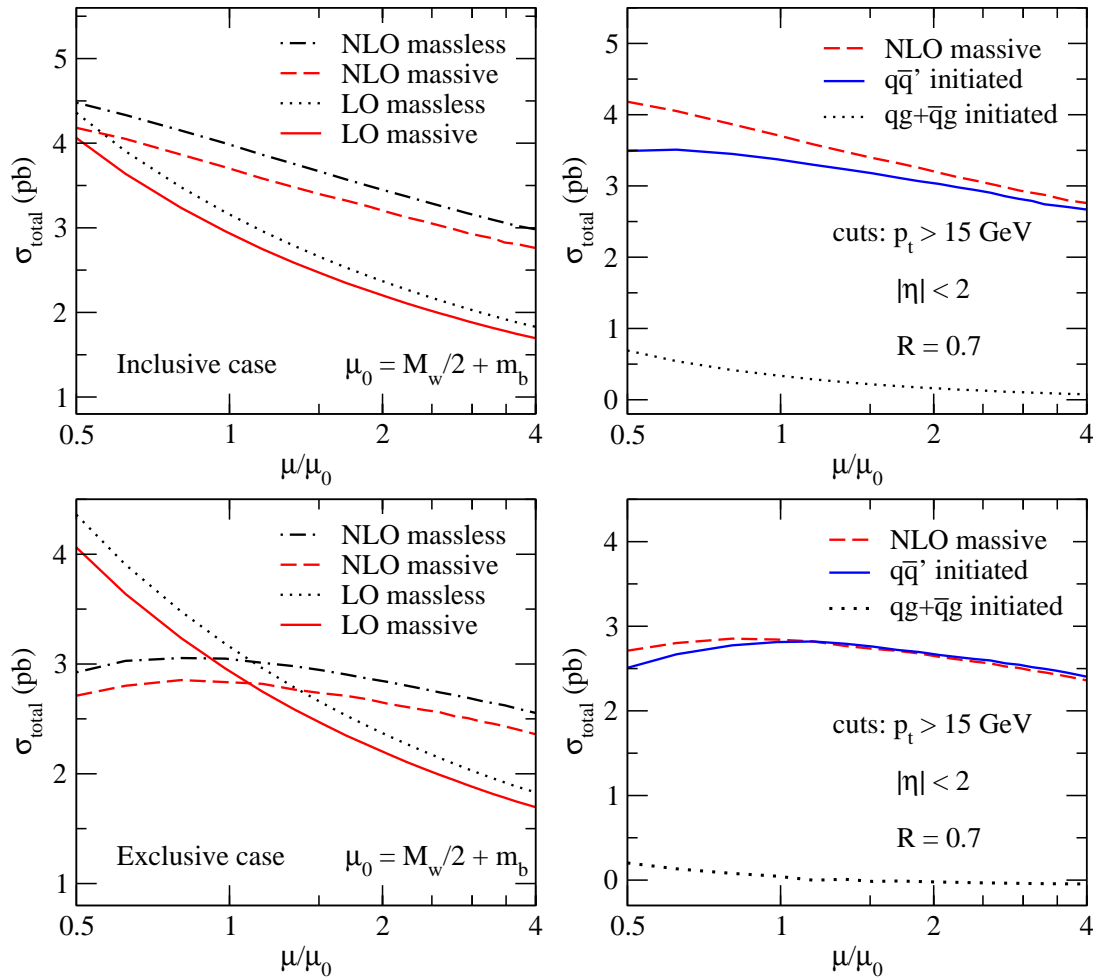


$Zb\bar{b}$: Tevatron (PRD 78 (2008) 074014)

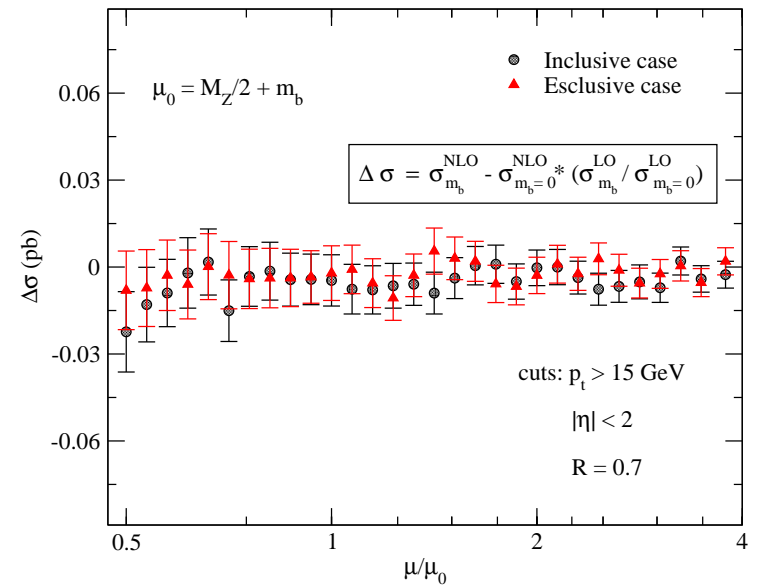
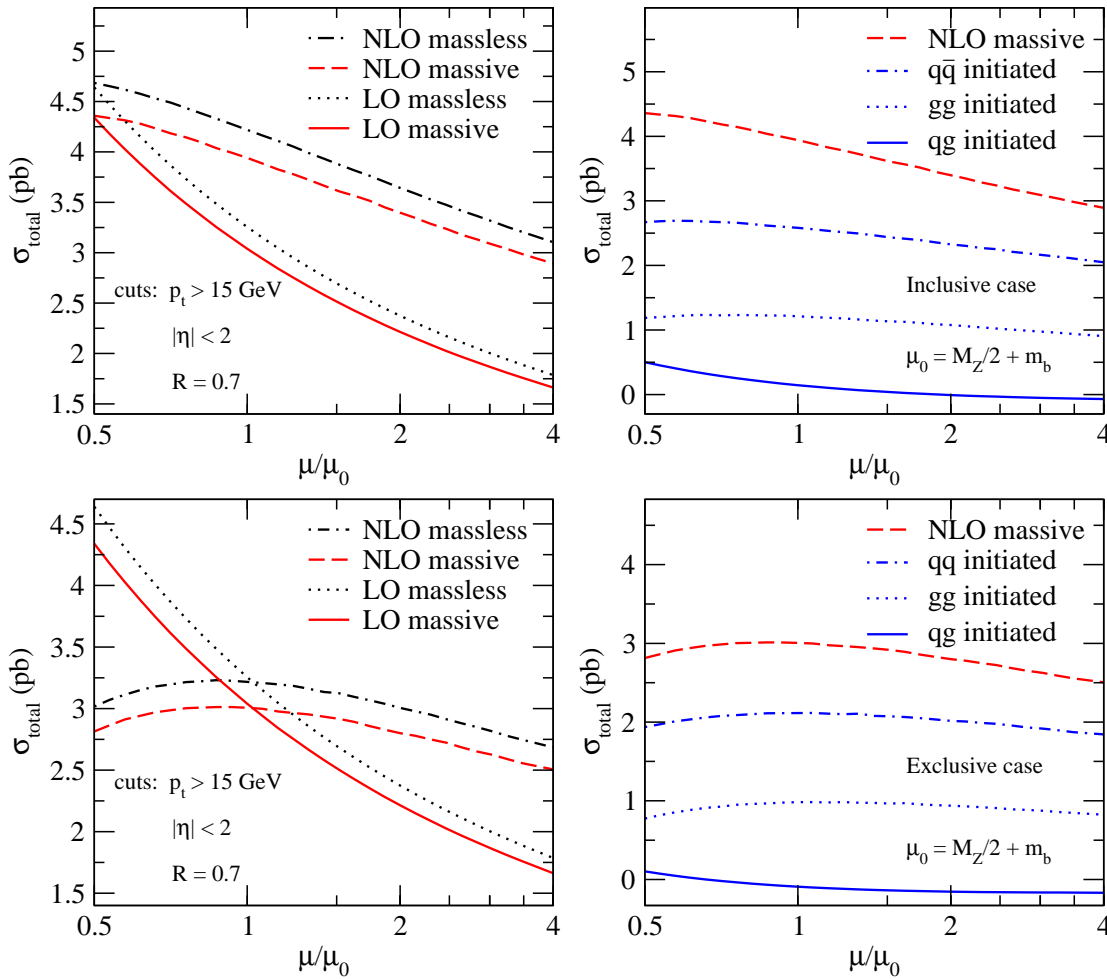
→ Bands obtained by varying both μ_R and μ_F between $\mu_0/2$ and $4\mu_0$ (with $\mu_0 = m_b + M_V/2$ ($V = W, Z$)).

- LO uncertainty $\sim 40\%$.
- Inclusive NLO uncertainty $\sim 20\%$.
- Exclusive NLO uncertainty $\sim 10\%$.

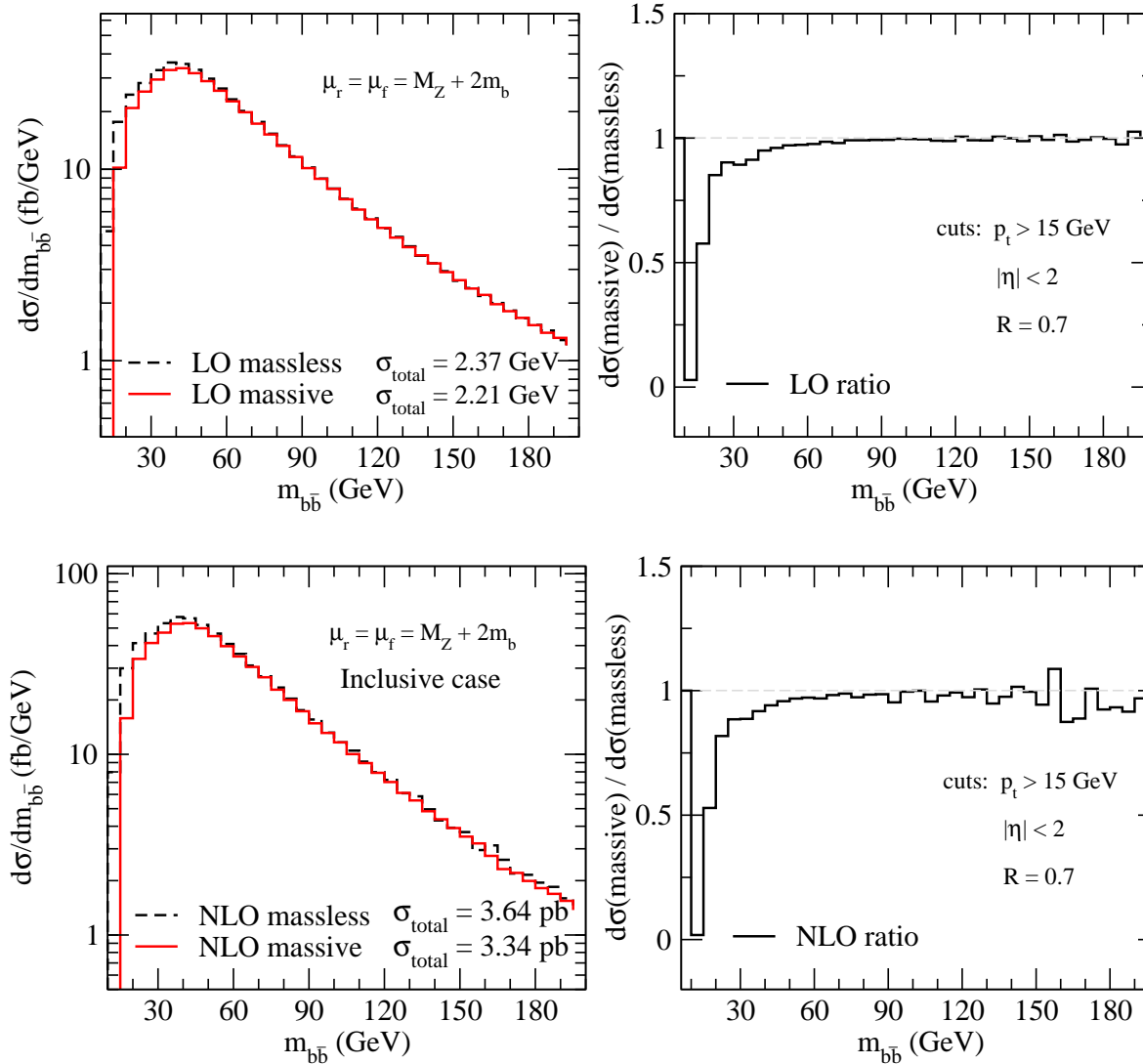
$Wb\bar{b}$, scale dependence: LO vs NLO and massless vs massive



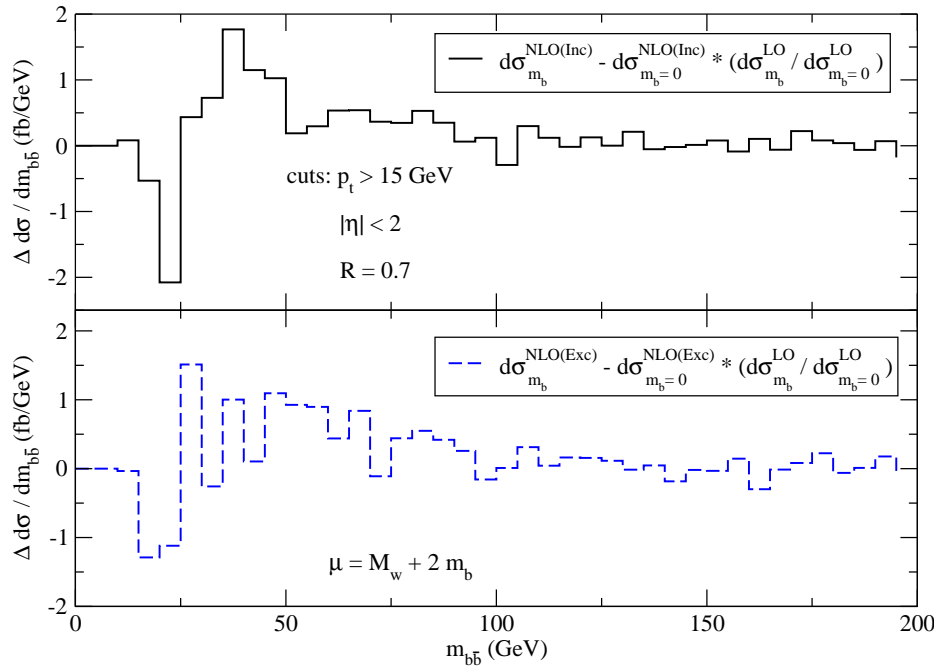
$Zb\bar{b}$, scale dependence: LO vs NLO and massless vs massive



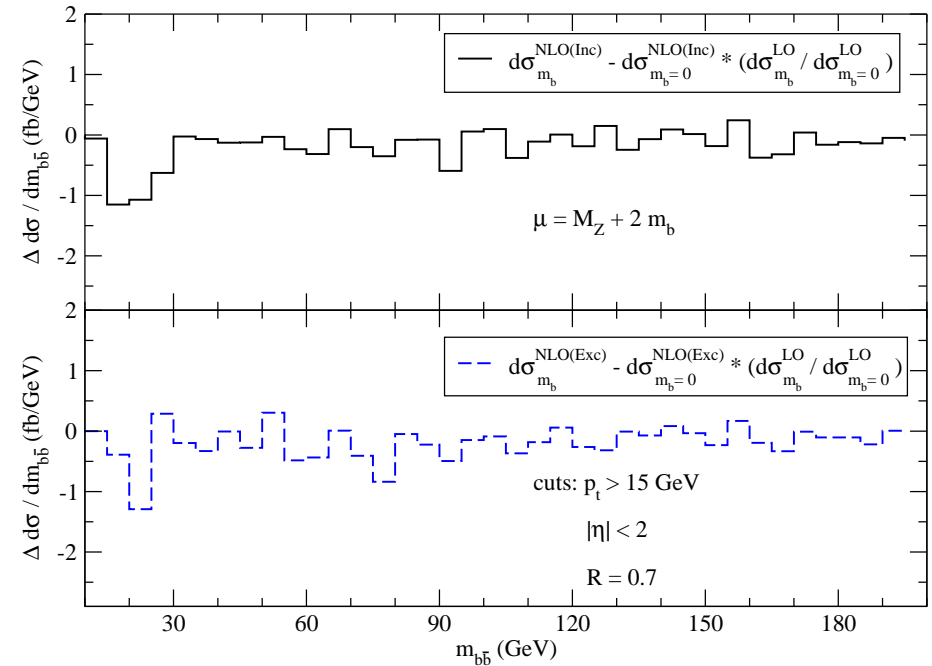
$Zb\bar{b}$: $m_{b\bar{b}}$ distributions, massive vs massless



$Wb\bar{b}/Zb\bar{b}$, $m_{b\bar{b}}$ distributions: testing rescaling LO \rightarrow NLO



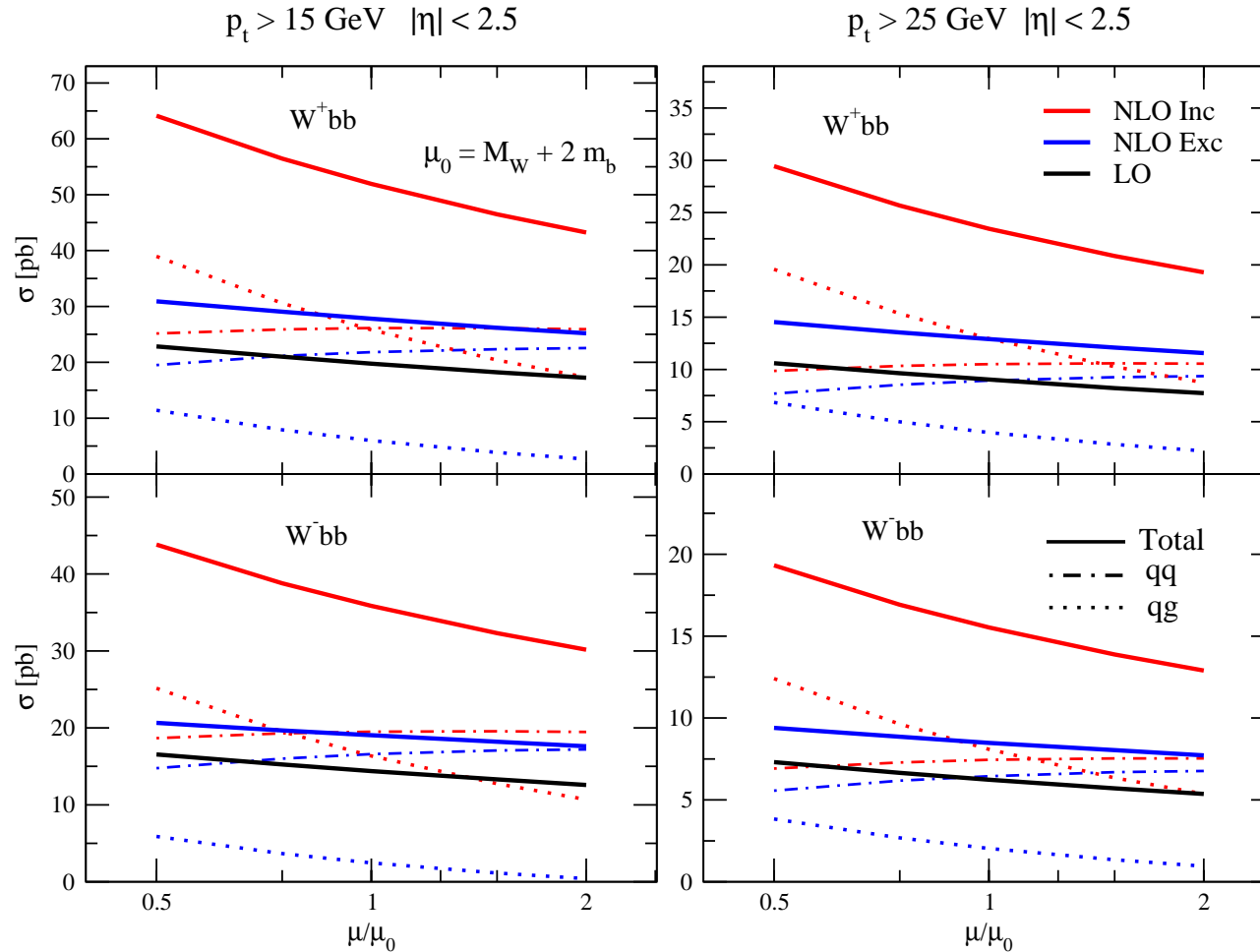
$Wb\bar{b}$: PRD 74 (2006) 034007



$Zb\bar{b}$: PRD 78 (2008) 074014

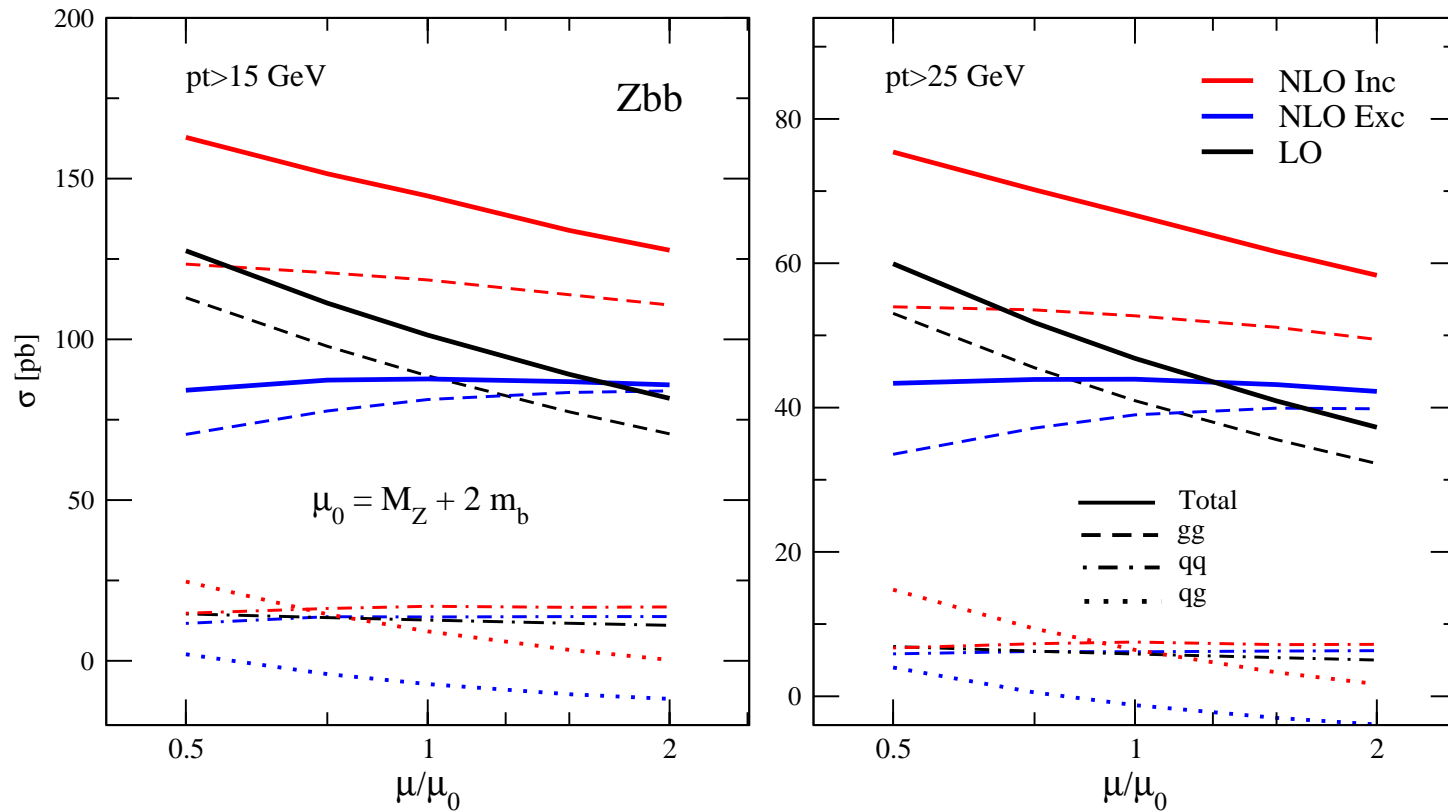
Clear effect in the low $m_{b\bar{b}}$ invariant mass region.

LHC: scale dependence and theoretical uncertainty at NLO



(PRD 80:034015, 2009)

- NLO corrections very large, particularly for inclusive production;
- large NLO scale-dependence (LO: 30%, NLO_{inc} : 50%, NLO_{exc} : 20%), induced by the opening of the $qg(\bar{q}g) \rightarrow Wb\bar{b} + q'(\bar{q}')$ channel;
- theoretical uncertainty not only given by scale-dependence!

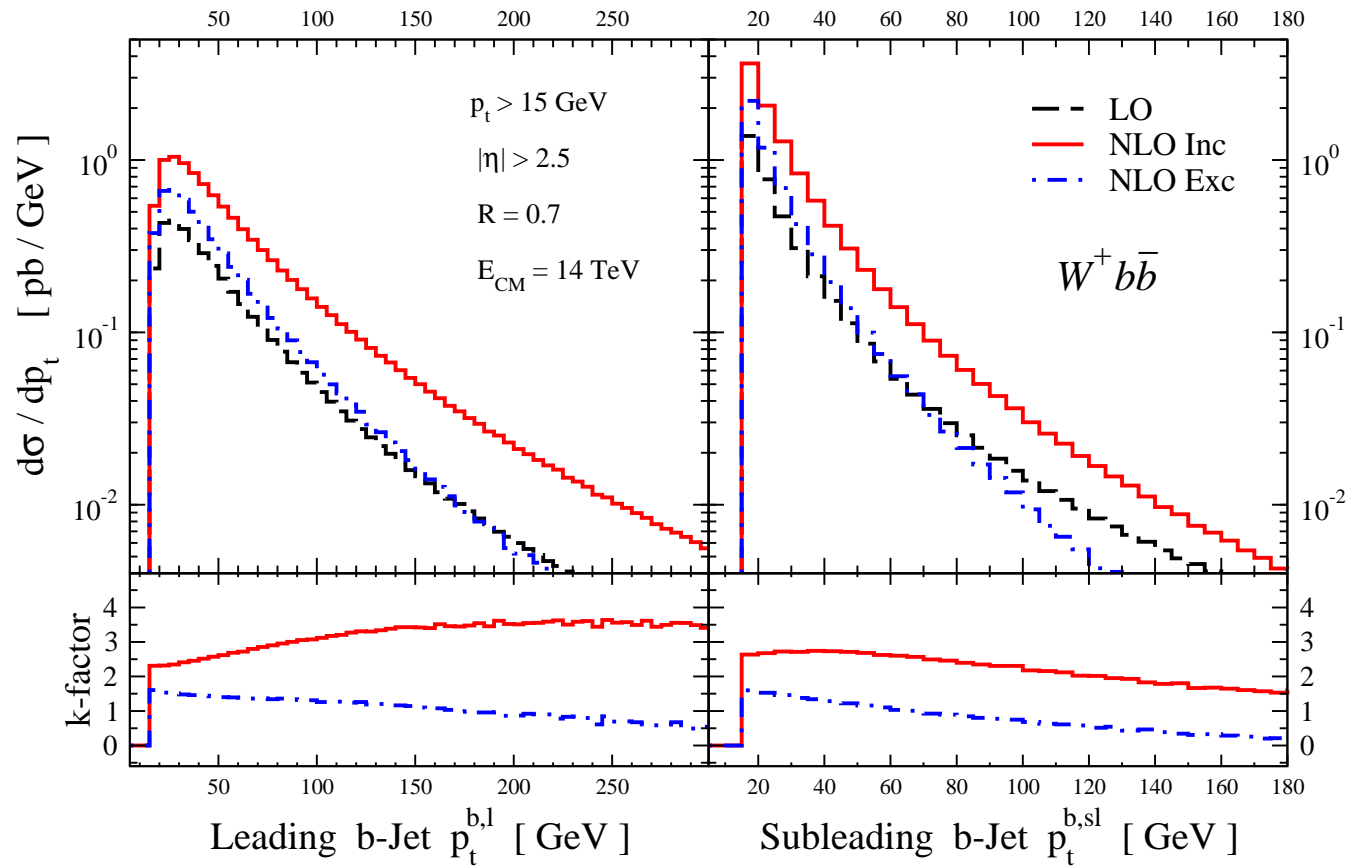


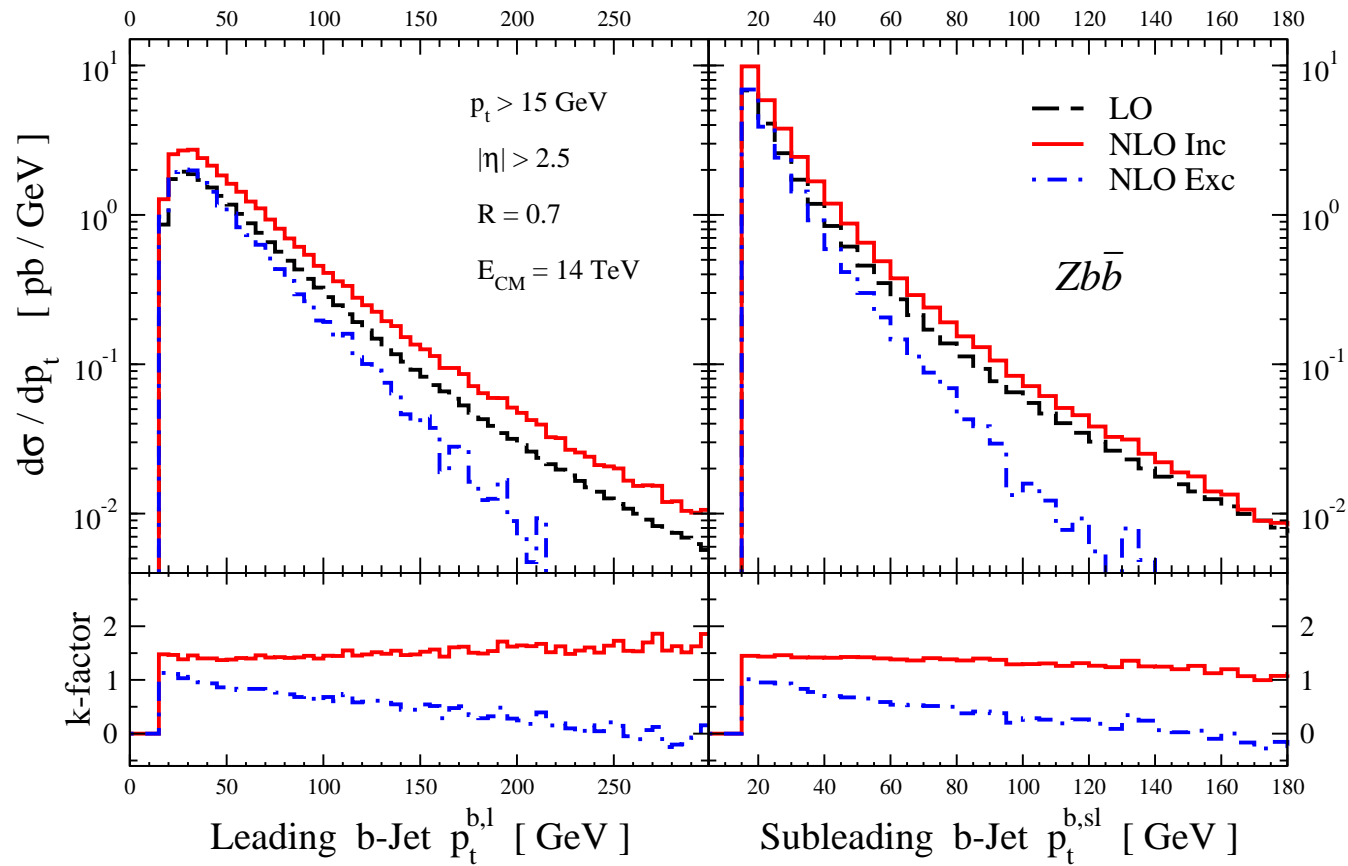
- NLO corrections still large, particularly for inclusive production;
- more moderate NLO scale-dependence (LO: 50%, NLO_{inc} : 30%, NLO_{exc} : 5%): $qg(\bar{q}g) \rightarrow Zb\bar{b} + q(\bar{q})$ channel not as dominant.

	$W^+ b\bar{b}$			
	$m_b \neq 0$	$m_b = 0$	$m_b \neq 0$	$m_b = 0$
	$p_T^b > 15 \text{ GeV}$		$p_T^b > 25 \text{ GeV}$	
	$\sqrt{s} = 10 \text{ TeV}$			
σ_{LO} (pb)	$14.4^{+2.6}_{-2.1}$	$15.5^{+2.7}_{-2.2}$	$6.49^{+1.3}_{-1.0}$	$6.67^{+1.3}_{-1.1}$
$\sigma_{NLO,inc}$ (pb)	$33.6^{+7.8}_{-5.6}$	$36.4^{+8.1}_{-6.2}$	$14.6^{+3.6}_{-2.5}$	$15.1^{+3.6}_{-2.7}$
$\sigma_{NLO,exc}$ (pb)	$18.6^{+1.7}_{-1.6}$	$20.3^{+1.7}_{-1.8}$	$8.37^{+0.84}_{-0.77}$	$8.67^{+0.85}_{-0.87}$
	$\sqrt{s} = 14 \text{ TeV}$			
σ_{LO} (pb)	$19.8^{+3.1}_{-2.5}$	$21.3^{+3.2}_{-2.7}$	$9.02^{+1.6}_{-1.3}$	$9.26^{+1.6}_{-1.3}$
$\sigma_{NLO,inc}$ (pb)	$51.9^{+12}_{-8.7}$	$56.3^{+13}_{-9.6}$	$23.4^{+6.0}_{-4.2}$	$24.3^{+5.9}_{-4.5}$
$\sigma_{NLO,exc}$ (pb)	$27.8^{+3.1}_{-2.5}$	$30.4^{+3.5}_{-2.8}$	$12.9^{+1.6}_{-1.3}$	$13.4^{+1.6}_{-1.5}$

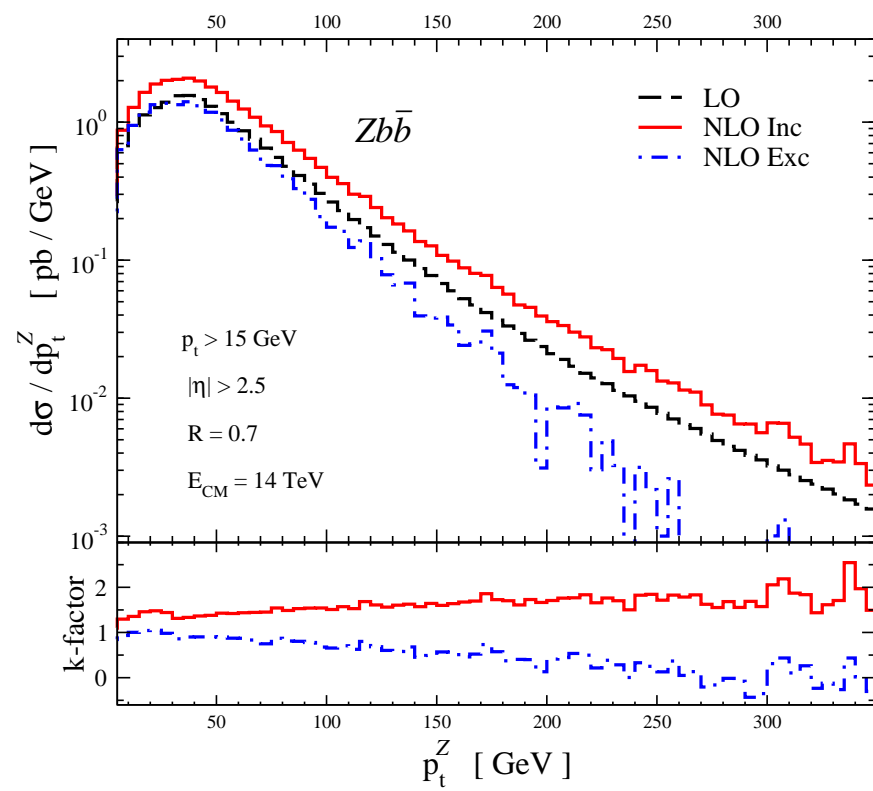
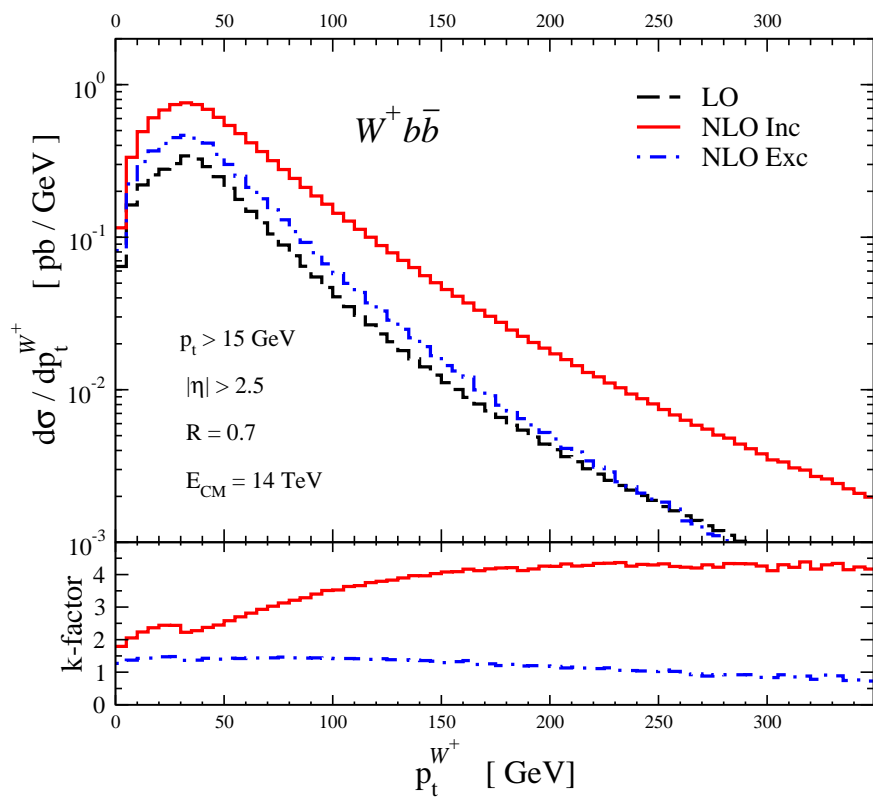
	$Zb\bar{b}$			
	$p_T^b > 15 \text{ GeV}$		$p_T^b > 25 \text{ GeV}$	
	$m_b \neq 0$	$m_b = 0$	$m_b \neq 0$	$m_b = 0$
	$\sqrt{s} = 10 \text{ TeV}$			
σ_{LO} (pb)	55.1^{+16}_{-12}	57.6^{+18}_{-13}	$24.6^{+7.6}_{-5.4}$	$25.1^{+8.4}_{-5.9}$
$\sigma_{NLO,inc}$ (pb)	82.5^{+12}_{-11}	84.5^{+14}_{-12}	$36.0^{+3.9}_{-4.6}$	$36.1^{+6.4}_{-5.2}$
$\sigma_{NLO,exc}$ (pb)	$52.1^{+0.0}_{-1.7}$	$53.5^{+0.2}_{-2.4}$	$24.6^{+0.3}_{-1.2}$	$24.7^{+0.3}_{-1.6}$
	$\sqrt{s} = 14 \text{ TeV}$			
σ_{LO} (pb)	101^{+26}_{-20}	106^{+30}_{-22}	$46.8^{+13.1}_{-9.6}$	$46.8^{+12.7}_{-9.9}$
$\sigma_{NLO,inc}$ (pb)	145^{+20}_{-17}	148^{+24}_{-19}	$66.6^{+8.8}_{-8.3}$	$66.1^{+10.5}_{-9.1}$
$\sigma_{NLO,exc}$ (pb)	$88.4^{+0.0}_{-3.0}$	$90.0^{+0.0}_{-1.6}$	$43.7^{+0.0}_{-1.6}$	$43.5^{+0.4}_{-1.9}$

LHC: some interesting NLO distributions ...

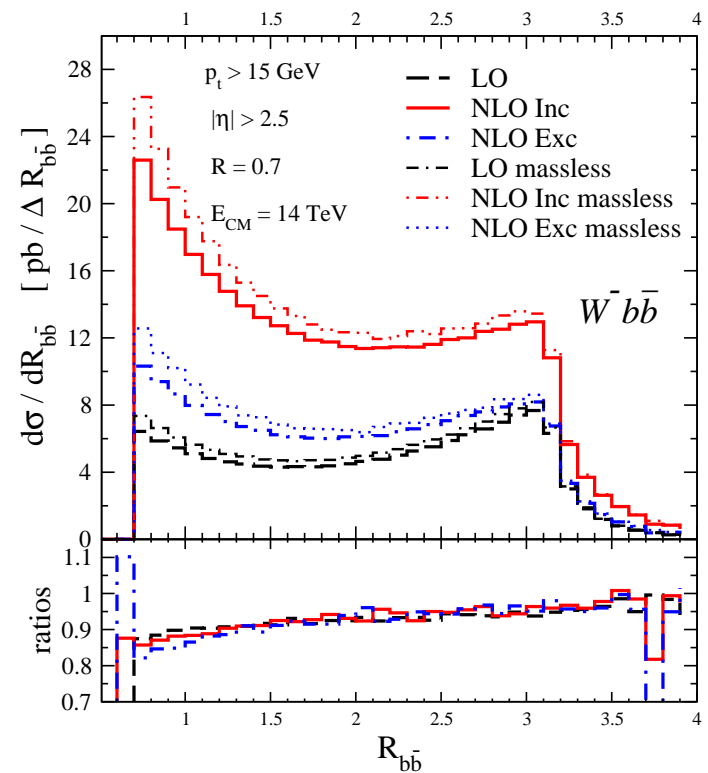
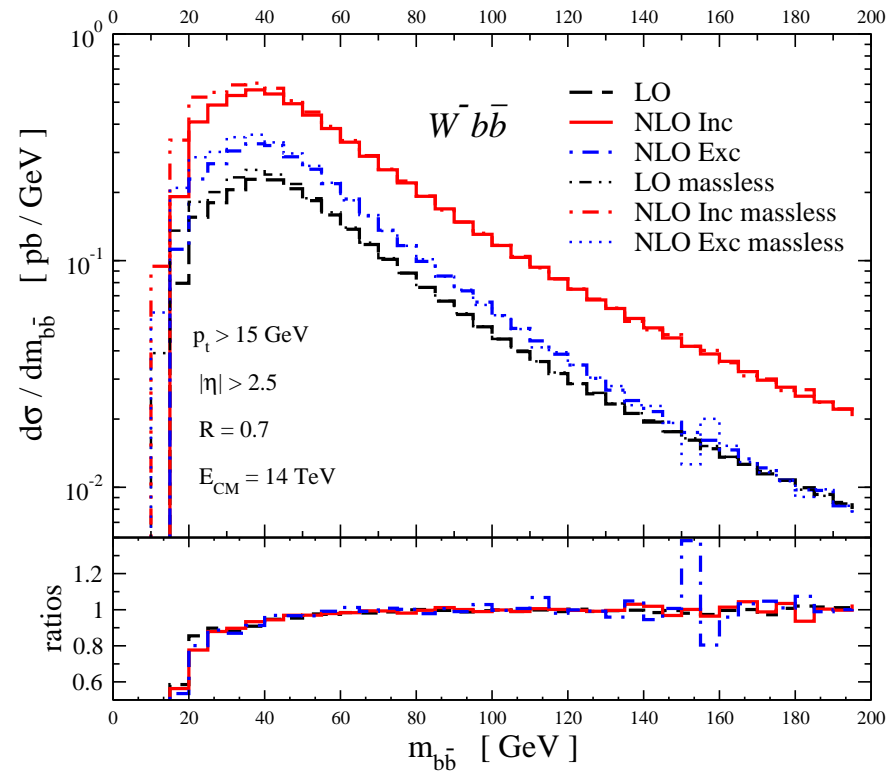


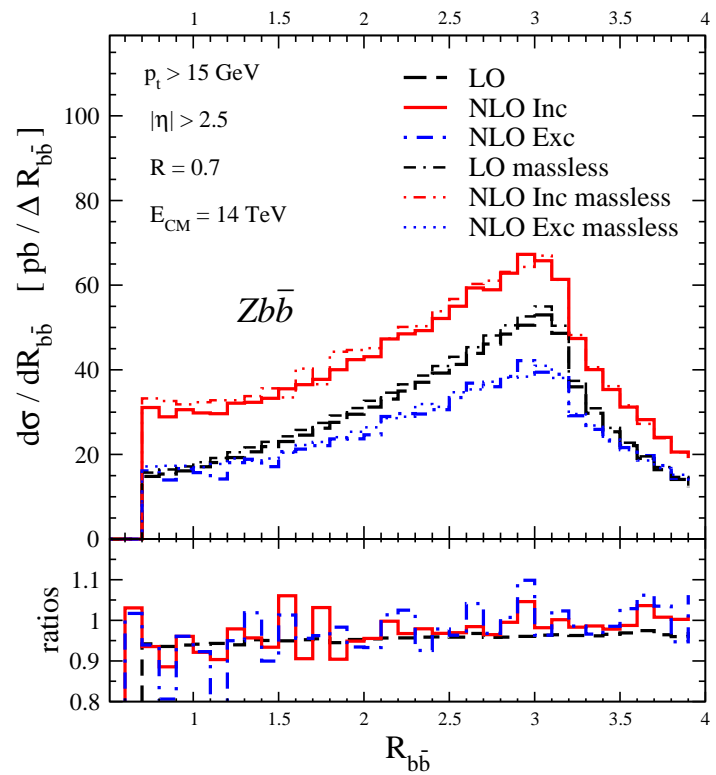
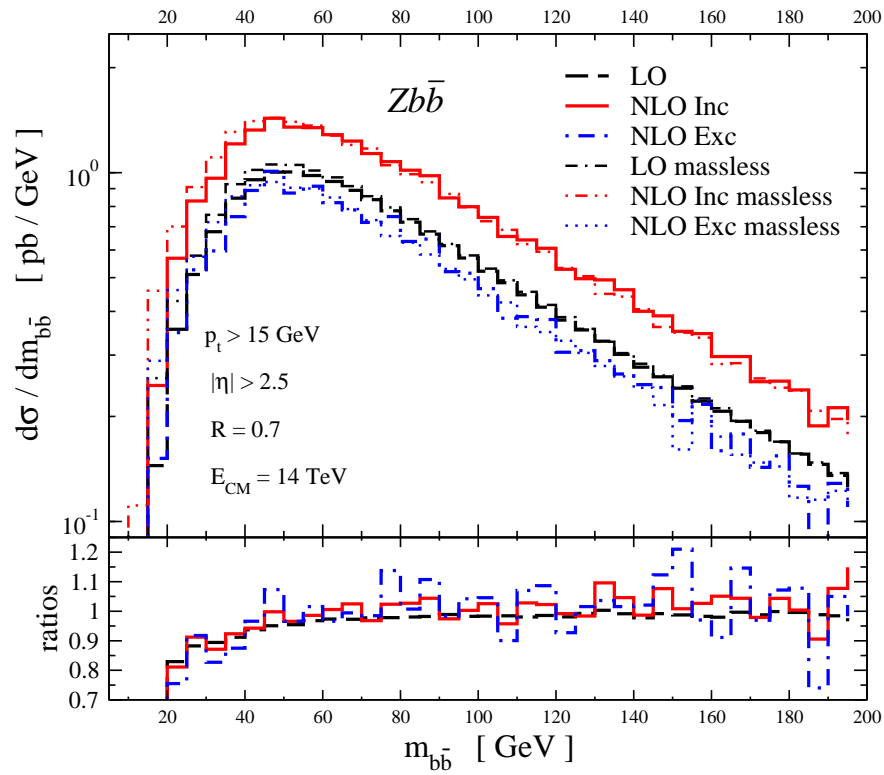


- NLO distributions cannot be rescaled from LO ones via a K-factor;
- large corrections in particular for $Wb\bar{b}$ and in the inclusive case.



Comparison with massless b -quark results ...



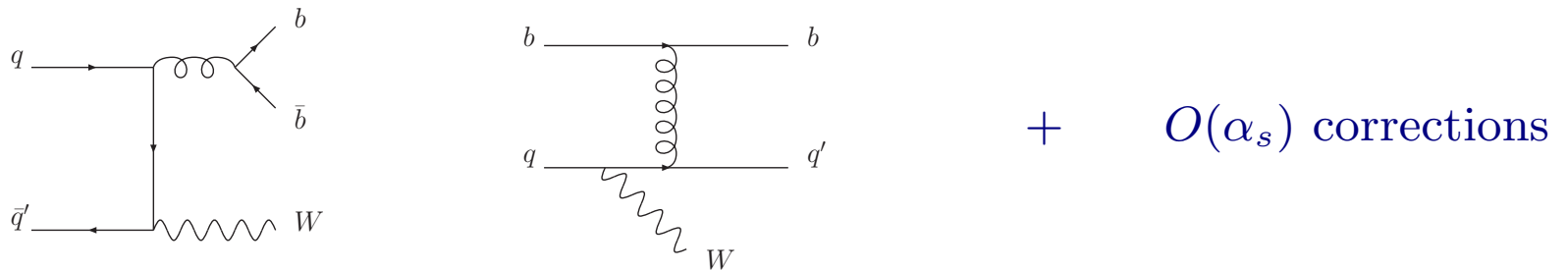


$W + 1 b\text{-jet}$

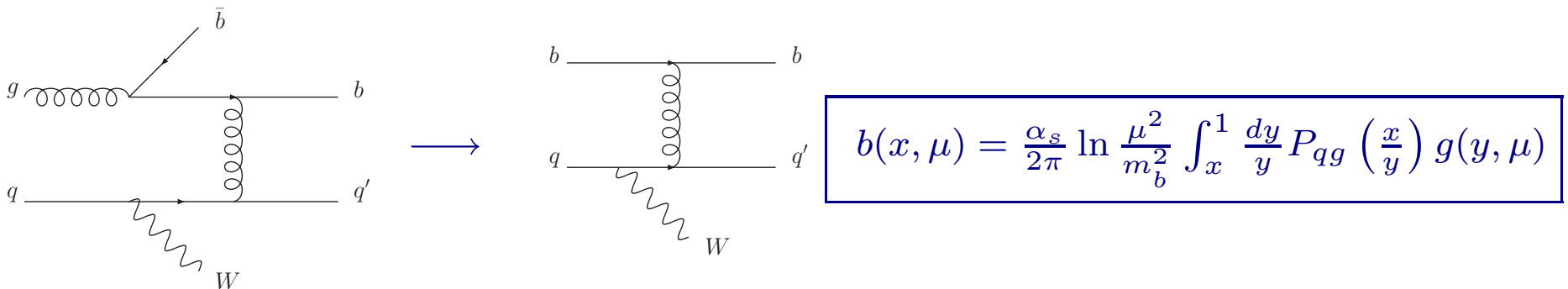
Campbell, Ellis, Febres Cordero, Maltoni, L.R., Wackerroth, Willenbrock

(PRD 79:034023, 2009)

Consistently combine 4FNS ($m_b \neq 0$) and 5FNS ($m_b = 0$) at NLO in QCD:

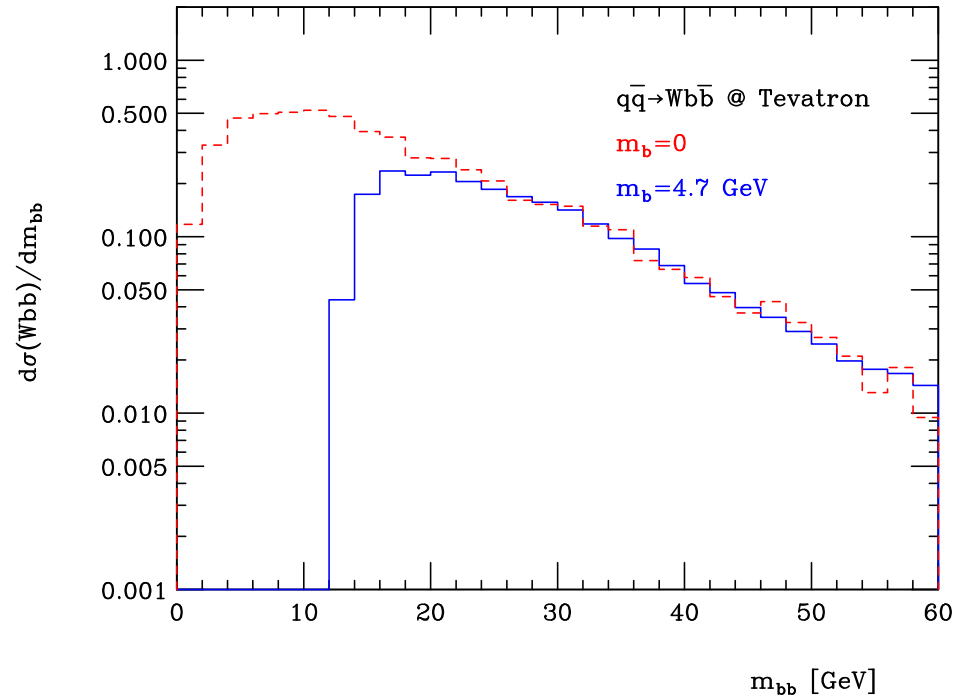


1. $q\bar{q}' \rightarrow Wb\bar{b}$ at tree level and one loop ($m_b \neq 0$)
2. $q\bar{q}' \rightarrow Wb\bar{b}g$ at tree level ($m_b \neq 0$)
3. $bq \rightarrow Wbq'$ at tree level and one loop ($m_b = 0$)
4. $bq \rightarrow Wbq'g$ and $bg \rightarrow Wbq'\bar{q}$ at tree level ($m_b = 0$)
5. $gq \rightarrow Wb\bar{b}q'$ at tree level ($m_b \neq 0$) \rightarrow avoiding double counting:



\rightarrow indeed: a fully consistent NLO 5FNS calculation (S-ACOT scheme).

- improved scale dependence: NLO corrections to $gq \rightarrow Wb\bar{b}q'$ partially included;
- need to keep $m_b \neq 0$ for final state b quarks (one b quark has low p_T)



- four signatures studied: exclusive/inclusive, with single and double- b jets, using $p_T^j > 15$ GeV, $|\eta^j| < 2 - 2.5$, cone algorithm with $\Delta R = 0.7$:
 - $Wb, W(b\bar{b})$ (exclusive)
 - Wb and $Wb + j, W(b\bar{b})$ and $W(b\bar{b}) + j$ (inclusive)
 which can be combined to obtain different backgrounds, ...
- both contributions play important complementary roles (Tevatron/LHC, inclusive/exclusive);

- NLO results at a glance:

	Exclusive cross sections (pb)	
Collider	Wb	$W(b\bar{b})$
TeV $W^+ (= W^-)$	$8.02+0.62[-0.05]=8.64$	$3.73-0.02[-0.02]=3.71$
LHC W^+	$40.0+48.4[22.6]=88.4$	$22.7+11.7[11.7]=34.4$
LHC W^-	$29.8+29.4[12.6]=59.2$	$17.2+6.5[6.5]=23.7$

	Inclusive cross sections (pb)	
Collider	$Wb + X$	$W(b\bar{b}) + X$
TeV $W^+ (= W^-)$	$11.77+2.40[0.77]=14.17$	$4.17+0.39[0.39]=4.56$
LHC W^+	$53.6+136.1[68.9]=189.7$	$25.1+35.9[35.9]=61.0$
LHC W^-	$39.3+88.2[44.6]=127.5$	$18.9+23.6[23.6]=42.5$

- first number: Processes 1 + 2 (pure 4FNS)
- second number: Processes 3 + \dots + 5 (pure 5FNS plus $qg \rightarrow Wb\bar{b} + q'$)
- number in square brackets: Process 5 alone ($qg \rightarrow Wb\bar{b} + q'$)

Comparison with CDF measurement: a puzzle?

CDF Note 9321:

$$\sigma_{b\text{-jet}}(W + b\text{jets}) \cdot Br(W \rightarrow l\nu) = 2.74 \pm 0.27(\text{stat}) \pm 0.42(\text{syst}) \text{ pb}$$

(Neu, Thomson, Heinrich)

ALPGEN prediction: 0.78 pb

From our $W + 1b$ calculation:

$$\sigma_{b\text{-jet}}(W + b\text{jets}) \cdot Br(W \rightarrow l\nu) = 1.22 \pm 0.14 \text{ pb}$$

- consistency of theoretical results (good!);
- waiting for experimental distributions to compare to;
- important to have a separate analysis from $D0$!

Ongoing and future activity on $W/ZQ\bar{Q}$ production ...

- Provide input to experimental studies:
 - ▷ DØ, CDF: $W + b$ study, Higgs and single-top working groups;
 - ▷ CMS Higgs working group: provide parton level distributions with specific cuts and interface with NLO parton shower Monte Carlo (POWHEG):
 - study $Z + 2b$ as background to $H \rightarrow ZZ$ inclusive production;
 - study $Z + 2b$ and $Z + 1b$ to measure b -PDF.
 - ▷ ATLAS study of WH with $H \rightarrow b\bar{b}$ in boosted regime (Butterworth, Piacquadio, et al.).
- $Z + 1b$ -jet using both 4FNS with $m_b \neq 0$ and 5FNS NLO calculations:
 - ▷ $bg \rightarrow Zb$ at tree level and one loop (with $m_b = 0$);
 - ▷ $bg \rightarrow Zb + g, bq \rightarrow Zb + q$ (with $m_b = 0$);
 - ▷ $q\bar{q}, gg \rightarrow Zb\bar{b}$ at tree level and one loop (with $m_b \neq 0$);
 - ▷ $q\bar{q}, gg \rightarrow Zb\bar{b} + g$ and $gq(g\bar{q}) \rightarrow Zb\bar{b} + q(\bar{q})$ (with $m_b \neq 0$).

quite different pattern: one loop corrections to $q\bar{q}, gg \rightarrow b\bar{b}Z$ are now a piece of the NNLO 5FNS calculation, comparable to two-loop corrections to (and double parton emission from) $bg \rightarrow Zb$ (?)

- Possible to use $Z + 1b$ to measure b -PDF? \rightarrow reduce the PDF error in $H + 1b$ production?
- Resum large final state collinear logs from $g \rightarrow b\bar{b}$ splitting in $Wg \rightarrow Wb\bar{b}$ (and $Zg \rightarrow Zb\bar{b}$).
- Need too investigate NLO corrections to $qg(\bar{q}g) \rightarrow Wb\bar{b} + q'(\bar{q}')$?
(now within reach)