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Typos, Errata and Additions Essential Graduate Electrodynamics Lecture Notes Bernd Berg, Fall 2016

p.11: After equation (1.38):

By reasons that will become clear in the later physical interpretation the variable ζ is called *rapidity*. For each value of d there is precisely one value of ζ and vice versa, because the hyperbolic sine performs a one to one mapping. Such a function is called *isomorphic* in calculus.

Using finally
$$\binom{x^0}{x^1} = \binom{1}{1}$$
 yields
$$0 = [\cosh(\zeta) - \sinh(\eta)]^2 - [-\sinh(\zeta) + \cosh(\eta)]^2$$

$$= -2 \cosh(\zeta) \sinh(\eta) + 2 \sinh(\zeta) \cosh(\eta)$$

$$= -2 \tanh(\eta) + 2 \tanh(\zeta) \implies \zeta = \eta. \tag{1.41}$$

p.28:

Consequently, we have

$$\lambda' = \lambda \sqrt{\frac{1+\beta}{1-\beta}} = \lambda \sqrt{\frac{\cosh \zeta + \sinh \zeta}{\cosh \zeta - \sinh \zeta}} = \lambda \exp(\zeta). \tag{1.130}$$

p.70:

$$Y_l^{-l} = \sqrt{\frac{(2l+1)(2l)!}{4\pi}} \frac{\sin^l \theta}{2^l l!} e^{-il\phi}. \tag{2.154}$$

p.77:

The traceless quadrupole tensor is defined by

$$Q^{ij} = Q_S^{ij} - Q_D^{ij} \text{ with } Q_S^{ij} = 3 \int x'^i x'^j \rho(\vec{x}') d^3 x', \ Q_D^{ij} = \int r'^2 \delta^{ij} \rho(\vec{x}') d^3 x'$$
 so that $\sum_{i=1}^3 Q_S^{ii} = \sum_{i=1}^3 Q_D^{ii}$ holds. (3.10)

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p.80:

$$\Phi(\vec{x}) = \frac{\rho(\vec{x}')}{|\vec{x} - \vec{x}'|} + \frac{\vec{p}(\vec{x}') \cdot (\vec{x} - \vec{x}')}{|\vec{x} - \vec{x}'|^3}.$$
(3.22)

p.101:

Faraday's Law law of induction

$$\epsilon_{\text{emf}} = \oint_C \left(\vec{E} + \vec{\beta} \times \vec{B} \right) \cdot d\vec{l} = -\frac{1}{c} \frac{d}{dt} \Phi_m \text{ with } \Phi_m = \int_S \vec{B} \cdot d\vec{a}, \qquad (4.4)$$

p.102:

$$\nabla \cdot \vec{B} = 0, \tag{4.7}$$

which implies due to Gauss' theorem that the flux through a closed surface is zero.

p.183:

$$v_g = \frac{d\omega}{dk_q} = \frac{c}{\sqrt{\mu\epsilon}} \frac{d}{dk_q} \sqrt{k_q^2 + \gamma_q^2} = \frac{c}{\sqrt{\mu\epsilon}} \frac{k_q}{\sqrt{k_q^2 + \gamma_q^2}} < \frac{c}{\sqrt{\mu\epsilon}}$$
 (6.48)

p.190:

$$\vec{E}_t = \frac{1}{2} \left(\vec{E}_t^+ + \vec{E}_t^- \right) = -\sin(kz) \frac{k}{\gamma^2} \nabla_t \psi(x, y) \quad (6.70)$$

p.212:

$$\vec{E}_{\rm in} = \hat{\epsilon}_{\rm in} E_{\rm in} e^{ik_{\rm in}\hat{k}_{\rm in} \cdot \vec{x}} , \qquad \vec{B}_{\rm in} = \hat{k}_{\rm in} \times \vec{E}_{\rm in} \qquad (7.87)$$

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p.242:

The equation

$$\begin{split} 0 &= L_- \, Y_l^{\ -l} = \, L_- \, P_l^{\ -l} \, e^{-i \, l \, \phi} \\ &= e^{-i \, \phi} \, \left(- \frac{\partial}{\partial \theta} + i \, \cot \theta \, \frac{\partial}{\partial \phi} \right) \, P_l^{\ -l} \, e^{-i \, l \, \phi} \ , \end{split}$$

p.256:

(3) Show that

$$\begin{split} \epsilon_{\alpha\beta_1\gamma_1\delta_1}\epsilon^{\alpha\beta_2\gamma_2\delta_2} \; &= -\delta_{\beta_1}^{\;\;\beta_2}\delta_{\gamma_1}^{\;\;\gamma_2}\delta_{\delta_1}^{\;\;\delta_2} - \delta_{\beta_1}^{\;\;\gamma_2}\delta_{\gamma_1}^{\;\;\delta_2}\delta_{\delta_1}^{\;\;\beta_2} - \delta_{\beta_1}^{\;\;\delta_2}\delta_{\gamma_1}^{\;\;\beta_2}\delta_{\delta_1}^{\;\;\gamma_2} \\ &+ \delta_{\beta_1}^{\;\;\beta_2}\delta_{\gamma_1}^{\;\;\delta_2}\delta_{\delta_1}^{\;\;\gamma_2} + \delta_{\beta_1}^{\;\;\delta_2}\delta_{\gamma_1}^{\;\;\gamma_2}\delta_{\delta_1}^{\;\;\beta_2} + \delta_{\beta_1}^{\;\;\gamma_2}\delta_{\gamma_1}^{\;\;\beta_2}\delta_{\delta_1}^{\;\;\delta_2} \,. \end{split}$$

p.285:

(4) Use the limit of figure E.8 to calculate the principal value integral defined by

$$\frac{1}{2\pi i} P \int_{-\infty}^{\infty} dx \, \frac{f(x)}{x - \omega_R} = \frac{1}{2\pi i} \lim_{\eta \to 0^+} \left[\int_{-\infty}^{\omega_R - \eta} \frac{f(x) \, dx}{x - \omega_R} + \int_{\omega_R + \eta}^{+\infty} \frac{f(x) \, dx}{x - \omega_R} \right]$$

$$= \frac{1}{2\pi i} \lim_{\eta \to 0^+} \left[\int_{-\infty}^{\omega_R + i\eta - \eta} \frac{f(z) \, dz}{z - \omega_R - i\eta} + \int_{\omega_R + i\eta + \eta}^{+\infty} \frac{f(z) \, dz}{z - \omega_R - i\eta} \right]$$