

Special and General Relativity (PHZ 4601/5606) Fall 2018 Solutions

Set 12

29. Electromagnetic current conservation. Exercise 14.2 of Rindler.

We start off with Rindler (14.2), p.298,

$$E_{\mu\nu} = \Phi_{\nu;\mu} - \Phi_{\mu;\nu} ,$$

where $:$ denotes the covariant derivative. The field equation may be written as

$$E_{\mu\nu}{}^{;\nu} = g^{\nu\hat{\nu}} E_{\mu\nu;\hat{\nu}} = g^{\nu\hat{\nu}} \Phi_{\nu;\mu\hat{\nu}} - g^{\nu\hat{\nu}} \Phi_{\mu;\nu\hat{\nu}} = \frac{4\pi}{c} J_{\mu} .$$

We want to show that the continuity equation

$$\begin{aligned} g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} E_{\mu\nu;\hat{\nu}\hat{\mu}} &= g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} \Phi_{\nu;\mu\hat{\nu}\hat{\mu}} - g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} \Phi_{\mu;\nu\hat{\nu}\hat{\mu}} \\ &= \frac{4\pi}{c} g^{\mu\hat{\mu}} J_{\mu;\hat{\mu}} = \frac{4\pi}{c} J_{\mu}{}^{;\mu} = 0 . \end{aligned}$$

Here we have used that covariant differentiation commutes with contraction. See Rindler (10.34), p.212. Let us interchange the μ and ν dummies in the second term:

$$g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} \Phi_{\mu;\nu\hat{\nu}\hat{\mu}} = g^{\mu\hat{\mu}} g^{\nu\hat{\nu}} \Phi_{\nu;\mu\hat{\nu}\hat{\mu}}$$

After this we use Rindler (10.57), p.218 and get

$$\begin{aligned} \frac{4\pi}{c} J_{\mu}{}^{;\mu} &= g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} (\Phi_{\nu;\mu\hat{\nu}\hat{\mu}} - \Phi_{\nu;\mu\hat{\mu}\hat{\nu}}) = g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} (\Phi_{\alpha;\mu} R^{\alpha}{}_{\nu\hat{\nu}\hat{\mu}} + \Phi_{\nu;\alpha} R^{\alpha}{}_{\mu\hat{\nu}\hat{\mu}}) \\ &= g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} (\Phi_{\alpha;\mu} R_{\hat{\alpha}\nu\hat{\nu}\hat{\mu}} + \Phi_{\nu;\alpha} R_{\hat{\alpha}\mu\hat{\nu}\hat{\mu}}) . \end{aligned}$$

Re-labelling the dummies in the second term by $\alpha \rightarrow \beta$, $\nu \rightarrow \alpha$ and $\mu \rightarrow \nu$ gives

$$\frac{4\pi}{c} J_{\mu}{}^{;\mu} = g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} \Phi_{\alpha;\mu} R_{\hat{\alpha}\nu\hat{\nu}\hat{\mu}} + g^{\alpha\hat{\alpha}} g^{\nu\hat{\nu}} g^{\beta\hat{\beta}} \Phi_{\alpha;\beta} R_{\hat{\beta}\nu\hat{\alpha}\hat{\nu}} .$$

Next, replacing in the second term $\beta \rightarrow \mu$ gives

$$\frac{4\pi}{c} J_{\mu}{}^{;\mu} = g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} \Phi_{\alpha;\mu} (R_{\hat{\alpha}\nu\hat{\nu}\hat{\mu}} + R_{\hat{\mu}\nu\hat{\alpha}\hat{\nu}}) .$$

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Using, respectively, identities (10.64), (10.62) of p.218 of Rindler, and relabelling we find

$$\begin{aligned} g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} R_{\hat{\mu}\nu\hat{\alpha}\hat{\nu}} &= g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} R_{\hat{\alpha}\hat{\nu}\hat{\mu}\nu} = \\ -g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} R_{\hat{\alpha}\hat{\nu}\nu\hat{\mu}} &= -g^{\nu\hat{\nu}} g^{\mu\hat{\mu}} g^{\alpha\hat{\alpha}} R_{\hat{\alpha}\nu\hat{\nu}\hat{\mu}} \end{aligned}$$

implying

$$\frac{4\pi}{c} J_{\mu;\mu} = 0.$$