Questions

Atwood machine: With

$$m_1 = 251 g = 0.251 kg$$
 and $m_2 = 250 g = 0.250 kg$

we measured we measured

$$\triangle t = 12 s$$
 for $\triangle x = 1 m$.

Approximately, this corresponds to the following lower bound on the gravitational constant g (pick one):

- 1. $5 m/s^2$
- 2. $6 m/s^2$
- 3. $7 m/s^2$
- 4. $8 m/s^2$
- 5. $9 m/s^2$
- 6. $10 m/s^2$

Hint: Calculate the accelerataion a from the our measurements of $\triangle t$ and $\triangle x$ and solve

$$(m_1 - m_2) g = (m_1 + m_2) a$$

for a.



Why is the estimate a lower bound?

- 1. Because of friction.
- 2. Because of the mass of the pulley.
- 3. Because of friction and the mass of the pulley.
- 4. Because of the tension in the rope.
- 5. Because of friction, the mass of the pulley and the tension in the rope.



Coefficient of static friction:

We measured

$$\theta_{\rm max} = 30^0$$

for a block of wood on an inclined plane of wood. This correspond to the coefficient of static friction (use your calculator and pick the result)

 $\mu_s =$

- 1. 0.55
- 2. 0.55 m
- 3. 0.58
- 4. 0.58 m
- 5. 0.61
- 6. 0.61 m

Circular Motion (Chapter 5-2 of Tipler)

Centripetal Acceleration:

Pythagorean theorem for figure 5-23 of Tipler:

$$(r+h)^{2} = r^{2} + (v t)^{2}$$

$$r^{2} + 2 h r + h^{2} = r^{2} + (v t)^{2}$$

$$2 h r + h^{2} = v^{2} t^{2}$$

Limit $h \to 0$ (neglect $O(h^2)$):

$$2 h r = v^{2} t^{2}$$

$$h = \frac{1}{2} \frac{v^{2}}{r} t^{2} = \frac{1}{2} a t^{2}$$

$$a = \frac{v^{2}}{r}$$



Position and Velocity Vectors: Tipler figure 5-24.

The angle $\Delta\theta$ between \vec{v}_1 and \vec{v}_2 is the same as that between \vec{r}_1 and \vec{r}_2 , because the position and velocity vectors must remain mutually perpendicular. The magnitude of the acceleration can be found from the following relations, which hold for in the limit $\Delta\theta \to 0$, i.e. for very small angles $\Delta\theta$.

$$\Delta \theta = \frac{\Delta r}{r} = \frac{\Delta v}{v}$$

$$\triangle v = \triangle \theta \ v = \triangle r \frac{v}{r}$$

$$\Delta r = v \, \Delta t$$
$$\Delta t = \frac{\Delta r}{v}$$

Therefore.

$$a = \lim_{\Delta t \to 0} \frac{\Delta v}{\Delta t} = \lim_{\Delta t \to 0} \frac{\Delta r}{\Delta r} \frac{v^2}{r} = \frac{v^2}{r}$$

$$a = \frac{v^2}{r}$$



The Period T:

The time T required for one complete revolution is called the period. For constant speed

$$v = \frac{2\pi r}{T}$$
 holds.

Components of acceleration for a particle moving along an arbitrary curve with varying speed: Figure 5-25 of Tipler.

We can treat a portion of the curve as an arc of a circle.

The has then a component of acceleration tangent to the circle

$$\frac{dv}{dt}$$

as well as a radially inward centripetal acceleration

$$\frac{v^2}{r}$$



Centripetal Force:

As with any acceleration, there must be a force in the direction of the acceleration. For centripetal accelerations it is called the centripetal force

$$ec{F}_{cp} = -m \, rac{v^2}{r} \, \hat{r}$$

where \hat{r} is the unit vector in the direction of \vec{r} :

$$\hat{r} = \frac{\vec{r}}{r}$$

Centrifugal Force:

This is the force opposite to the centripetal force, which acts on the entity, which pulls the object towards the center of the circle.

$$ec{F}_{cf} = + m \, rac{v^2}{r} \, \hat{r} = - ec{F}_{cp}$$



Example: bucket of water in vertical, circular motion.

The force of the water onto the bottom of the bucket is:

$$F_{top} = m \frac{v^2}{r} - m g$$

at the top of the circle, and

$$F_{bot} = m \, \frac{v^2}{r} + m \, g$$

at the bottom of the circle.

How fast must velocity be, such that the water does not spill?

$$\begin{array}{ll} 1. & v > \sqrt{g/r} \\ 2. & v < \sqrt{g/r} \end{array}$$

$$2. \ v < \sqrt{g/r}$$

Example: Circular Pendulum: Figures 5-28 and 5-29 of Tipler.

$$\vec{T} + \vec{F}_{cf} + m \, \vec{g} = 0$$

$$\vec{T} = T_r \, \hat{r} + T_y \, \hat{y}$$

$$T_r = -T \, \sin(\theta) = -m \, \frac{v^2}{r}$$

$$T_y = T \, \cos(\theta) = m \, g$$

$$\frac{\sin(\theta)}{\cos(\theta)} = \frac{v^2}{g \, r}$$

$$\tan(\theta) = \frac{v^2}{g \, r}$$

$$v = \sqrt{g \, r \, \tan(\theta)}$$



Example: Forces on a car in a banked curve:

Figures 5-32 and 5-33 of Tipler. The optimal angle θ is the one for which the centrifugal force is balanced by the inward component of the normal force. Then:

$$F_n \cos(\theta) - m g = 0$$

$$F_n \sin(\theta) - m \frac{v^2}{r} = 0$$

$$\tan(\theta) = \frac{v^2}{g r}$$

