

SOLUTIONS FINAL ADVANCED DYNAMICS

PHY 4241 (Spring, 2010)

PROBLEM 1

See Homework 6 and $L(q_i + \epsilon_i, \dot{q}_i, t) = L(q_i, \dot{q}_i, t)$ implies

$$0 = \frac{\partial L}{\partial q_i} = \frac{d}{dt} \frac{\partial L}{\partial \dot{q}_i} \Rightarrow p_i = \frac{\partial L}{\partial \dot{q}_i} \text{ conserved (generalized momentum).}$$

PROBLEM 2

The electromagnetic field tensor is given by Find (the Einstein summation convention is assumed for identical indices):

1. $F^\alpha{}_\alpha = 0$.

2.

$$(F_{\alpha\beta}) = \begin{pmatrix} 0 & +E^x & +E^y & +E^z \\ -E^x & 0 & -B^z & B^y \\ -E^y & B^z & 0 & -B^x \\ -E^z & -B^y & B^x & 0 \end{pmatrix}.$$

3. $F_{\alpha\beta}F^{\alpha\beta} = 2(\vec{B}^2 - \vec{E}^2)$.

PROBLEM 3

1. Yes, because L does not depend explicitly on time.

2.

$$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}.$$

The center of mass frame is defined by $\vec{R} = 0$.

3.

$$\begin{aligned} \vec{r}_1 - \vec{r}_2 + \left(\frac{m_1 + m_2}{m_2} \right) \vec{R} &= +\vec{r}_1 + \frac{m_1}{m_2} \vec{r}_1 = \frac{m_1 + m_2}{m_2} \vec{r}_1, \\ \vec{r}_1 - \vec{r}_2 - \left(\frac{m_1 + m_2}{m_1} \right) \vec{R} &= -\vec{r}_2 - \frac{m_2}{m_1} \vec{r}_2 = \frac{m_1 + m_2}{m_1} \vec{r}_2, \end{aligned}$$

Therefore, in the center of mass frame

$$\vec{r}_1 = \frac{m_2}{m_1 + m_2} \vec{r} \quad \text{and} \quad \vec{r}_2 = -\frac{m_1}{m_1 + m_2} \vec{r}.$$

The Lagrangian becomes

$$L = \mu \dot{\vec{r}}^2 - \frac{\alpha}{|\vec{r}|} \quad \text{with} \quad \mu = \frac{m_1 m_2}{m_1 + m_2} \text{ reduced mass.}$$

4. In spherical coordinates we have

$$L = \frac{1}{2} \mu \dot{r}^2 + \frac{1}{2} \mu r^2 \left(\dot{\theta}^2 + \sin^2(\theta) \dot{\phi}^2 \right) - \frac{\alpha}{|\vec{r}|}.$$

5. The Euler-Lagrange equations for the angles give

$$0 = \frac{d}{dt} \frac{\partial L}{\partial \dot{\theta}} - \frac{\partial L}{\partial \theta} = 2\mu r \dot{r} \dot{\theta} + \mu r^2 \left(\ddot{\theta} - \sin(\theta) \cos(\theta) \dot{\theta} \dot{\phi}^2 \right),$$

$$0 = \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} - \frac{\partial L}{\partial \phi} = \frac{d}{dt} \left(\frac{1}{2} \mu r^2 \sin^2(\theta) \dot{\phi} \right).$$

From the last equation we find that angular momentum is conserved:

$$L_\phi = \mu r^2 \sin^2(\theta) \dot{\phi} = \text{constant}.$$

6. Let $\dot{\phi}(0) = 0$ and $\theta \neq 0, \pi$. Then for *all* times

$$\sin^2[\theta(t)] \dot{\phi}(t) = 0 \Rightarrow \dot{\phi}(t) = 0 \Rightarrow \phi(t) = \text{constant}.$$

The motion is in a plane.

7. Yes, just choose $\hat{\phi}(0)$ orthogonal to $\vec{r}(0) \times \dot{\vec{r}}(0)$.