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## Electrodynamics A (PHY 5346) Fall 2016 Solutions

## Set 3:

(7) Exercise (CW 5 points): Lorentz group Lie matrix generator. The Matrix L is:

$$L = \begin{pmatrix} l_0^0 & l_1^0 & l_2^0 & l_3^0 \\ l_0^1 & l_1^1 & l_2^1 & l_3^1 \\ l_0^2 & l_1^2 & l_2^2 & l_3^2 \\ l_0^3 & l_1^3 & l_2^3 & l_3^3 \end{pmatrix}$$

(1) The matrix g is:

$$g = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}.$$

By matrix multiplication:

$$-gL = \begin{pmatrix} -l_0^0 & -l_1^0 & -l_2^0 & -l_3^0 \\ l_0^1 & l_1^1 & l_2^1 & l_3^1 \\ l_0^2 & l_1^2 & l_2^2 & l_3^2 \\ l_0^3 & l_1^3 & l_2^3 & l_3^3 \end{pmatrix}.$$

(2) The transpose of L is:

$$\widetilde{L} = \begin{pmatrix} \widetilde{l}_{0}^{0} \ \widetilde{l}_{0}^{1} \ \widetilde{l}_{0}^{2} \ \widetilde{l}_{0}^{3} \\ \widetilde{l}_{1}^{0} \ \widetilde{l}_{1}^{1} \ \widetilde{l}_{1}^{2} \ \widetilde{l}_{1}^{3} \\ \widetilde{l}_{2}^{0} \ \widetilde{l}_{2}^{1} \ \widetilde{l}_{2}^{2} \ \widetilde{l}_{2}^{3} \\ \widetilde{l}_{3}^{0} \ \widetilde{l}_{3}^{1} \ \widetilde{l}_{3}^{2} \ \widetilde{l}_{3}^{3} \end{pmatrix} = \begin{pmatrix} l_{0}^{0} \ l_{0}^{1} \ l_{0}^{2} \ l_{0}^{3} \\ l_{0}^{1} \ l_{1}^{1} \ l_{1}^{2} \ l_{3}^{3} \\ l_{0}^{0} \ l_{1}^{1} \ l_{2}^{2} \ l_{3}^{3} \\ l_{0}^{0} \ l_{1}^{1} \ l_{3}^{2} \ l_{3}^{3} \end{pmatrix}.$$

(3) By matrix multiplication:

$$\widetilde{L}g = \begin{pmatrix} l_0^0 - l_1^0 - l_0^2 - l_0^3 \\ l_1^0 - l_1^1 - l_1^2 - l_1^3 \\ l_2^0 - l_2^1 - l_2^2 - l_2^3 \\ l_3^0 - l_3^1 - l_3^2 - l_3^3 \end{pmatrix}.$$

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(4) Set  $-gL = \widetilde{L}g$  and compare the components:

$$\begin{pmatrix} -l_0^0 = l_0^0 - l_1^0 = -l_1^1 - l_2^0 = -l_2^2 - l_3^0 = -l_3^0 \\ l_0^1 = l_1^0 - l_1^1 = -l_1^1 - l_2^1 = -l_2^1 - l_3^1 = -l_3^1 \\ l_0^2 = l_2^0 - l_1^2 = -l_2^1 - l_2^2 = -l_2^2 - l_3^2 = -l_3^2 \\ l_0^3 = l_3^0 - l_3^1 = -l_3^1 - l_3^2 = -l_3^2 - l_3^2 = -l_3^2 \end{pmatrix}$$

The resulting matrix is:

$$\begin{pmatrix} 0 & l_1^0 & l_2^0 & l_3^0 \\ l_1^0 & 0 & l_2^1 & l_3^1 \\ l_2^0 - l_2^1 & 0 & l_3^2 \\ l_3^0 - l_3^1 - l_3^2 & 0 \end{pmatrix}$$

(8) Exercise (5 points): Obtain the result of the previous Exercise discussing the elements of  $g^{\alpha\beta} \tilde{l}^{\ \gamma}_{\beta} g_{\gamma\delta} = -l^{\alpha}_{\delta}$ . Performing the contractions we get  $\tilde{l}^{\alpha}_{\ \delta} = -l^{\alpha}_{\ \delta}$ . Using now the definition  $\tilde{l}^{\alpha}_{\ \delta} = l^{\ \alpha}_{\ \delta}$ , the following equations are obtained (no summations, a = 0, 1, 2, 3, i = 1, 2, 3 and j = 1, 2, 3):

$$\begin{split} \tilde{l}^{a}_{\ a} &= l^{\ a}_{a} = + l^{a}_{\ a} \Rightarrow l^{a}_{\ a} = - l^{a}_{\ a} = 0 \ , \\ \\ \tilde{l}^{i}_{\ 0} &= l^{\ i}_{0} = - l^{0}_{\ i} \Rightarrow l^{0}_{\ i} = + l^{i}_{\ 0} \ , \\ \\ \\ \tilde{l}^{i}_{\ j} &= l^{\ i}_{j} = + l^{j}_{\ i} \Rightarrow l^{j}_{\ i} = - l^{i}_{\ j} \end{split}$$

and the resulting matrix is the same as before.

(9) Exercise: A rotation and a boost generator.

For  $S_3$  we have

$$S_3 = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \quad \text{and} \quad (S_3)^2 = -\mathbf{1}_{2,3} \quad \text{with} \quad \mathbf{1}_{2,3} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} .$$

Therefore, for n = 1, 2, 3, ...

$$(S_3)^{2n} = (-1)^n \mathbf{1}_{2,3}$$
 and  $(S_3)^{2n+1} = (-1)^n S_3$ 

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and the result follows:

$$e^{-\omega S_3} = \mathbf{1} + \sum_{n=1}^{\infty} (-1)^n \frac{\omega^{2n}}{(2n)!} \mathbf{1}_{2,3} - \sum_{n=0}^{\infty} (-1)^n \frac{\omega^{2n+1}}{(2n+1)!} S_3$$
$$= \begin{pmatrix} 1 & 0 & 0 & 0\\ 0 & \cos \omega & \sin \omega & 0\\ 0 & -\sin \omega & \cos \omega & 0\\ 0 & 0 & 0 & 1 \end{pmatrix}.$$

Similarly we deal with  $K_1$ :

Therefore, for n = 1, 2, 3, ...

$$(K_1)^{2n} = \mathbf{1}_{1,2}$$
 and  $(K_1)^{2n+1} = K_1$ 

and

$$e^{-\zeta K_1} = \mathbf{1} + \sum_{n=1}^{\infty} \frac{\omega^{2n}}{(2n)!} \mathbf{1}_{1,2} - \sum_{n=0}^{\infty} \frac{\omega^{2n+1}}{(2n+1)!} K_1 = \begin{pmatrix} \cosh \zeta & -\sinh \zeta & 0 & 0 \\ -\sinh \zeta & \cosh \zeta & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

## (10) Exercise: Four-dimensional Levi-Civita tensor.

(1) Using the Einstein summation convention of the left-hand side, but not on the right-hand side, of the following equation we have

$$g_{\alpha_1\alpha_2}g_{\beta_1\beta_2}g_{\gamma_1\gamma_2}g_{\delta_1\delta_2}\epsilon^{\alpha_2\beta_2\gamma_2\delta_2} = g_{\alpha_1\alpha_1}g_{\beta_1\beta_1}g_{\gamma_1\gamma_1}g_{\delta_1\delta_1}\epsilon^{\alpha_1\beta_1\gamma_1\delta_1}.$$

The right-hand side is zero unless  $\alpha_1\beta_1\gamma_1\delta_1$  are a permutation  $\pi_0\pi_1\pi_2\pi_3$ of the numbers 0, 1, 2, 3 and in that case

$$\epsilon_{\pi_0\pi_1\pi_2\pi_3} = g_{\pi_0\pi_0}g_{\pi_1\pi_1}g_{\pi_1\pi_2}g_{\pi_1\pi_3}\epsilon^{\pi_0\pi_1\pi_2\pi_3} = (-1)^3 \epsilon^{\pi_0\pi_1\pi_2\pi_3},$$

which proves  $\epsilon_{\alpha\beta\gamma\delta} = -\epsilon^{\alpha\beta\gamma\delta}$ .

(2) We have  $\epsilon_{\alpha\beta\gamma_1\delta_1}\epsilon^{\alpha\beta\gamma_2\delta_2} = 0$  unless either  $\gamma_1 = \gamma_2$ ,  $\delta_1 = \delta_2$  or  $\gamma_1 = \delta_2, \, \delta_1 = \gamma_2$  or (otherwise one of the already used numbers will be repeated by  $\alpha$  or  $\beta$  of the sums). Therefore,

$$\epsilon_{\alpha\beta\gamma_1\delta_1}\epsilon^{\alpha\beta\gamma_2\delta_2} = a\,\delta_{\gamma_1}^{\ \gamma_2}\,\delta_{\delta_1}^{\ \delta_2} + b\,\delta_{\gamma_1}^{\ \delta_2}\,\delta_{\delta_1}^{\ \gamma_2}$$

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holds. With no summation in the permutation indices  $\pi_0$  to  $\pi_3$  the constants follow from

$$\epsilon_{\alpha\beta\pi_{2}\pi_{3}}\epsilon^{\alpha\beta\pi_{2}\pi_{3}} = \epsilon_{\pi_{0}\pi_{1}\pi_{2}\pi_{3}}\epsilon^{\pi_{0}\pi_{1}\pi_{2}\pi_{3}} + \epsilon_{\pi_{1}\pi_{0}\pi_{2}\pi_{3}}\epsilon^{\pi_{1}\pi_{0}\pi_{2}\pi_{3}} = -2 = a ,$$
  

$$\epsilon_{\alpha\beta\pi_{2}\pi_{3}}\epsilon^{\alpha\beta\pi_{3}\pi_{2}} = \epsilon_{\pi_{0}\pi_{1}\pi_{2}\pi_{3}}\epsilon^{\pi_{0}\pi_{1}\pi_{3}\pi_{2}} + \epsilon_{\pi_{1}\pi_{0}\pi_{2}\pi_{3}}\epsilon^{\pi_{1}\pi_{0}\pi_{3}\pi_{2}} = +2 = b .$$

(3) We have  $\epsilon_{\alpha\beta_1\gamma_1\delta_1}\epsilon^{\alpha\beta_2\gamma_2\delta_2} = 0$  unless  $\beta_2\gamma_2\delta_2$  is a permutation of  $\beta_1\gamma_1\delta_1$ . There are six such permutations, so that the results is a sum of the form

$$\begin{split} \epsilon_{\alpha\beta_{1}\gamma_{1}\delta_{1}} \epsilon^{\alpha\beta_{2}\gamma_{2}\delta_{2}} &= a_{1}\,\delta_{\beta_{1}}^{\ \beta_{2}}\delta_{\gamma_{1}}^{\ \gamma_{2}}\delta_{\delta_{1}}^{\ \delta_{2}} + a_{2}\,\delta_{\beta_{1}}^{\ \gamma_{2}}\delta_{\gamma_{1}}^{\ \delta_{2}}\delta_{\delta_{1}}^{\ \beta_{2}} + a_{3}\,\delta_{\beta_{1}}^{\ \delta_{2}}\delta_{\gamma_{1}}^{\ \beta_{2}}\delta_{\delta_{1}}^{\ \gamma_{2}} \\ &+ b_{1}\,\delta_{\beta_{1}}^{\ \beta_{2}}\delta_{\gamma_{1}}^{\ \delta_{2}}\delta_{\delta_{1}}^{\ \gamma_{2}} + b_{2}\,\delta_{\beta_{1}}^{\ \delta_{2}}\delta_{\gamma_{1}}^{\ \gamma_{2}}\delta_{\delta_{1}}^{\ \beta_{2}} + b_{3}\,\delta_{\beta_{1}}^{\ \gamma_{2}}\delta_{\gamma_{1}}^{\ \beta_{2}}\delta_{\delta_{1}}^{\ \delta_{2}} \,, \end{split}$$

where it follows from (no summation)  $\epsilon_{\pi_0\pi_1\pi_2\pi_3}\epsilon^{\pi_0\pi_1\pi_2\pi_3} = -1$  that  $a_i = -1$  and  $b_i = 1$  holds for i = 1, 2, 3.

(11) Exercise: Addition theorem for transverse velocity components.

In K' the motion with velocity  $\vec{u'}$  is

$$x^{\prime i} \; = \; c^{-1} \, u^{\prime i} \, x^{\prime 0} \, .$$

With respect to frame K the origin of K' frame moves with speed v along the  $x^1$  axis of K. For i = 2, 3 the Lorentz transformations give then

$$x^i \; = \; x'^i \; = \; c^{-1} \, u'^i \, \gamma \, (x^0 - \beta \, x^1) \, .$$

Dividing by  $x^0$  gives the velocity components in K:

$$c^{-1} u^{i} = c^{-1} u^{\prime i} \gamma \left( 1 - \beta c^{-1} u^{1} \right) = c^{-1} u^{\prime i} \gamma \left( 1 - \beta c^{-1} \frac{u^{\prime 1} + v}{1 + u^{\prime 1} v/c^{2}} \right),$$

where in the second step the already calculated equation for  $u^1$  has been inserted. Bringing everything to the common denominator gives

$$\begin{aligned} u^{i} &= u'^{i} \gamma \, \frac{(1 + \beta \, u'^{1}/c) - \beta \, c^{-1} \, (u'^{1} + v)}{1 + u'^{1} v/c^{2}} \\ &= u'^{i} \gamma \, \frac{1 - \beta^{2}}{1 + u'^{1} v/c^{2}} \, = \, \frac{u'^{i}}{\gamma \left(1 + u'^{1} v/c^{2}\right)} \,, \end{aligned}$$

where  $1 - \beta^2 = 1/\gamma^2$  was used in the last step.