

Electrodynamics A (PHY 5346) Fall 2016 Solutions

Set 9:

(27) Exercise E.41: Potential over a plane.

The Dirichlet Green function is obtained with the method of images

$$G_D(\vec{r}, \vec{r}') = -\frac{1}{4\pi\sqrt{(\vec{\rho} - \vec{\rho}')^2 + (z - z')^2}} + \frac{1}{4\pi\sqrt{(\vec{\rho} - \vec{\rho}')^2 + (z + z')^2}}.$$

and

$$\left. \frac{\partial G_D}{\partial z'} \right|_{z'=0} = -\frac{z}{2\pi [(\vec{\rho} - \vec{\rho}')^2 + z^2]^{3/2}}.$$

The potential on the z axis is then

$$\Phi(\vec{r}) = -\int_{-\infty}^{+\infty} dx' \int_{-\infty}^{+\infty} dy' \Phi(x', y', 0) \left. \frac{\partial G_D}{\partial z'} \right|_{z'=0}$$

where the potential $\Phi(x', y', 0)$ in the x - y -plane is given. In our case the integration is only over the upper half-plane as $\Phi = 0$ on the lower half-plane. On the z -axis $\rho = 0$ holds and

$$\begin{aligned} \Phi(z) &= \int_{-\infty}^{+\infty} dx' \int_0^{\infty} dy' \frac{z \Phi_1}{2\pi [(\vec{\rho}')^2 + z^2]^{3/2}} \\ &= \frac{z \Phi_1}{2\pi} \int_0^{\pi} d\phi' \int_0^{\infty} \rho' d\rho' \frac{1}{(\rho'^2 + z^2)^{3/2}} = \frac{\Phi_1}{2}. \end{aligned}$$

(28) Exercise E.44: Angled plates.

(a) The potential satisfies $\nabla^2 \Phi = 0$. Choosing cylindrical coordinates and recognizing that by symmetry reasons Φ does not depend on ρ and z , we have

$$\nabla^2 \Phi = \frac{1}{\rho^2} \frac{d^2 \Phi}{d\phi^2} = 0.$$

(b) The boundary conditions are $\Phi(0) = 0$ and $\Phi(\beta) = \Phi_0$ and the solution is

$$\Phi(\phi) = a\phi + b = \Phi_0 \frac{\phi}{\beta}.$$

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(c) The electric field is

$$\vec{E} = -\nabla\Phi = -\frac{\hat{\phi}}{\rho} \frac{d\Phi}{d\phi} = -\frac{\Phi_0}{\beta\rho} \hat{\phi}.$$

(d) Charge densities and total charge follow from Gauss' law. On each plate with opposite signs:

$$\int_S d\vec{a} \cdot \vec{E} = 4\pi Q, \quad -\hat{\phi} \cdot \vec{E} = 4\pi\sigma.$$

Therefore,

$$\frac{\Phi_0}{\beta\rho} \hat{\phi} = 4\pi\sigma \Rightarrow \sigma = \frac{\Phi_0}{4\pi\beta\rho}$$

and the total charge on each plate is

$$Q = \int_{\rho_1}^{\rho_2} d\rho \int_0^h dz \frac{\Phi_0}{4\pi\beta\rho} = \frac{\Phi_0 h}{4\pi\beta} \ln(\rho_2/\rho_1).$$

(d) The capacitance is

$$C = \frac{Q}{\Phi_0} = \frac{h}{4\pi\beta} \ln(\rho_2/\rho_1).$$

(e) The energy is

$$U = \frac{1}{2} C \Phi_0^2 = \frac{h \Phi_0^2}{8\pi\beta} \ln(\rho_2/\rho_1).$$

From this the torque follows to be

$$\vec{\tau} = \frac{dU}{d\beta} \hat{z} = -\frac{h \Phi_0^2}{8\pi\beta^2} \ln(\rho_2/\rho_1) \hat{z}.$$

(29) Exercise: Expansion in cylindrical coordinates.

As there is no azimuthal dependence in the BC, there will be not ϕ dependence in the result and the expansion reads

$$\Phi(\rho, z) = \sum_{m=0}^{\infty} \sum_{n=1}^{\infty} A_{mn} J_m(k_{mn}\rho) \sinh(k_{mn}z)$$

where due to the BC the only surviving coefficient is A_{01} . So,

$$\Phi(\rho, z) = A_{01} J_0(k_{01}\rho) \sinh(k_{01}z).$$

Now,

$$1 = \rho_0 = \frac{x_{01}}{k_{01}} \Rightarrow k_{01} = x_{01} \approx 2.405$$

and the BC at $z = 1$ is matched with $A_{01} = 1/\sinh(x_{01})$.

(a) The potential inside the cylinder is

$$\Phi(\rho, z) = J_0(x_{01}\rho) \frac{\sinh(x_{01}z)}{\sinh(x_{01})}.$$

(b) For $\rho = 0$, $z = 1/2$: $J_0(x_{01}\rho) = J_0(0) = 1$, $\sinh(x_{01}/2) = 1.514$.
Further, $\sinh(x_{01}) = 5.494$ and we find

$$\Phi\left(0, \frac{1}{2}\right) = J_0(0) \frac{1.514}{5.495} = 0.2756.$$

(c) For $\rho = 1/2$ we have $J_0(x_{01}/2) = 0.670$ and find

$$\Phi\left(\frac{1}{2}, \frac{1}{2}\right) = 0.670 \frac{1.514}{5.495} = 0.1846, .$$

(30) Exercise: Calculate L_- . See (C.17) to (C.23) lecture notes.