1. Electron-positron annihilation.

Let us use natural units with c = 1 and denote the four-vectors of the photons by p and q. Energy conservation:

$$p^{0} + q^{0} = 2E$$
 with $E = 0.511 [MeV]$.

Momentum conservation and massless photons:

$$\vec{p} = -\vec{q}$$
 and $p^2 = q^2 = 0$.

(a) Therefore, we have in the frame where the photon with momentum p moves along the positive x^1 direction

$$(p^{\alpha}) = \begin{pmatrix} E \\ E \\ 0 \\ 0 \end{pmatrix}$$
 and $(q^{\alpha}) = \begin{pmatrix} E \\ -E \\ 0 \\ 0 \end{pmatrix}$.

(b) $\beta = 3/5 \rightarrow \gamma = 1/\sqrt{1-\beta^2} = 5/4$ and the energy is the zero component of a four-vector transforms according to

$$E' = \gamma (1 - \beta) E = \frac{1}{2} E = 0.2555 [MeV].$$

2. Coaxial cable.

Ampére's law reads

$$\nabla \times \vec{B} = \frac{4\pi}{c} \vec{J}$$

where in our case $\vec{J} = J \hat{z}$ and

$$J = J(\rho) = \begin{cases} J_1 = I/(\pi \rho_1^2) & \text{for } \rho \le \rho_1, \\ 0 & \text{for } \rho_1 < \rho < \rho_2, \\ J_2 = -I/(\pi \rho_3^2 - \pi \rho_2^2) & \text{for } \rho_2 \le \rho \le \rho_3, \\ 0 & \text{for } \rho > \rho_3. \end{cases}$$

Now $\vec{B} = B_{\phi} \hat{\phi}$ with $B_{\phi} = B = |\vec{B}|$ and

$$2\pi \rho B_{\phi} = \rho \int_{0}^{2\pi} d\phi B_{\phi} = \oint d\vec{s} \cdot \vec{B} = \int_{S} da \,\hat{a} \cdot (\nabla \times \vec{B}) = \frac{4\pi}{c} \int_{S} da \,\hat{a} \cdot \vec{J}$$

$$\begin{split} \rho \leq \rho_1: \\ & 2\pi \,\rho \, B_\phi = \frac{4\pi}{c} \, J_1 \,\pi \,\rho^2 \,\Rightarrow\, B = \frac{2}{c} \, J_1 \,\pi \,\rho = \frac{2 \, I \,\rho}{c \,\rho_1^2} \,\,. \\ \rho_1 < \rho < \rho_2: \\ & 2\pi \,\rho \, B_\phi = \frac{4\pi}{c} \, J_1 \,\pi \,\rho_1^2 \,\Rightarrow\, B = \frac{2 \, I}{c \,\rho} \,\,. \\ \rho_2 \leq \rho \leq \rho_3: \\ & B = \frac{2 \, I}{c \,\rho} + \frac{2 \, J_2}{c \,\rho} \,\pi \left(\rho^2 - \rho_2^2\right) = \frac{2 \, I}{c \,\rho} \,\frac{\left(\rho_3^2 - \rho^2\right)}{\left(\rho_3^2 - \rho_2^2\right)} \,\,. \end{split}$$
Finally, $\rho > \rho_3: B = 0$.

3. Potential from distinct BCs on half-spheres.

(a) We have to calculate

$$\Phi(r) = R^2 \int_S d\Omega' \, \Phi_0 \left. \frac{\partial G_D}{\partial r'} \right|_{r'=R} \,,$$

where the surface is the upper half-sphere and the derivative of the Green function is given by

$$\left. \frac{\partial G_D}{\partial r'} \right|_{r'=R} = \frac{R^2 - r^2}{4\pi R \left(R^2 + r^2 - 2Rr \cos \gamma \right)^{3/2}}$$

with $\cos \gamma = \cos(\theta) \cos(\theta') + \sin(\theta) \sin(\theta') \cos(\phi - \phi')$. For r = z on the z-axis we have $\theta = 0$ and, hence, $\cos \gamma = \cos \theta'$. The integral over the upper half-sphere becomes

$$\Phi(z) = \Phi_0 \frac{R}{2} (R+z) (R-z) \int_0^1 \frac{d\cos\theta'}{(R^2 + z^2 - 2Rz\,\cos\theta')^{3/2}}.$$

With $a = R^2 + z^2$ and b = 2Rz the integral is

$$\int_0^1 \frac{dx}{(a-bx)^{3/2}} = \frac{2}{b(a-bx)^{1/2}} \Big|_0^1 = \frac{2}{b\sqrt{a-b}} - \frac{2}{b\sqrt{a}}$$
$$= \frac{1}{Rz\sqrt{(R-z)^2}} - \frac{1}{Rz\sqrt{R^2+z^2}}$$

and

$$\Phi(z) = \Phi_0 \frac{1}{2z} \left(R + z \right) \left(1 - \frac{R - z}{\sqrt{R^2 + z^2}} \right) \,.$$

(b) Special values:

$$z = -R \Rightarrow \Phi(-R) = 0,$$

$$z = -\frac{R}{2} \Rightarrow \Phi\left(-\frac{R}{2}\right) = \Phi_0\left(-\frac{1}{2}\right)\left(1 - \frac{3}{\sqrt{5}}\right) = 0.17082 \Phi_0$$

$$z = \frac{R}{2} \Rightarrow \Phi\left(\frac{R}{2}\right) = \Phi_0\left(\frac{3}{2}\right)\left(1 - \frac{1}{\sqrt{5}}\right) = 0.82918 \Phi_0$$

$$z = R \Rightarrow \Phi(R) = \Phi_0\frac{2R}{2R} = \Phi_0,$$

$$z \to 0 \Rightarrow \Phi(0) = \Phi_0\frac{R}{2z}\left(1 - \frac{R-z}{R}\right) \rightarrow \frac{1}{2}\Phi_0.$$

The plot is shown in the figure.

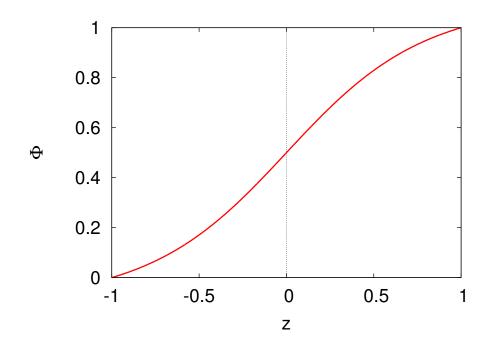


FIG. 1: Potential as function of z.