PHY 5667 : Quantum Field Theory A, Fall 2015

September 
$$10^{th}$$
, 2015  
Assignment # 3  
(due Thursday September  $24^{th}$ , 2015)

1. Consider a system of fields  $\phi_i(x)$  and a Lorentz transformation acting on both field and coordinates as follows,

$$x^{\mu} = \Lambda^{\mu}_{\nu} x^{\nu} \simeq x^{\mu} + \sum_{k=1}^{6} \alpha_k X^{\mu}_k + O(\alpha^2) ,$$
  
$$\phi'_i(x') = L_{ij}(\Lambda) \phi_i(x) \simeq \phi_i(x) + \sum_{k=1}^{6} \alpha_k A_{ij,k} \phi_i(x) + O(\alpha^2) ,$$

where  $L(\Lambda)$  denote the representation of the Lorentz group on the space of the fileds  $\phi_i(x)$ . Show that if the action S,

$$S = \int d^4x \mathcal{L}(\phi_i(x), \partial_\mu \phi_i(x), x) ,$$

associated to the system of fields is invariant under such transformation, there are six conserved currents (Noether's currents) of the form,

$$M^{\mu}_{\rho\sigma} = T^{\mu}_{\rho} x_{\sigma} - T^{\mu}_{\sigma} x_{\rho} - \frac{\partial \mathcal{L}}{\partial(\partial_{\mu}\phi_i(x))} A_{ij,\rho\sigma} \phi_j(x) ,$$

where  $T^{\mu\nu}$  is the energy-momentum tensor associated to the system of fields, and six conserved charges. How are the charges defined? How can you interpret the components of  $M^{\mu}_{\rho\sigma}$  due to the transformation of the coordinates and to the transformation of the fields respectively? Explain your reasoning.

2. Classical electromagnetism (with no sources) follows from the action,

$$S = \int d^4x \left( -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} \right) , \quad \text{where } F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

as you have derived in your first homework. Construct the energy-momentum tensor for this theory. Notice that the usual procedure does not result in a symmetric tensor. Show that you can remedy that using a tensor  $K^{\lambda\mu\nu}$  antisymmetric in its first two indices (you will have to build a new energy-momentum tensor  $\hat{T}^{\mu\nu} = T^{\mu\nu} + \ldots$  where the dots stays for a function of  $K^{\lambda\mu\nu}$ ). In particular show that

$$K^{\lambda\mu\nu} = F^{\mu\lambda}A^{\nu} \ ,$$

leads to the standard formulae for the electromagnetic energy and momentum densities:

$$\mathcal{E} = \frac{1}{2}(E^2 + B^2)$$
,  $\mathbf{S} = \mathbf{E} \times \mathbf{B}$ 

- **3.** Recall that  $T(a)^{-1}\phi(x)T(a) = \phi(x-a)$ , where  $T(a)\exp(iP^{\mu}a_{\mu})$  is the spacetime translation operator, and  $P^0$  is identified as the hamiltonian H.
  - **3.a)** Let  $a^{\mu}$  be infinitesimal, and derive an expression for  $[\phi(x), P^{\mu}]$ .
  - **3.b)** Show that the time component of your result is equivalent to the Heisenberg equation of  $i\dot{\phi} = [\phi, H]$ .
  - **3.c)** For a free field, use the Heisenberg equation to derive the Klein- Gordon equation.
  - 3.d) Define a spatial momentum operator

$$\mathbf{P} = -\int d^3x \Pi(x) \nabla \phi(x) \ ,$$

Use the canonical commutation relations to show that P obeys the relation you derived in part **3.a**).

- **3.e)** Express **P** in terms of  $a(\mathbf{k})$  and  $a^{\dagger}(\mathbf{k})$ .
- 4. Problem 2.2 of Peskin and Schroeder's book.