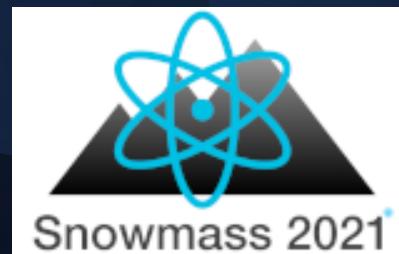


Building the case for future explorations at the Energy Frontier



Laura Reina (FSU)
The University of Oklahoma
October 26, 2023



From the work of the Snowmass 2021-22 Energy Frontier

Conveners: Meenakshi Narain, Laura Reina, Alessandro Tricoli

Ten Topical Groups focused on Electroweak, QCD, BSM physics

QCD and S.I. BSM Higgs, EW, top

| Topical Group | Co-Conveners | | |
|---|--------------------------------------|------------------------------|-----------------------------------|
| EF01: EW Physics: Higgs Boson properties and couplings | Sally Dawson (BNL) | Caterina Vernieri (SLAC) | |
| EF02: EW Physics: Higgs Boson as a portal to new physics | Patrick Meade (Stony Brook) | Isobel Ojalvo (Princeton) | |
| EF03: EW Physics: Heavy flavor and top quark physics | Reinhard Schwienhorst (MSU) | Doreen Wackerlo (Buffalo) | |
| EF04: EW Physics: EW Precision Physics and constraining new physics | Alberto Belloni (Maryland) | Ayres Freitas (Pittsburgh) | Junping Tian (Tokyo) |
| EF05: QCD and strong interactions: Precision QCD | Michael Begel (BNL) | Stefan Hoeche (FNAL) | Michael Schmitt (Northwestern) |
| EF06: QCD and strong interactions: Hadronic structure and forward QCD | Huey-Wen Lin (MSU) | Pavel Nadolsky (SMU) | Christophe Royon (Kansas) |
| EF07: QCD and strong interactions: Heavy Ions | Yen-Jie Lee (MIT) | Swagato Mukherjee (BNL) | |
| EF08: BSM: Model specific explorations | Jim Hirschauer (FNAL) | Elliot Lipeles (UPenn) | Nausheen Shah (Wayne State) |
| EF09: BSM: More general explorations | Tulika Bose (U Wisconsin-Madison) | Zhen Liu (Maryland) | Simone Griso (LBL) |
| EF10: BSM: Dark Matter at colliders | Caterina Doglioni (Lund) | LianTao Wang (Chicago) | Antonio Boveia (Ohio State) |

Liaisons, task forces, cross-frontier fora, contributed papers

| Other Frontier | Liaisons |
|-------------------------------|---|
| Neutrino Physics Frontier | André de Gouvêa (Northwestern) |
| Rare Processes and Precision | Manuel Franco Sevilla (Maryland) |
| Cosmic Frontier | Caterina Doglioni (Lund), Antonio Boveia (Ohio State) |
| Theory Frontier | Laura Reina (FSU) |
| Accelerator Frontier | Dmitri Denisov (BNL), Meenakshi Narain (Brown) |
| Computational Frontier | Peter Onyisi (U.Texas) |
| Instrumentation Frontier | Caterina Vernieri (SLAC), Maksym Titov (CEA Saclay) |
| Community Engagement Frontier | Daniel Whiteson (UCI), Sergei Gleyzer (Alabama) |

Muon Collider Forum Coordinators

EF: **Kevin Black** (U. Wisconsin-Madison), **Sergo Jindariani** (Fermilab)
AF: **Derun Li** (LBNL), **Diktys Stratakis** (Fermilab)
TF: **Patrick Meade** (Stony Brook U.), **Fabio Maltoni** (Louvain U., Bologna)

e+e- Collider Forum Coordinators

EF: **Maria Chamizo Llatas** (BNL), **Sridhara Dasu** (Wisconsin)
AF: **Emilio Nanni** (SLAC), **John Power** (ANL)
IF: **Ulrich Heintz** (Brown), **Steve Wagner** (Colorado)

Monte Carlo task force and production team

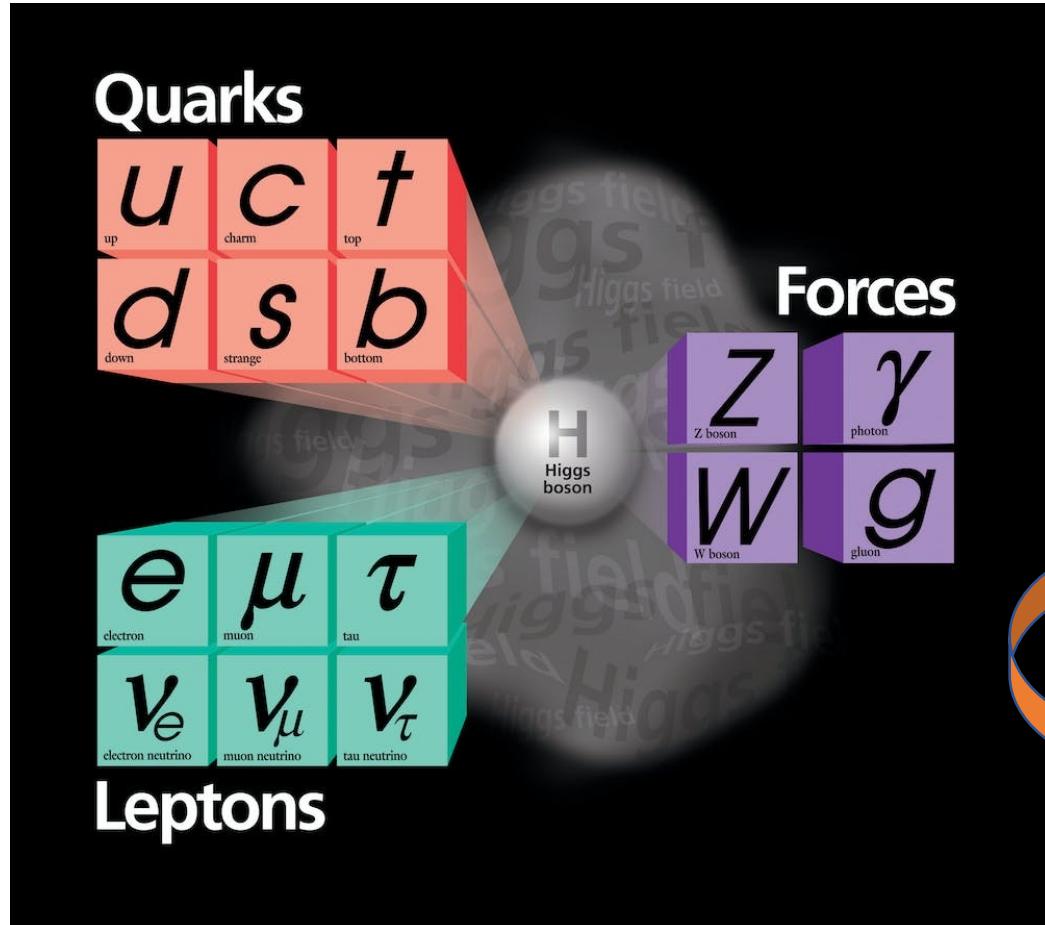
Coordinated by **John Stupak** (U. Oklahoma)
1) Assess the MC needs ⇒ “**Task force**”
2) Produce MC samples ⇒ “**Production Team**”

>160 Contributed Papers, TG and general frontier meetings, and
fora's activities served as basis to the EF main reports.

Energy Frontier: Exploring the TeV Scale and beyond

Through the breadth and multitude of
collider physics signatures

Exploring beyond the Standard Model of particle physics



A very minimal quantum field theory describing strong, weak, and electromagnetic interactions, based on a local (gauge) symmetry

$$SU(3)_C \times SU(2)_L \times U(1)_Y \rightarrow SU(3)_C \times U(1)_Q$$

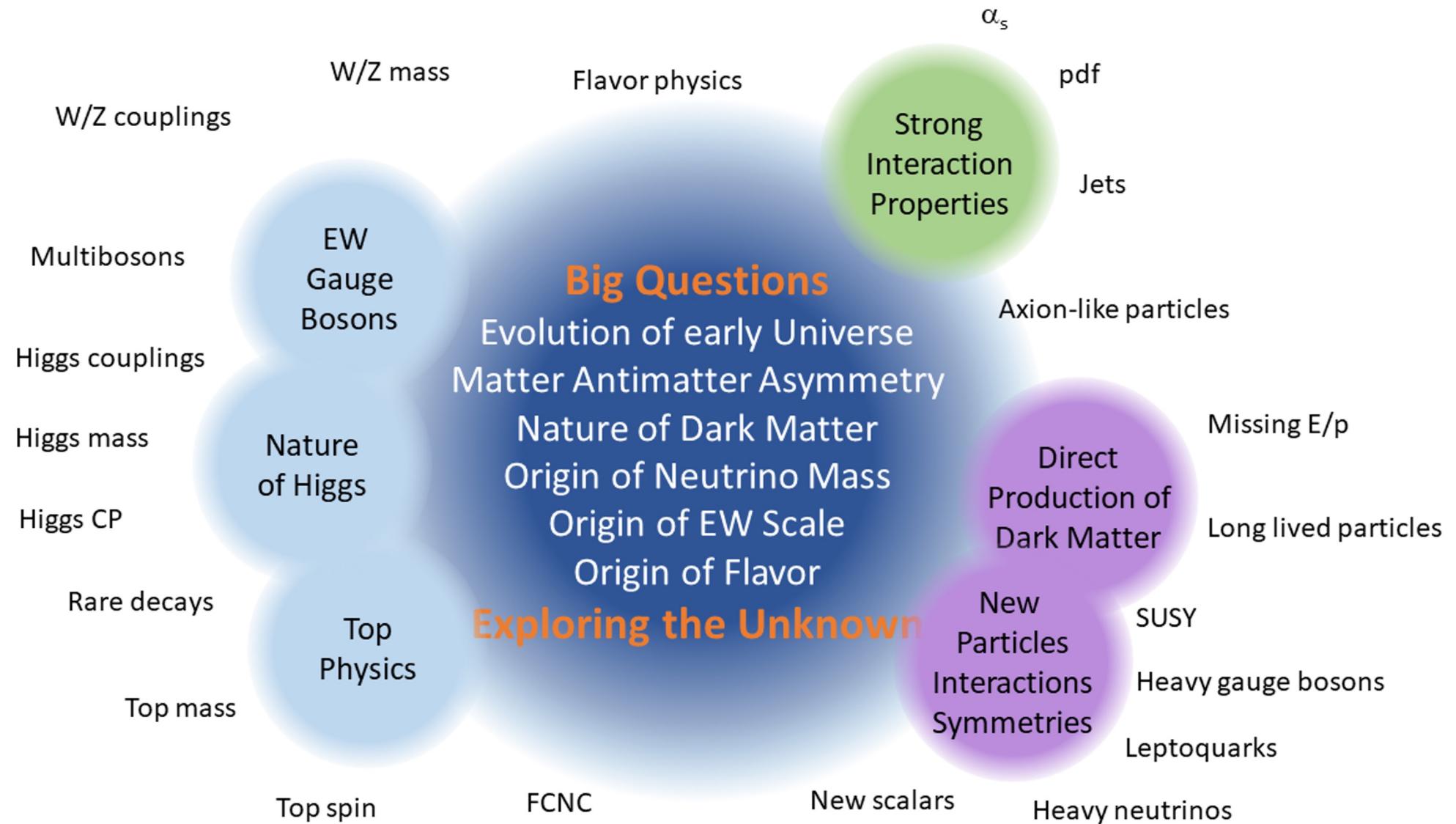
Strong interactions: gluons $\rightarrow m_g = 0$

Electromagnetic interactions: photon $\rightarrow m_\gamma = 0$

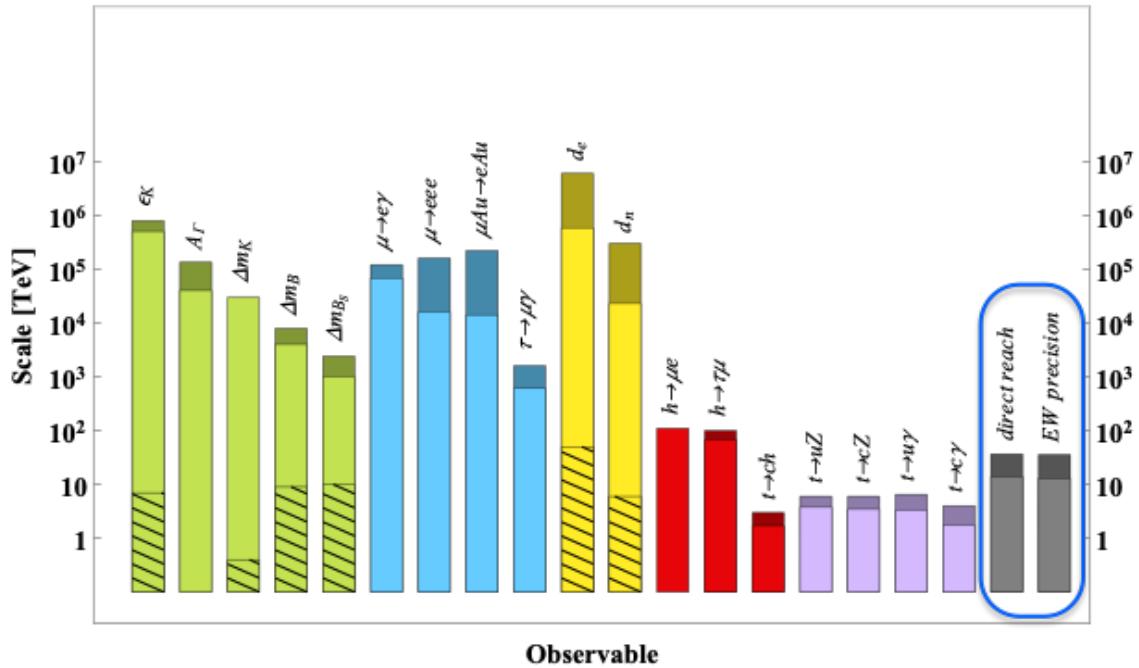
Weak interactions: W^\pm and $Z \rightarrow M_W, M_Z \neq 0$

Due to the presence of a scalar field whose potential spontaneously breaks the gauge symmetry of electroweak interactions and gives origin to massive gauge bosons (W,Z)

The Higgs boson (H) is the physical particle associated with such field



Emphasizing the breadth of collider physics



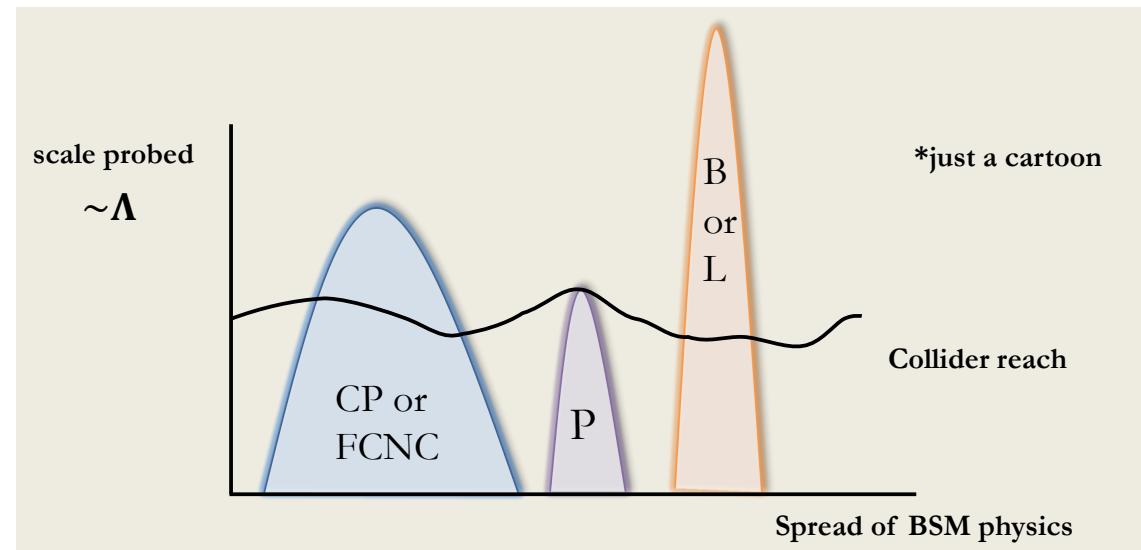
[European Strategy, arXiv:1910.11775]

Any new physics hypothesis will have to stand the test of colliders

Colliders may not be able to indirectly probe scales as high as e.g. flavor physics, but they provide a huge spectrum of measurements



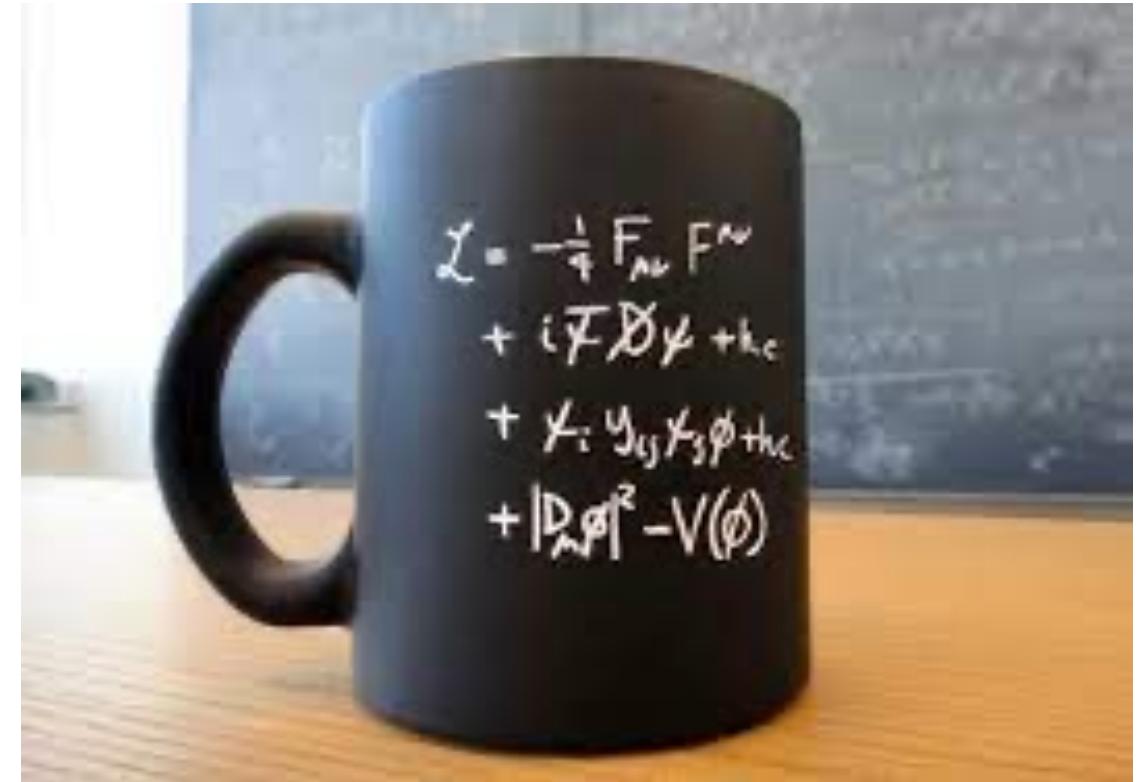
complementarity



[J. de Vries, talk at Snowmass 21 CPM]

The Standard Model: so simple that it can fit on a mug!

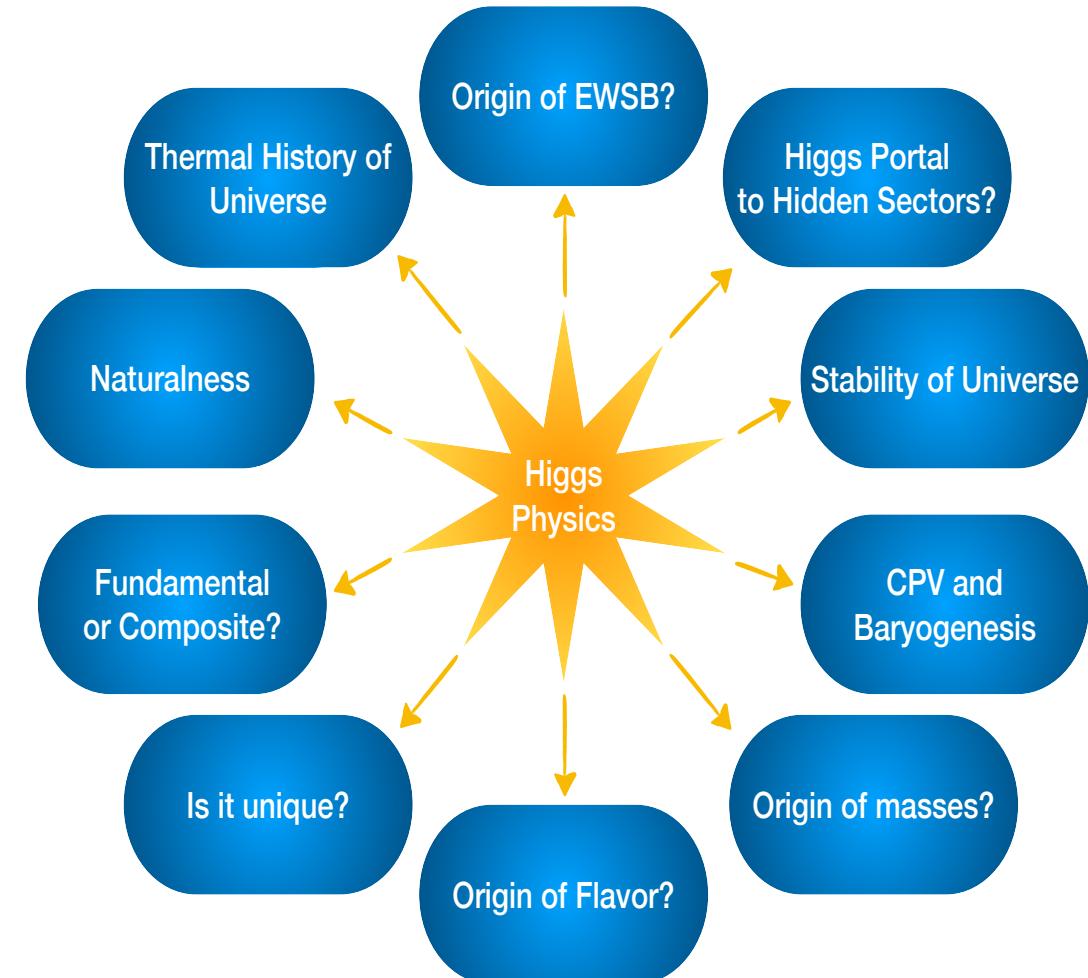
$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + i \bar{\psi} \gamma^\mu \psi + h.c. + Y_i Y_{ij} Y_j \phi + h.c. + |D_\mu \phi|^2 - V(\phi)$$



... and half of it is about Higgs!

Higgs physics
identified as
central to the
Energy Frontier
physics program

Unique link to
BSM physics



The origin of the SM Higgs pattern escapes the SM itself

The origin of SSB and ultimately of the EW scale is unexplained by the SM

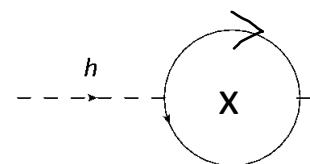
➤ Why the Higgs potential? Why $\mu^2 < 0$?

- Dynamical origin? What induces it?
- Cubic and quartic couplings, same λ ?

$$V(\phi) = \mu^2 \phi^\dagger \phi + \lambda(\phi^\dagger \phi)^2$$

➤ Why $M_H = 125$ GeV? \rightarrow Hierarchy problem – Naturalness

- Mass of scalar not protected by symmetry, receives large quantum corrections


$$\Delta M_H^2 \propto \pm \frac{\lambda_X}{16\pi^2} M_X^2$$

Yukawa interactions depends on arbitrary parameters, unexplained by the SM

➤ Why the hierarchy of Yukawa couplings \leftrightarrow fermion masses?

➤ Why flavor diagonal scalar couplings? Why only one scalar?

➤ Other sources of flavor mixing and CP violation?

➤ A new force all together?

$$L_{Yukawa} = y_{ij} \bar{\psi}_L^i \phi \psi_R^j + h.c.$$

fermion masses

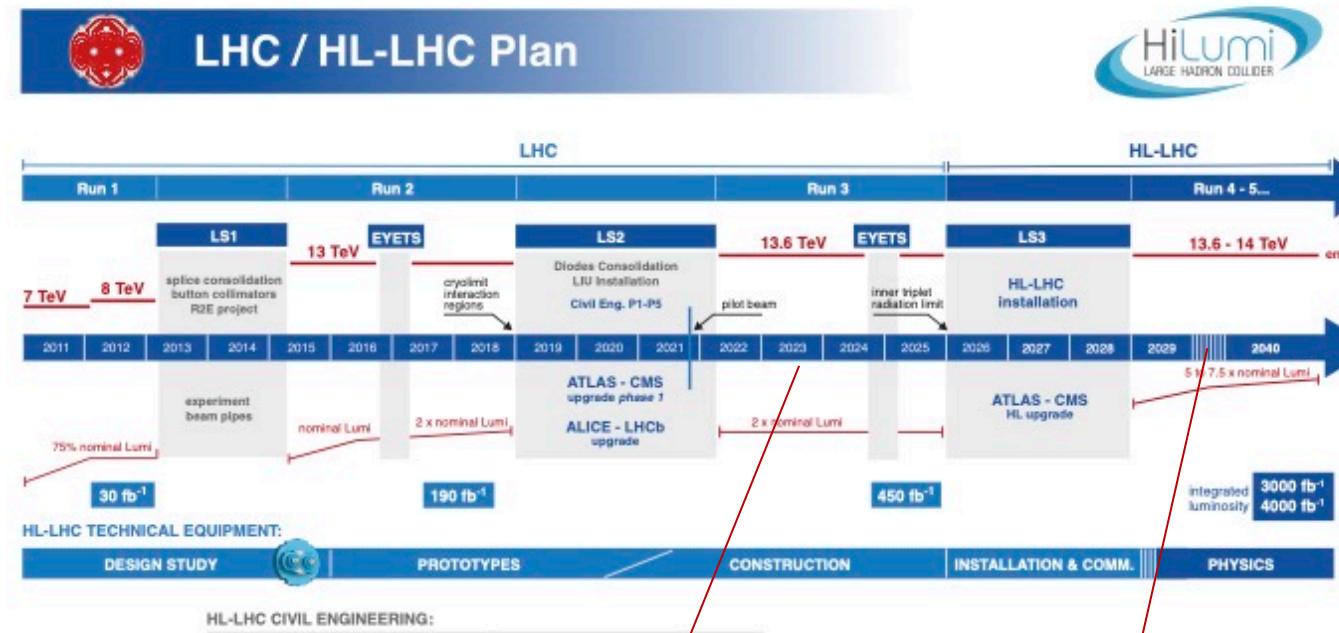
$$y_{ij} \rightarrow \frac{m_f}{v} \delta_{ij}$$

Yukawa couplings

$$\phi \rightarrow H + v$$

The LHC and its legacy

Ten years of LHC physics and looking ahead



We are only here

Many years of HL running ahead of us

- 2-fold increase in statistics by the end of Run 3
- 20-fold increase in statistics by the end of HL-LHC!

Higgs physics has been at the core of the LHC physics program

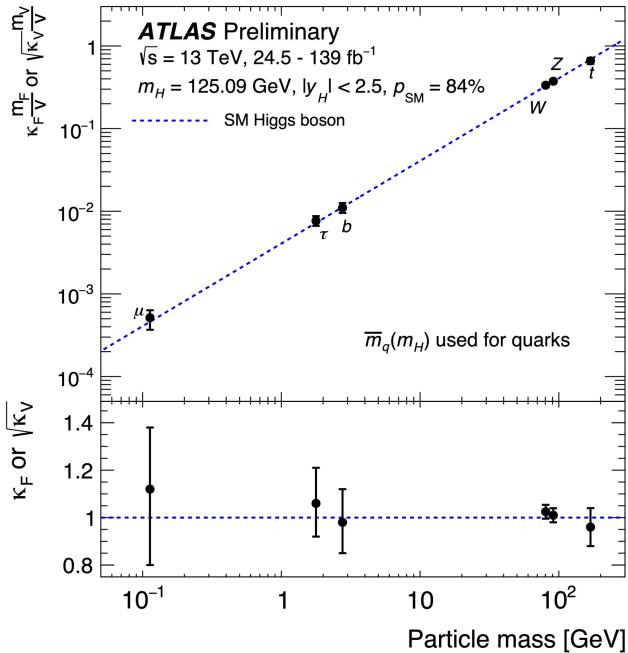
Snowmass 2013/Previous P5

- **Run 1: Higgs discovery**
- **Run 2: Higgs couplings**
 - outperformed expectations
- **Run 3 to HL-LHC**
- **Higgs precision program**

Snowmass 2021/Current P5

Run 2

Zooming in on couplings to probe the TeV scale



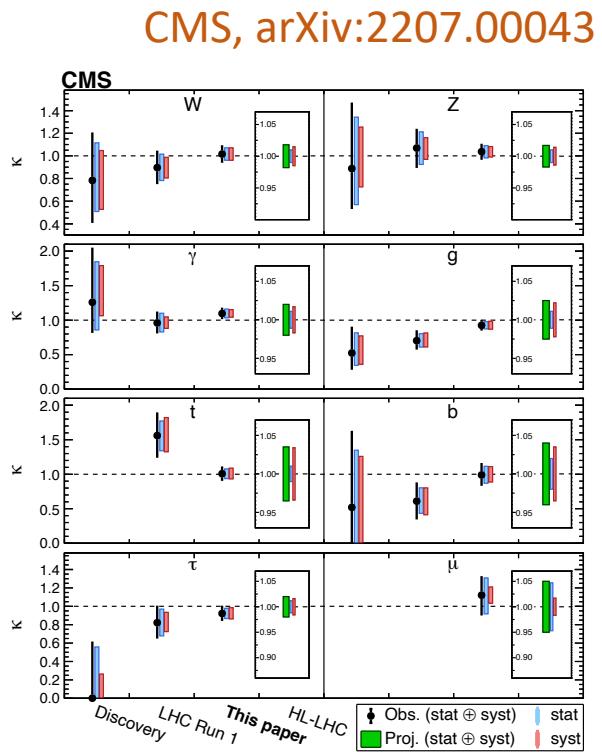
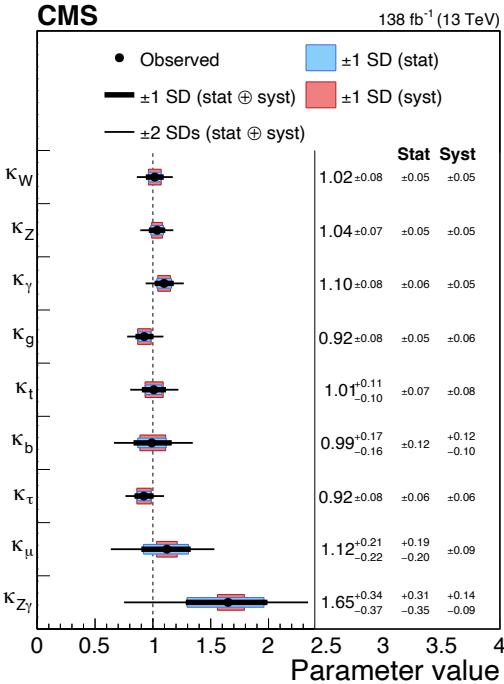
$$\kappa = g_x / g_x^{\text{SM}} = 1 + \Delta \kappa$$

$$\Delta \kappa \propto v^2 / \Lambda_{\text{BSM}}^2$$

Precision on $\Delta \kappa$



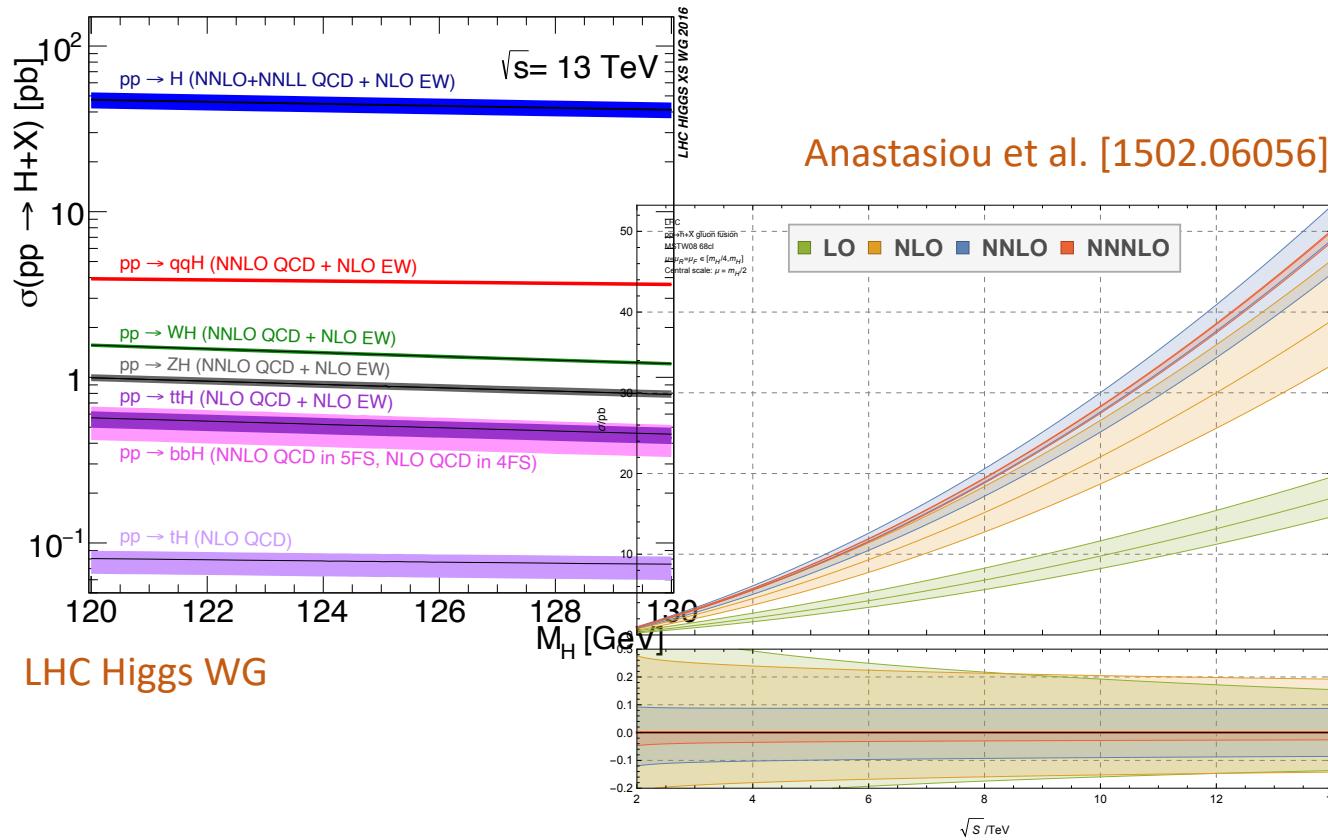
reach for Λ_{BSM}



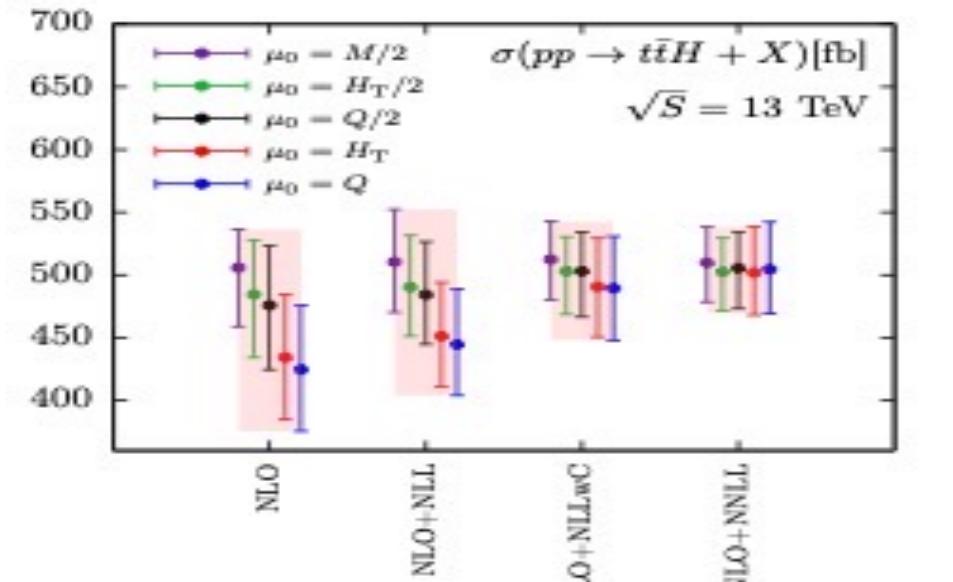
- Couplings to W/Z at 5-10 %
- Couplings to 3rd generation to 10-20%
- First measurements of 2nd generation couplings

- HL-LHC projections from partial Run 2 data (YR):
 - 2-5 % on most couplings
 - < 50% on Higgs self-coupling.
- Full Run2 results drastically improve partial Run 2 results: better projections expected

Theory has also come a long way



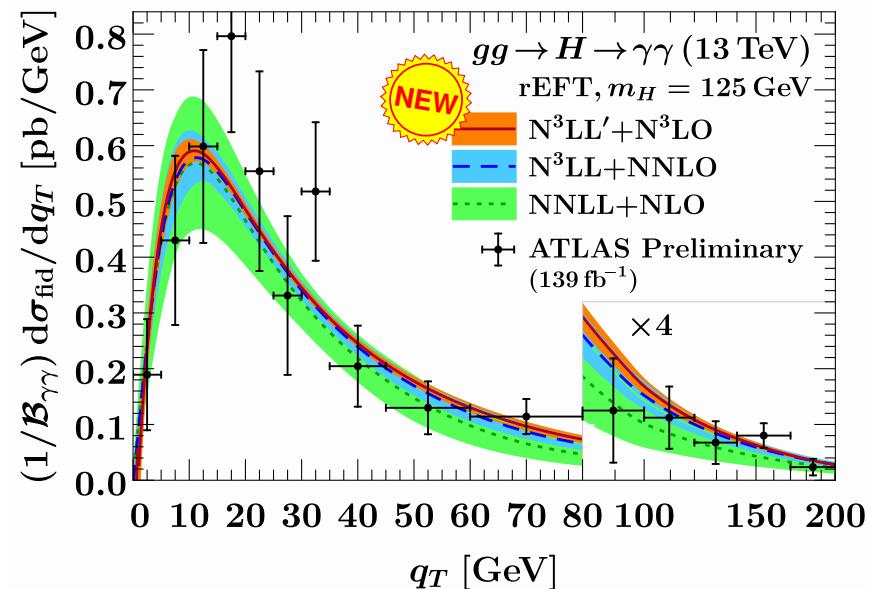
Several backgrounds also known at
NLO QCD+EW or improved NLO (+NNLL)
(e.g. W/Z+j, ttbb, ttW, ttZ, tt γ , ...)



Anastasiou et al. [1502.06056]

Kulesza et al. [1812.08622]

Bliss et al. [2102.08039]



Beyond SM-coupling rescaling

Model new physics by extending the SM Lagrangian by effective interactions (ex. SM EFT)

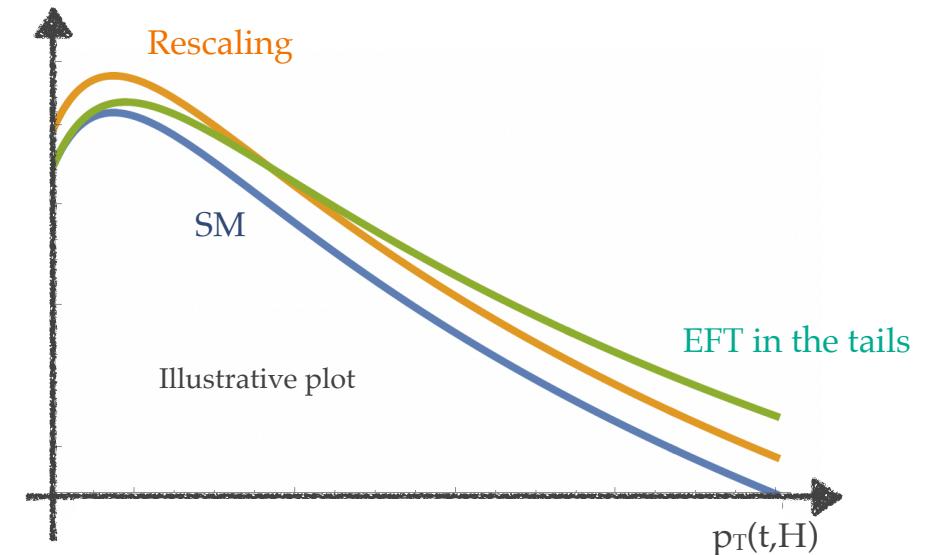
$$\mathcal{L}_{\text{SM}}^{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_{d>4} \frac{1}{\Lambda^{d-4}} \mathcal{L}_d = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda} \mathcal{L}_5 + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots$$

$$\mathcal{L}_d = \sum_i C_i^{(d)} \mathcal{O}_i^{(d)}, \quad \left[\mathcal{O}_i^{(d)} \right] = d$$

Under the assumption that new physics leaves at scales $\Lambda > \sqrt{s}$

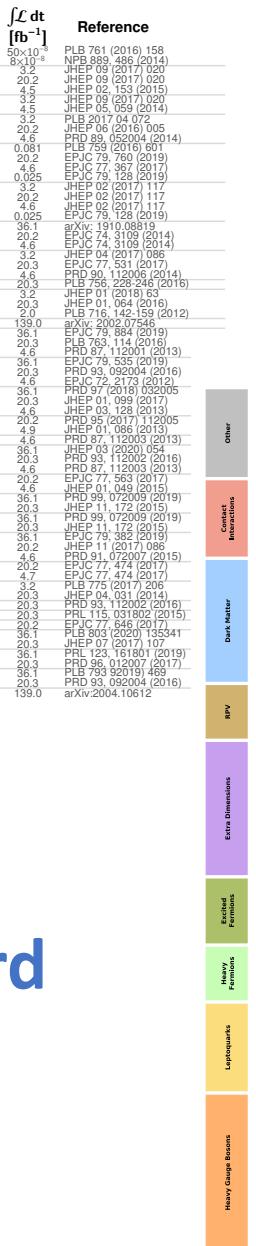
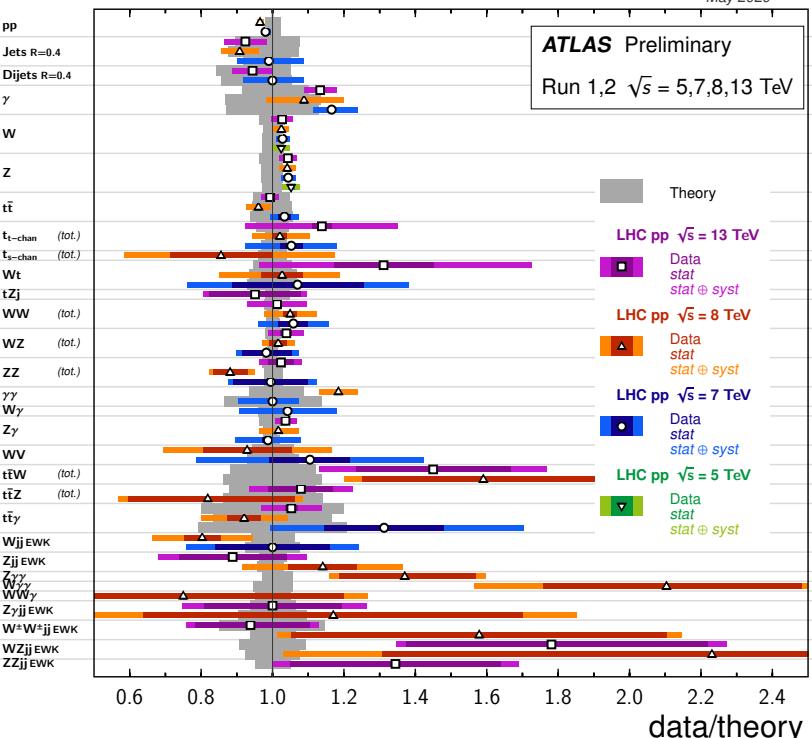
Expansion in $(v, E)/\Lambda$: affects all SM observables at both low and high energy

- **SM masses and couplings** → **rescaling**
- **Shapes of distributions** → more visible in **tails of distributions**

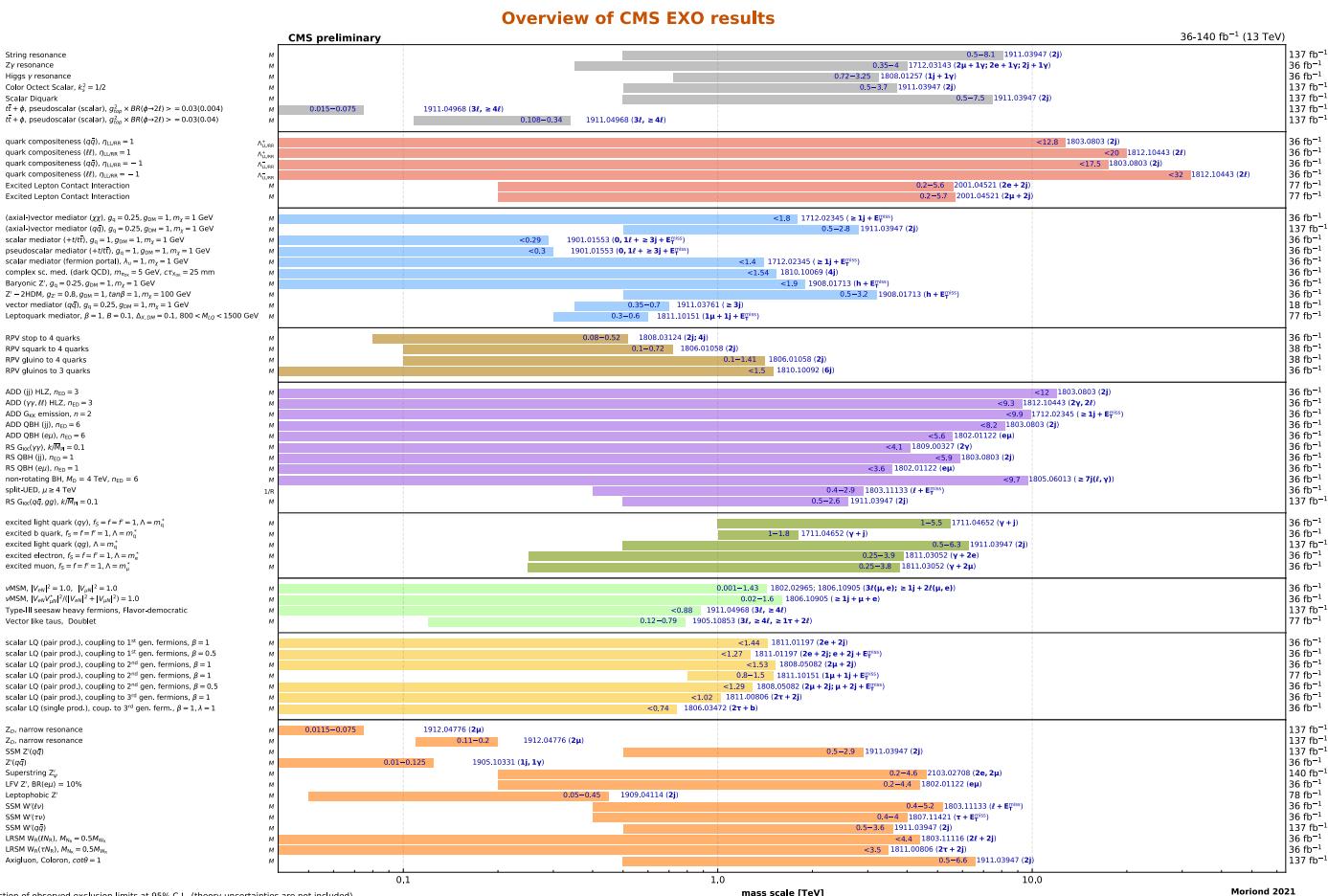


Standard Model Production Cross Section Measurements

Status:
May 2020



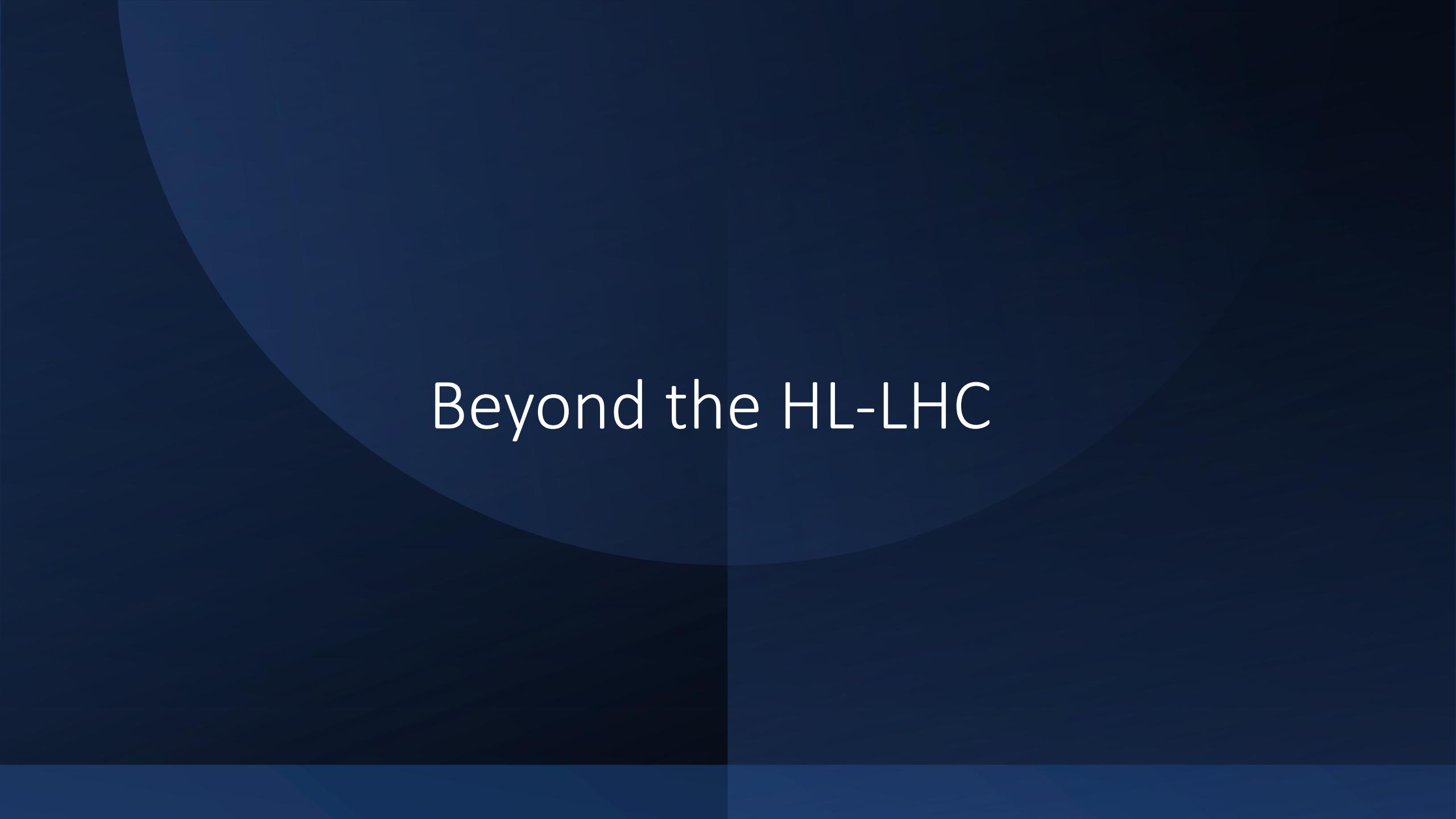
... to which we should add a unique spectrum of SM measurements and BSM direct searches!



What is the path forward
beyond the HL-LHC?

Selection of observed exclusion limits at 95% C.L. (theory uncertainties are not included).

Moriond 2021



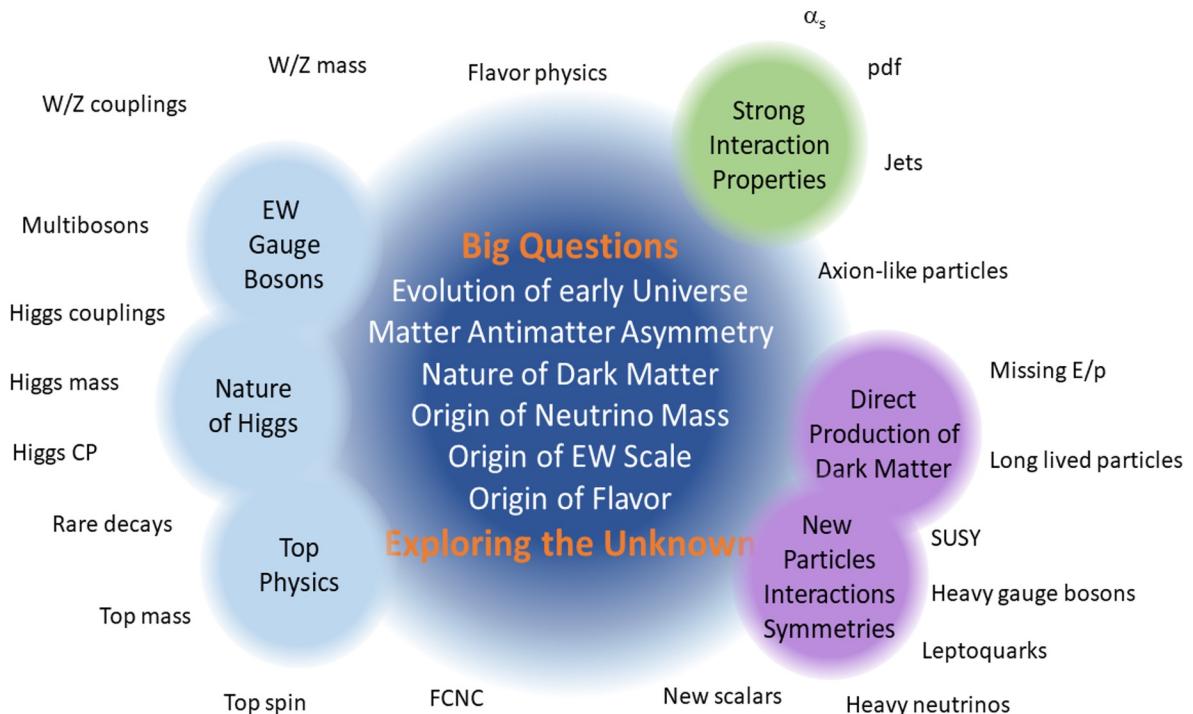
Beyond the HL-LHC

Addressing the “Big Questions” and “Exploring the unknown” are the main scientific goals of the EF

Should be pursued following

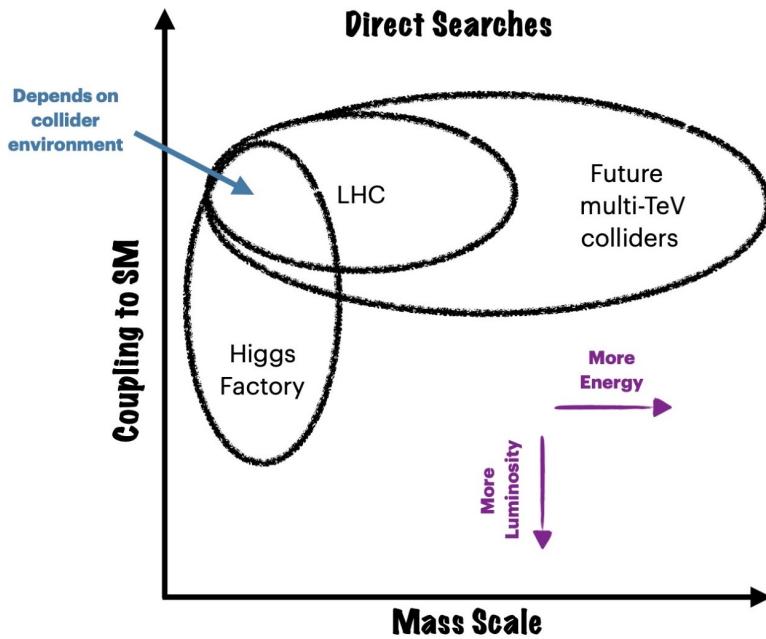
Two main avenues

- **Study known phenomena at high energies looking for indirect evidence of BSM physics**
 - Need **factories of Higgs bosons** (and other SM particles) to probe the TeV scale via precision measurements
- **Search for direct evidence of BSM physics at the energy frontier**
 - Need **multi-TeV colliders**



Beyond the HL-LHC: Precision and Energy

New physics can be at low as at high mass scales,
Naturalness would prefer scales close to the EW scale, but
the LHC has already placed **strong bounds** around 1-2 TeV.

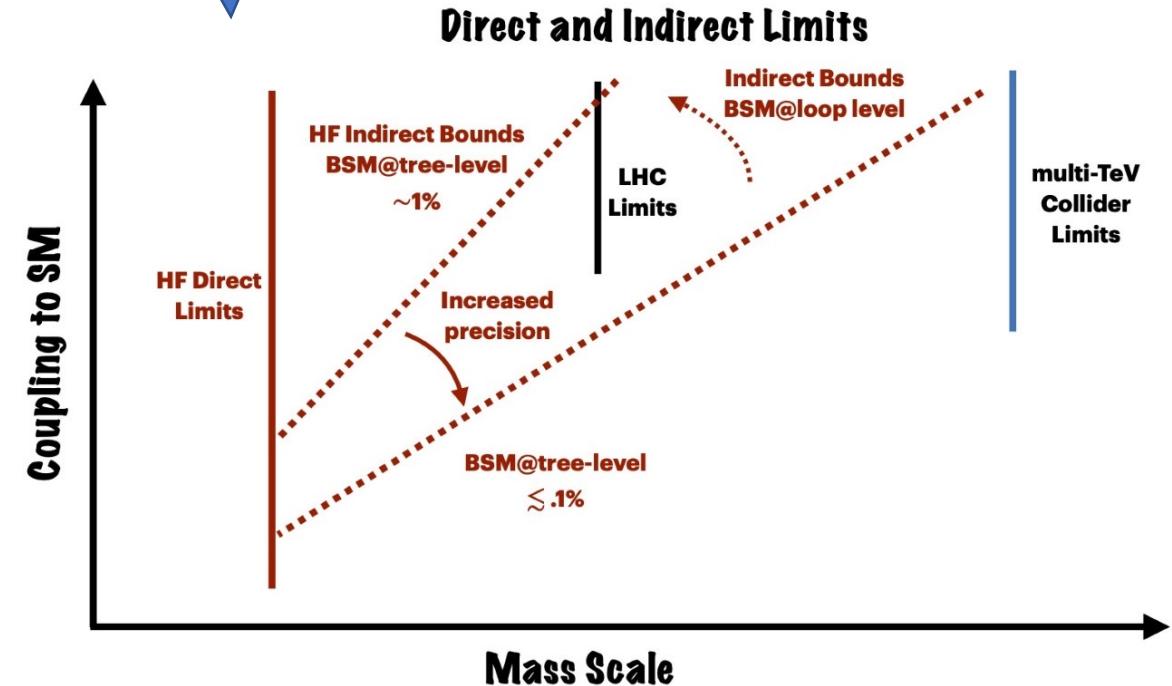


Higgs coupling measurements and direct searches
will complement each other in exploring the
1-10 TeV scale and beyond.

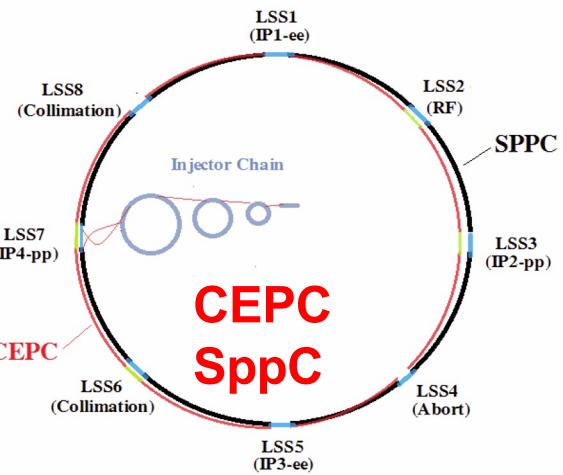
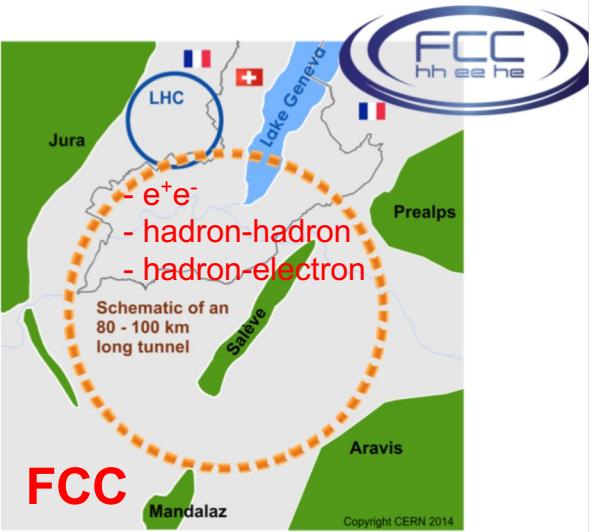
In a simplified picture:

New physics at **tree level**:
 $\delta\eta_{SM} \sim g_{BSM}^2 E^2/M^2$

New physics at **loop level**:
 $\delta\eta_{SM} \sim 1/16\pi^2 \times g_{BSM}^2 E^2/M^2$



EF future colliders



Hadrons

- large mass reach \Rightarrow exploration?
- $S/B \sim 10^{-10}$ (w/o trigger)
- $S/B \sim 0.1$ (w/ trigger)
- requires multiple detectors (w/ optimized design)
- only pdf access to \sqrt{s}
- \Rightarrow couplings to quarks and gluons

Circular

- higher luminosity
- several interaction points
- precise E-beam measurement ($O(0.1\text{MeV})$ via resonant depolarization)
- \sqrt{s} limited by synchrotron radiation

Leptons

- $S/B \sim 1 \Rightarrow$ measurement?
- polarized beams (handle to chose the dominant process)
- limited (direct) mass reach
- identifiable final states
- \Rightarrow EW couplings

Linear

- easier to upgrade in energy
- easier to polarize beams
- “greener”: less power consumption*
- large beamstrahlung
- one IP only

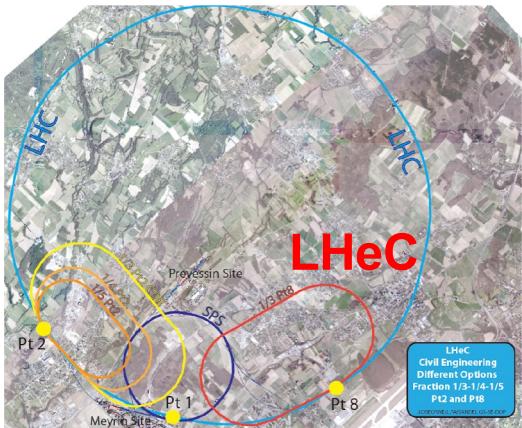
*energy consumption per integrated luminosity is lower at circular colliders but the energy consumption per GeV is lower at linear colliders

Inst. Pascal, Dec. 4, 2019

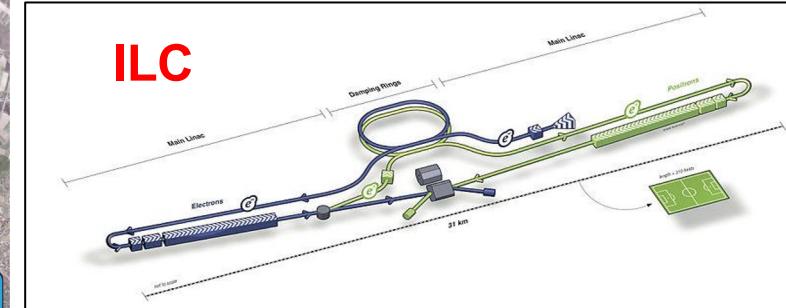
Christophe Grojean

Future Measurements

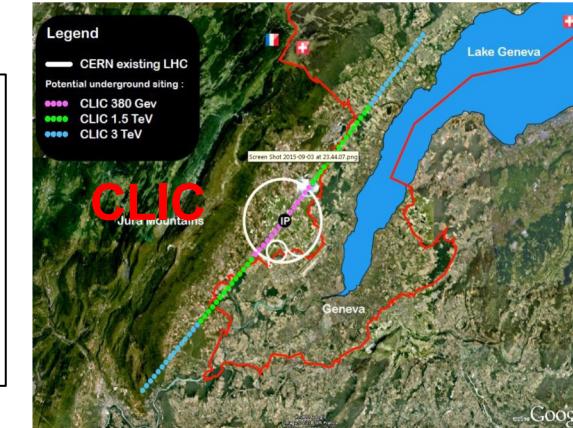
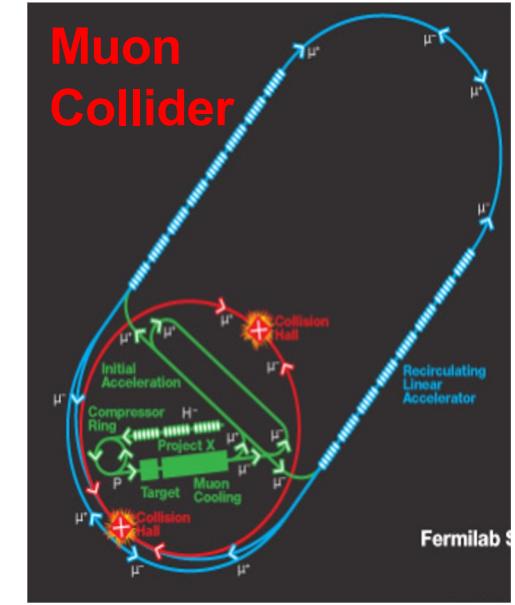
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ILC



SM21 added CCC (C³)



Higgs-boson factories (up to 1 TeV c.o.m. energy)

| Collider | Type | \sqrt{s} | $\mathcal{P}[\%]$ e^-/e^+ | \mathcal{L}_{int} ab^{-1}/IP | Start Date | Const. | Physics |
|----------------------|----------|-------------|--------------------------------|--|------------|--------|---------|
| HL-LHC | pp | 14 TeV | | 3 | | | 2027 |
| ILC & C ³ | ee | 250 GeV | $\pm 80/\pm 30$ | 2 | 2028 | 2038 | |
| | | 350 GeV | $\pm 80/\pm 30$ | 0.2 | | | |
| | | 500 GeV | $\pm 80/\pm 30$ | 4 | | | |
| | | 1 TeV | $\pm 80/\pm 20$ | 8 | | | |
| CLIC | ee | 380 GeV | $\pm 80/0$ | 1 | 2041 | | 2048 |
| CEPC | ee | M_Z | | 50 | 2026 | 2035 | |
| | | $2M_W$ | | 3 | | | |
| | | 240 GeV | | 10 | | | |
| | | 360 GeV | | 0.5 | | | |
| FCC-ee | ee | M_Z | | 75 | 2033 | 2048 | |
| | | $2M_W$ | | 5 | | | |
| | | 240 GeV | | 2.5 | | | |
| | | $2 M_{top}$ | | 0.8 | | | |
| μ -collider | $\mu\mu$ | 125 GeV | | 0.02 | | | |

Snowmass EF wiki: <https://snowmass21.org/energy/start>

Snowmass 21: EF Benchmark Scenarios

Multi-TeV colliders (> 1 TeV c.o.m. energy)

| Collider | Type | \sqrt{s} | $\mathcal{P}[\%]$ e^-/e^+ | \mathcal{L}_{int} ab^{-1}/IP | Start Date | Const. | Physics |
|-----------------|----------|------------|--------------------------------|--|------------|--------|---------|
| HE-LHC | pp | 27 TeV | | | 15 | | |
| FCC-hh | pp | 100 TeV | | | 30 | 2063 | 2074 |
| SppC | pp | 75-125 TeV | | | 10-20 | | 2055 |
| LHeC | ep | 1.3 TeV | | | 1 | | |
| FCC-eh | | 3.5 TeV | | | 2 | | |
| CLIC | ee | 1.5 TeV | $\pm 80/0$ | 2.5 | 2052 | 2058 | |
| | | 3.0 TeV | $\pm 80/0$ | 5 | | | |
| μ -collider | $\mu\mu$ | 3 TeV | | 1 | 2038 | 2045 | |
| | | 10 TeV | | 10 | | | |

Timelines are taken from the Collider ITF report (arXiv: 2208.06030)

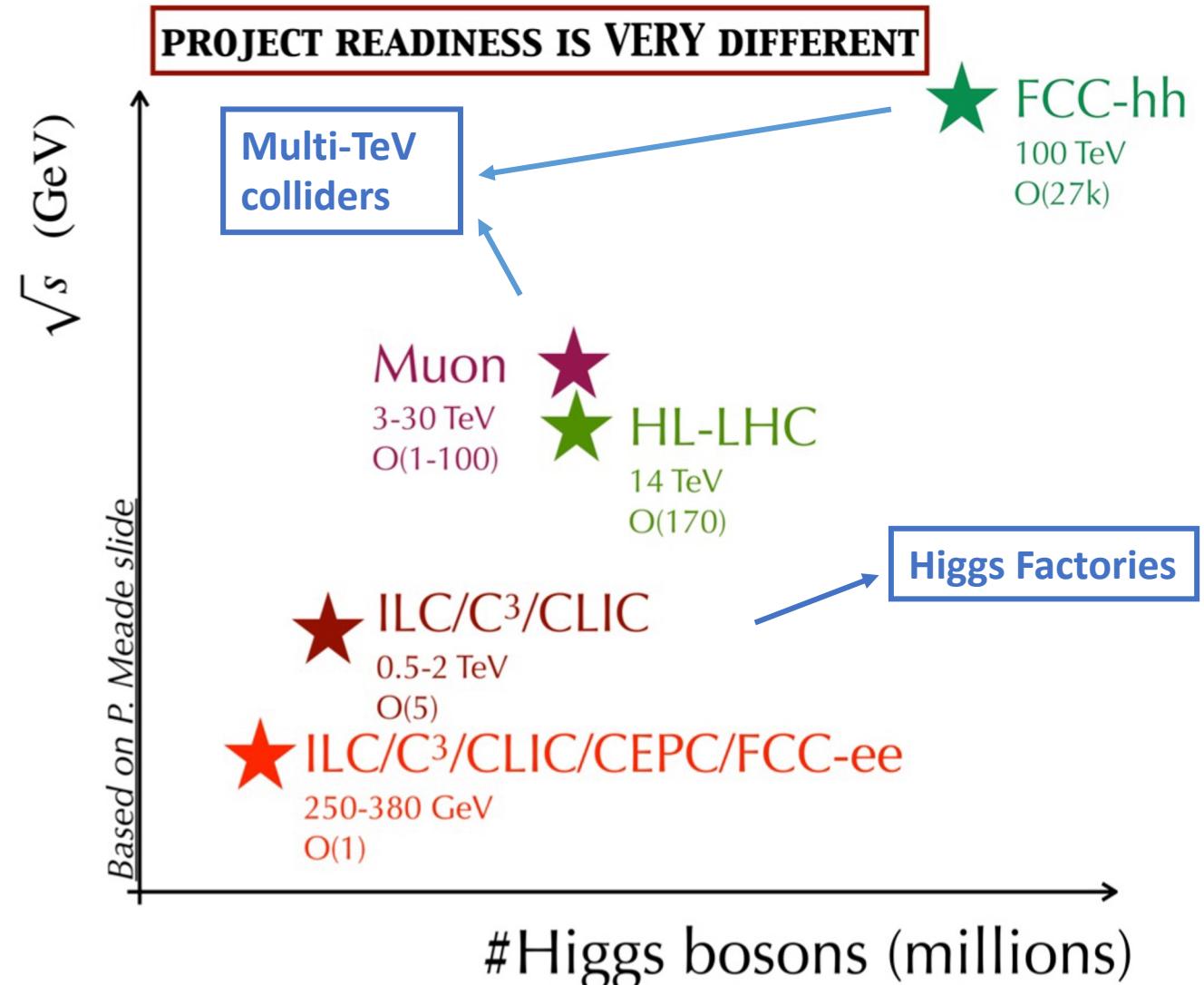
Beyond the HL-LHC: proposed future colliders

LEPTON COLLIDERS

- **Circular e+e-** (CEPC, FCC-ee)
 - **90-350 GeV**
 - *strongly limited by synchrotron radiation above 350– 400 GeV*
- **Linear e+e-** (ILC, CLIC, C³)
 - **250 GeV — > 1 TeV**
 - *Reach higher energies, and can use polarized beams*
- **$\mu^+\mu^-$**
 - **3-30 TeV**

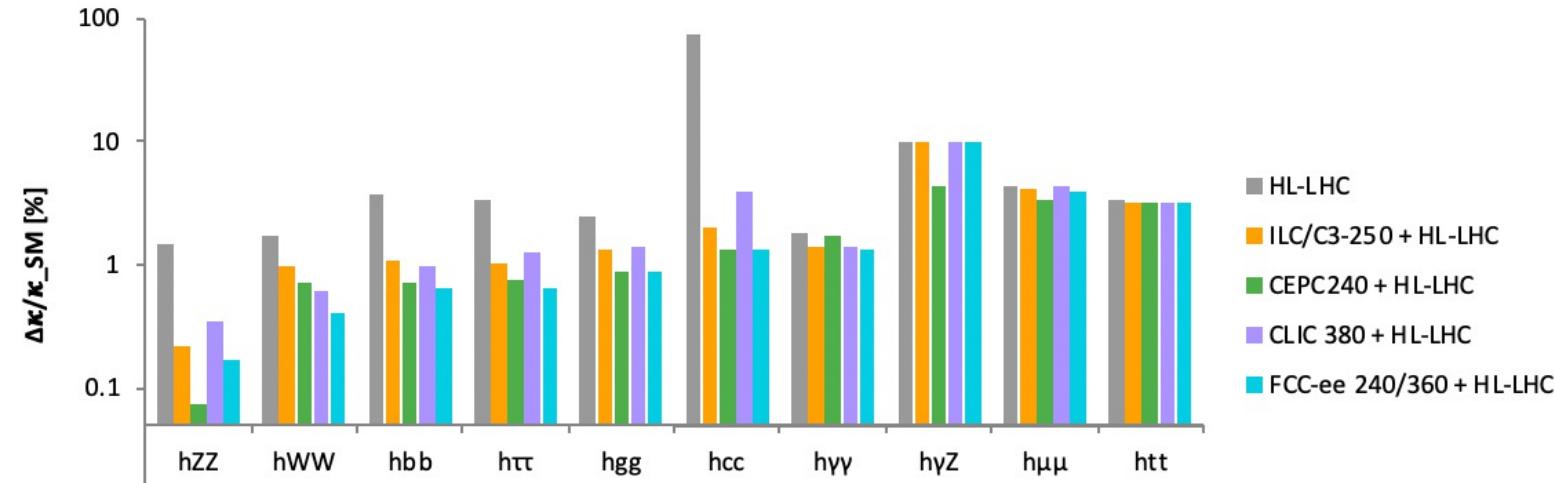
HADRON COLLIDERS

- **75-200 TeV** (FCC-hh)

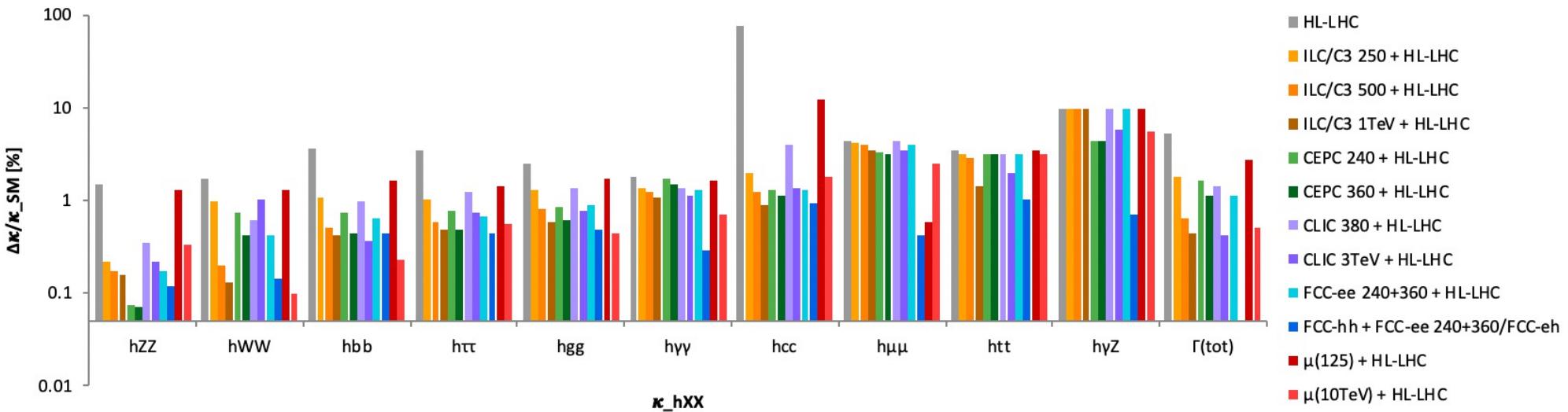


Reach of future colliders for Higgs couplings: a closer look

Based on full Run 2 dataset analyses



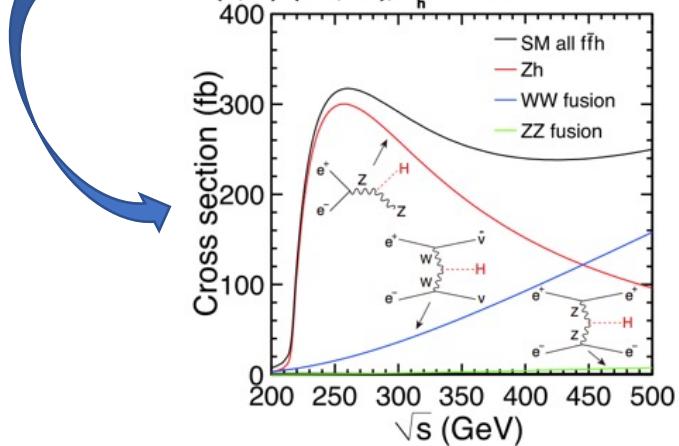
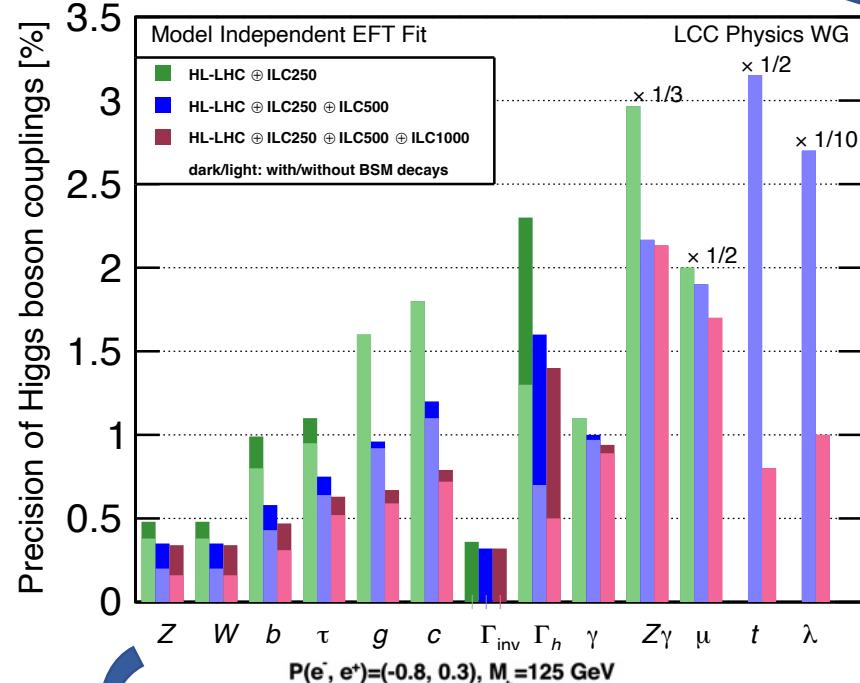
Initial stages of future
 e^+e^- machines



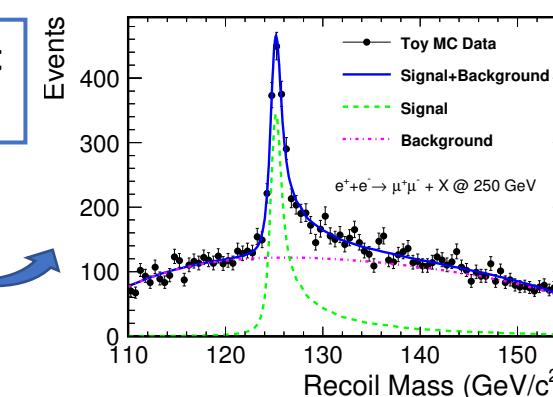
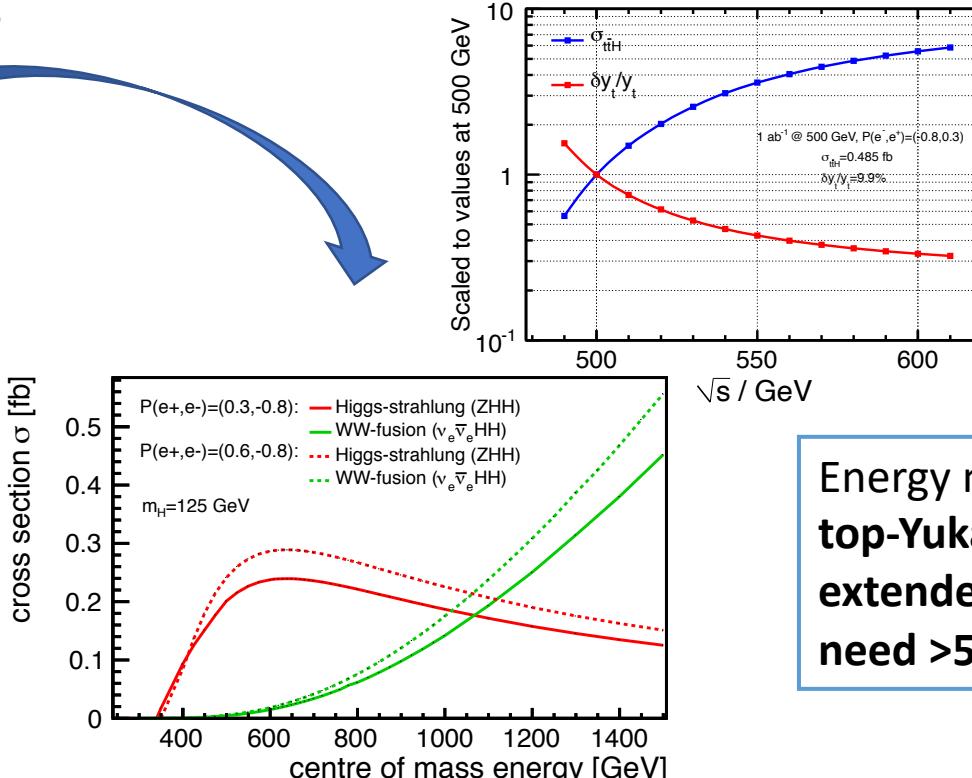
Final reach of all
considered
future colliders

From Snowmass 2021 EF
Higgs Topical Group Report
arXiv:2209.07510

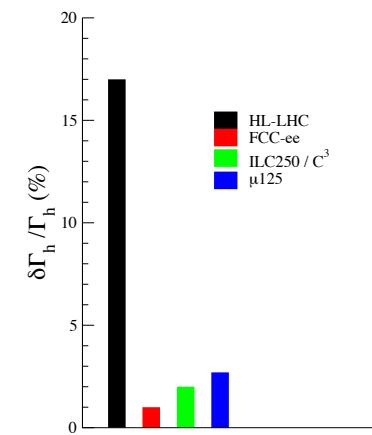
The case of e^+e^- Higgs factories



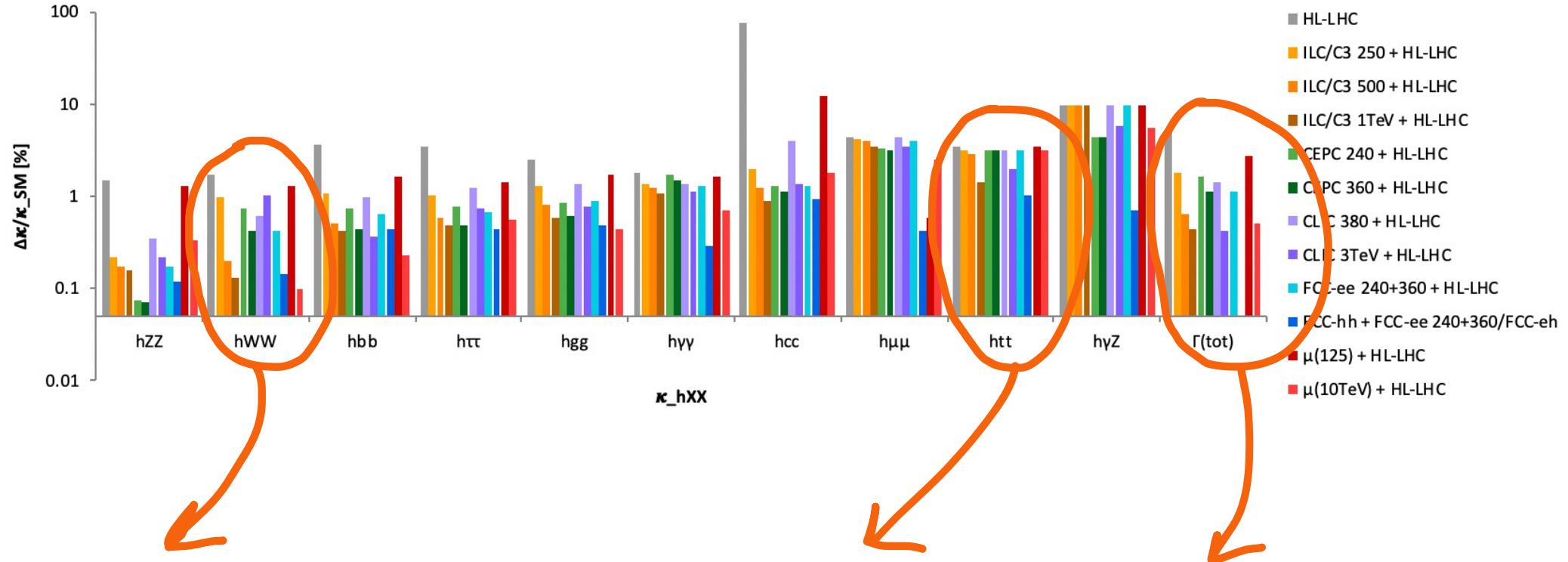
Model-independent
 Γ_h measurement



Energy matters
top-Yukawa, HH,
extended Higgs sectors
need >500 GeV



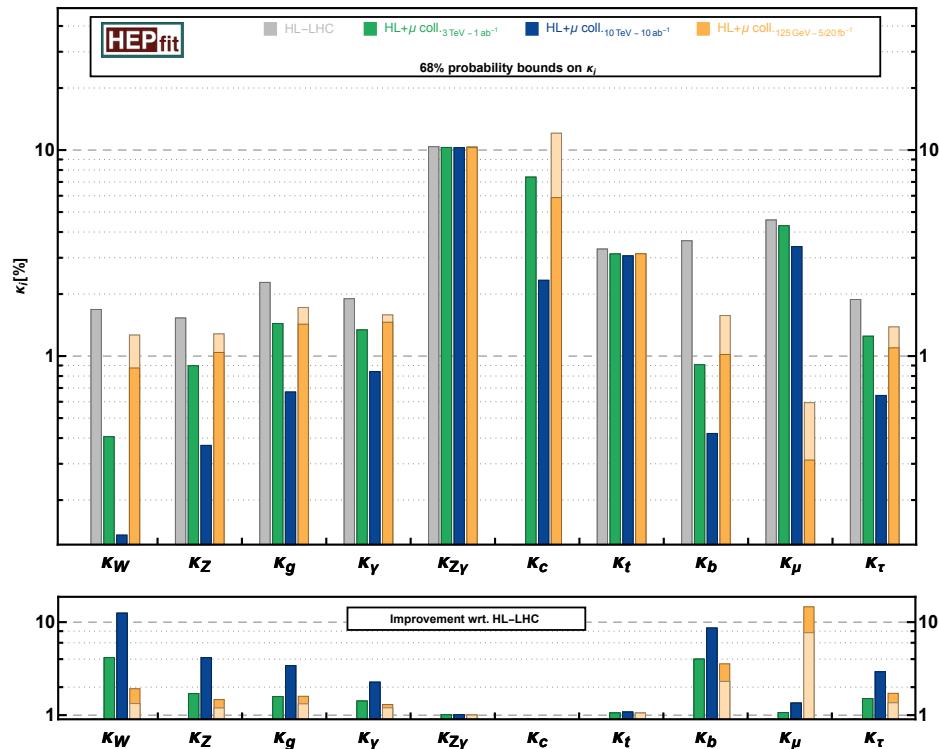
Focusing on final reach of e^+e^- machines



ILC/C³ reach beyond 500 GeV and upgrade to 1 TeV allows drastic improvements in measuring couplings to W and top as well as more precision in a model independent measurement of the total width.

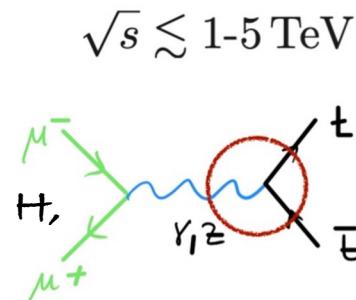
What about Higgs self-coupling?

The case of a Muon Collider



Snowmass 21 EF Higgs TG Report
(arXiv:2209.07510) &
MuC Forum Report
(arXiv:2209.01318)

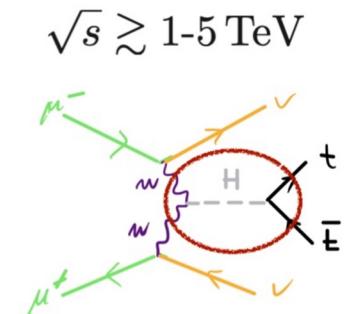
- Many stages/upgrades:
 - 125 GeV on-Higgs resonance
 - 3 TeV
 - 10 TeV
 - >10 TeV (14, 30, ... TeV)
- Lepton collider
 - Cleaner environment → [precision](#)
- ... but high energy
 - Pushing the EF → [discovery](#)
- Competitive/complementary to ~ 100 TeV hadron collider
- Contained size
 - $M_\mu \sim 200 m_e$ → reduced synchrotron radiation ($\times 1.6 \times 10^{-9}$)
- New physics regimes
 - $E > \Lambda_{EW}$
 - EW radiation



$$\sigma_s \sim \frac{1}{s}$$



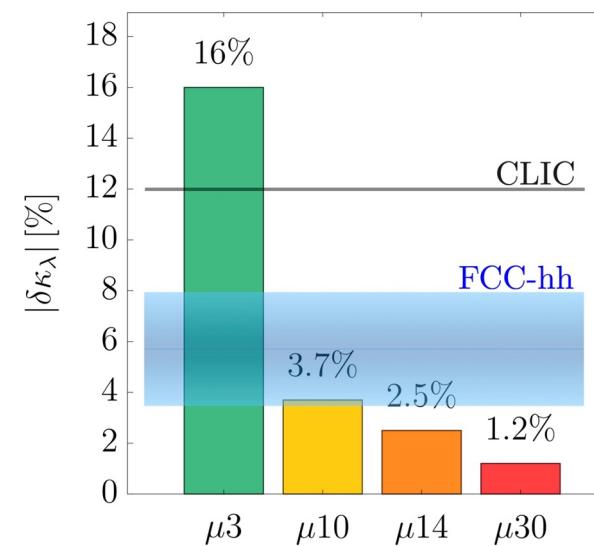
$$\sigma_s \sim \frac{1}{M^2} \log^n \frac{s}{M}$$



Reach for Higgs self-coupling

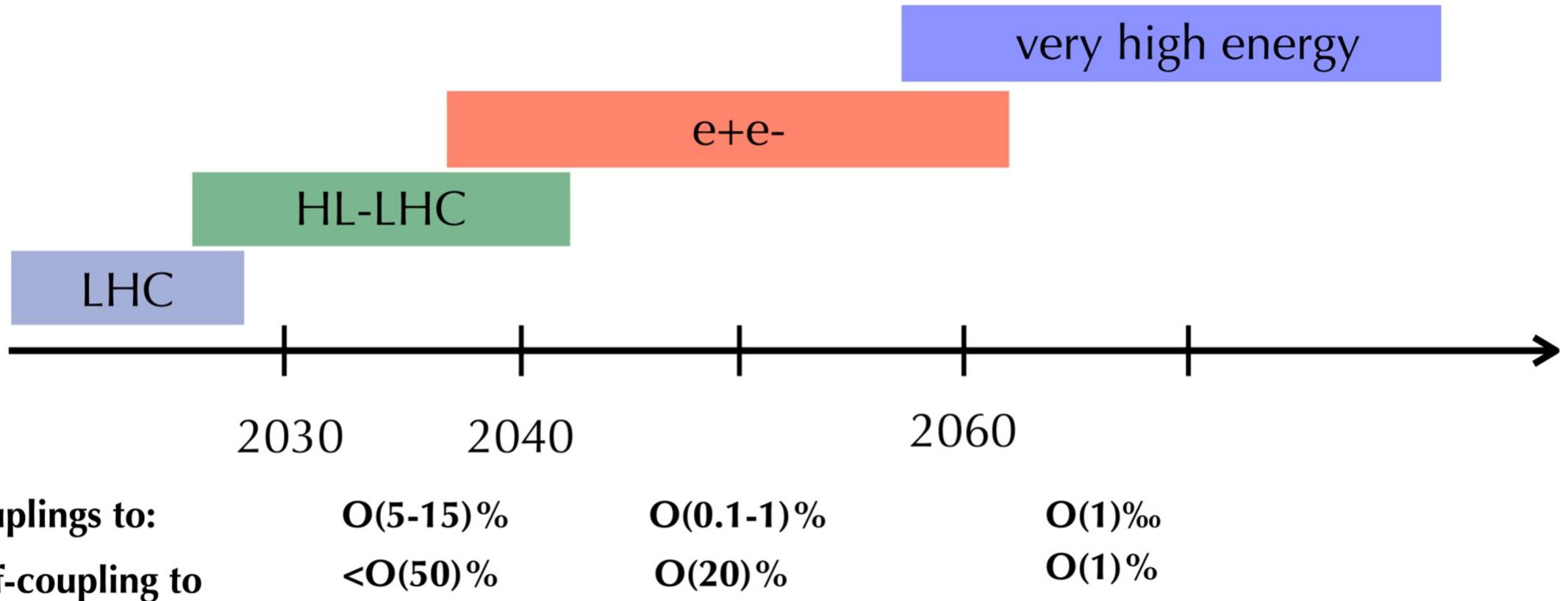
| collider | Indirect- h | hh | combined |
|---|---------------|----------|----------|
| HL-LHC | 100-200% | 50% | 50% |
| ILC ₂₅₀ /C ³ -250 | 49% | — | 49% |
| ILC ₅₀₀ /C ³ -550 | 38% | 20% | 20% |
| CLIC ₃₈₀ | 50% | — | 50% |
| CLIC ₁₅₀₀ | 49% | 36% | 29% |
| CLIC ₃₀₀₀ | 49% | 9% | 9% |
| FCC-ee | 33% | — | 33% |
| FCC-ee (4 IPs) | 24% | — | 24% |
| FCC-hh | - | 2.9-5.5% | 2.9-5.5% |
| μ (3 TeV) | - | 15-30% | 15-30% |
| μ (10 TeV) | - | 4% | 4% |

- ATLAS and CMS HL-LHC updated
- FCC-hh updated [arXiv:2004.03505](https://arxiv.org/abs/2004.03505)
- Added MuC reach:



[arXiv:2203.07256](https://arxiv.org/abs/2203.07256)

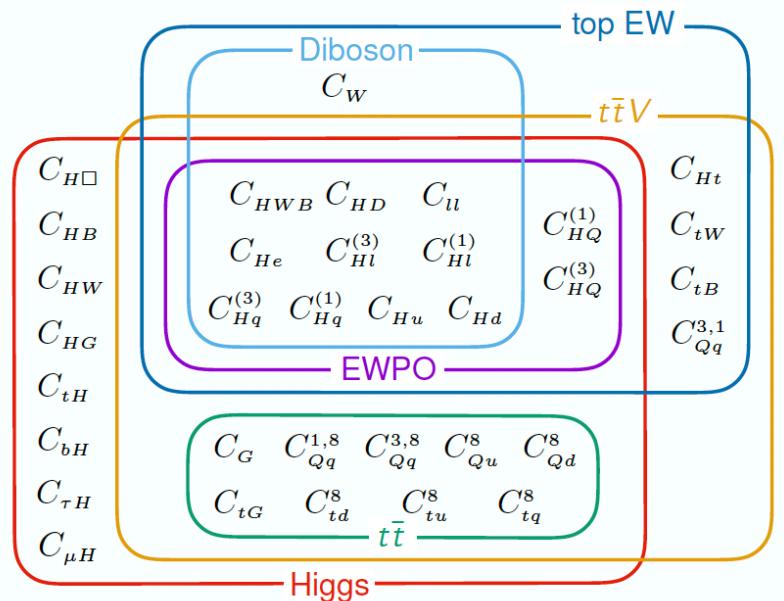
Beyond the HL-LHC: projections for Higgs couplings



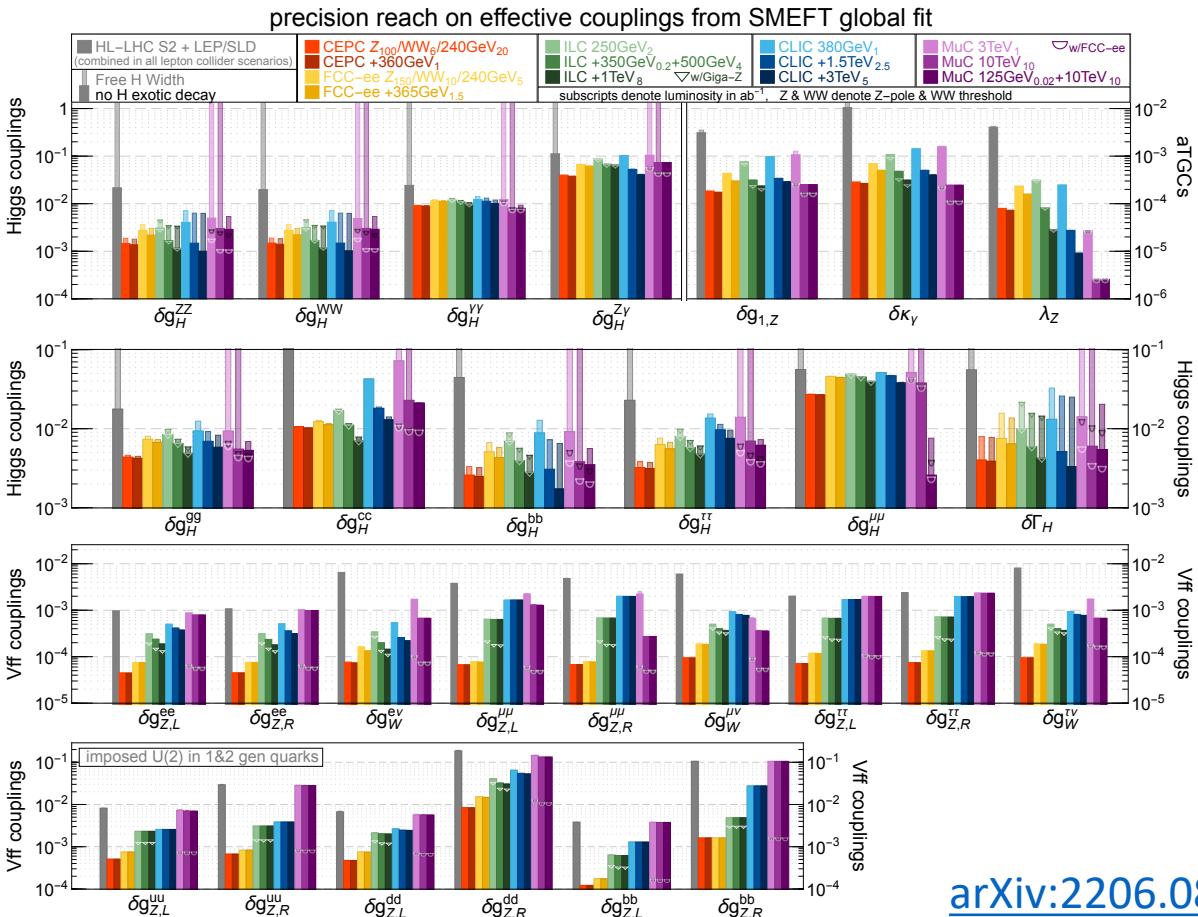
From C. Vernieri – Snowmass 21 EF Workshop - Brown U. - March 2022

Constraining BSM via global EFT fits

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \left(\frac{1}{\Lambda^2} \sum_i C_i O_i + \text{h.c.} \right) + O(\Lambda^{-4})$$



EW + Higgs

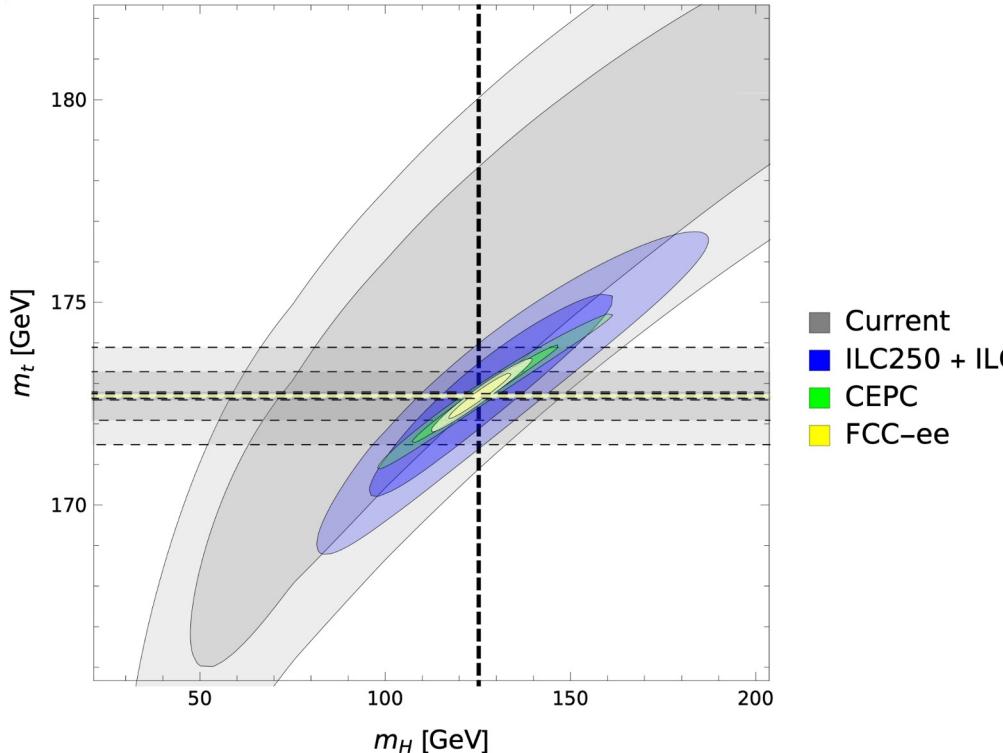


arXiv:2206.08326

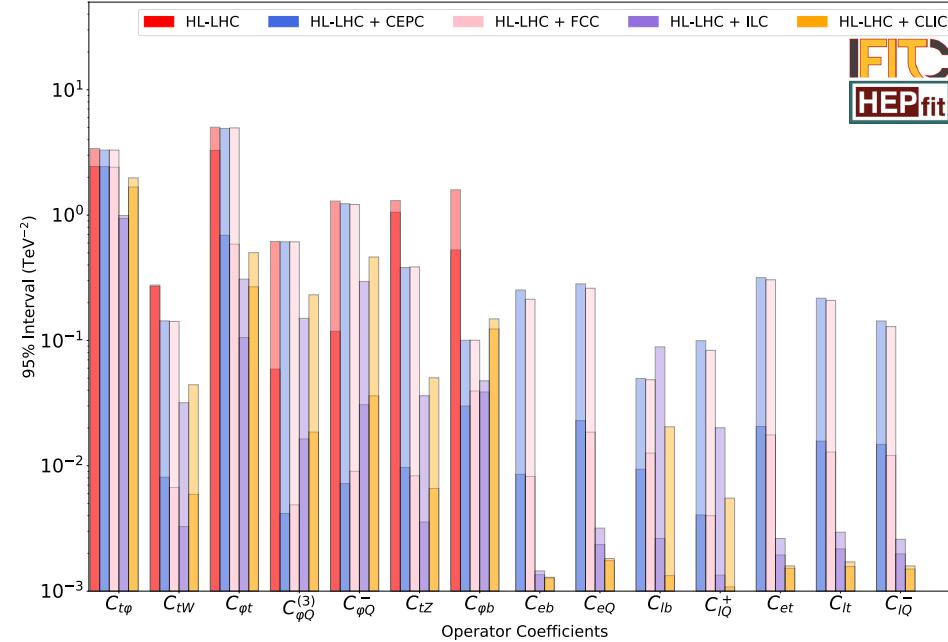
EFT connects different processes with large correlations: pattern of coefficients give insights on underlying BSM model

Interplay with top-quark precision measurements

Stress testing the SM and exploring anomalous couplings



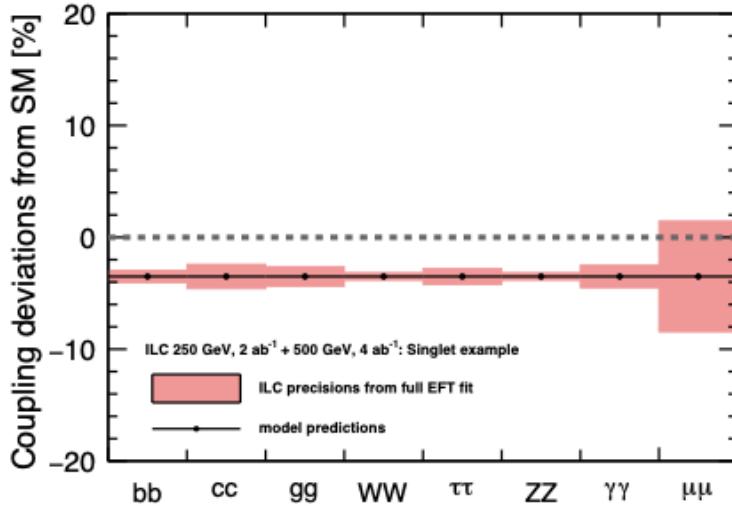
| Parameter | HL-LHC | ILC 500 | FCC-ee | FCC-hh |
|--|--------|---------|--------|--------|
| \sqrt{s} [TeV] | 14 | 0.5 | 0.36 | 100 |
| Yukawa coupling y_t (%) | 3.4 | 2.8 | 3.1 | 1.0 |
| Top mass m_t (%) | 0.10 | 0.031 | 0.025 | — |
| Left-handed top- W coupling $C_{\phi Q}^3$ (TeV^{-2}) | 0.08 | 0.02 | 0.006 | — |
| Right-handed top- W coupling C_{tW} (TeV^{-2}) | 0.3 | 0.003 | 0.007 | — |
| Right-handed top- Z coupling C_{tZ} (TeV^{-2}) | 1 | 0.004 | 0.008 | — |
| Top-Higgs coupling $C_{\phi t}$ (TeV^{-2}) | 3 | 0.1 | 0.6 | — |
| Four-top coupling c_{tt} (TeV^{-2}) | 0.6 | 0.06 | — | 0.024 |



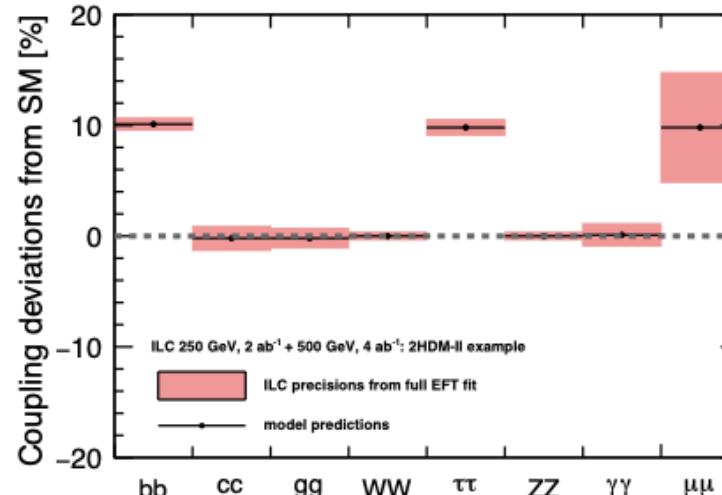
From Snowmass 2021 EF
HF and EW TG's Reports
arXiv:2209.11267,
arXiv:2209.08078

Disentangling models from EFT patterns

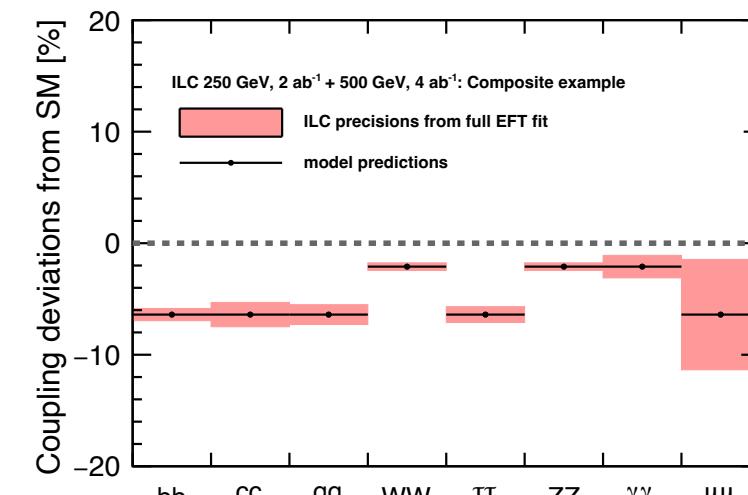
The “inverse Higgs” problem



additional scalar singlet
($m_S=2.8$ TeV, max mixing)



2HDM-II
($M_H=600$ GeV, $\tan\beta=7$)

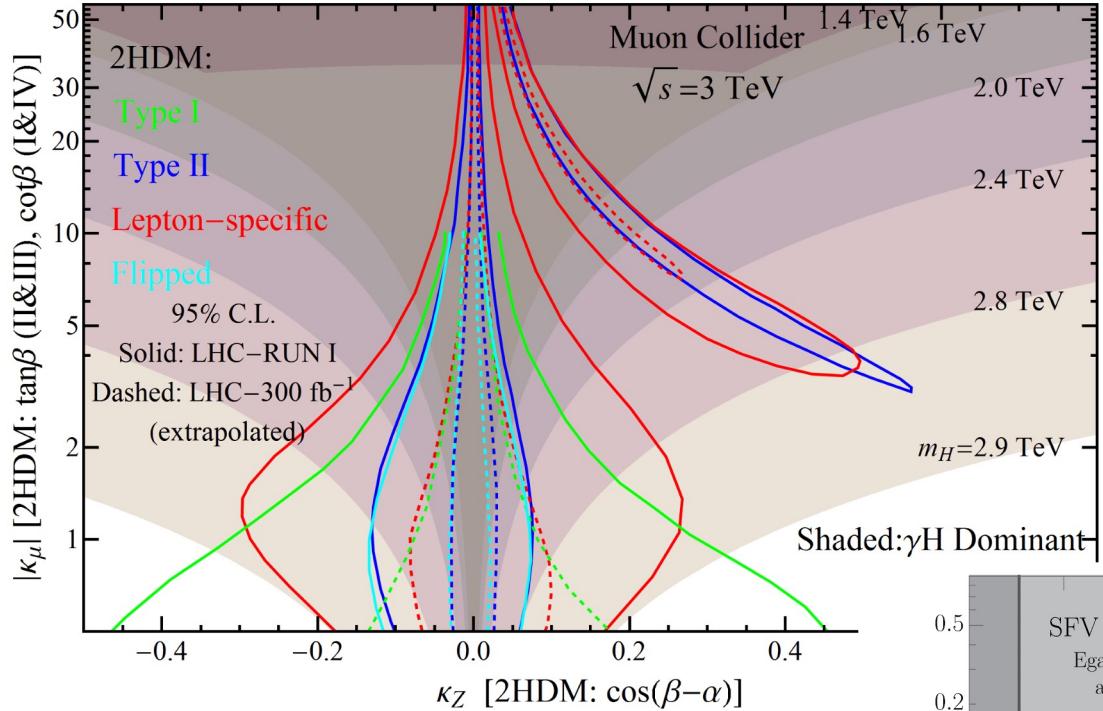


Composite Higgs
($f=1.2$ TeV)

Snowmass 2021: ILC white paper (arXiv: 2203.07622)

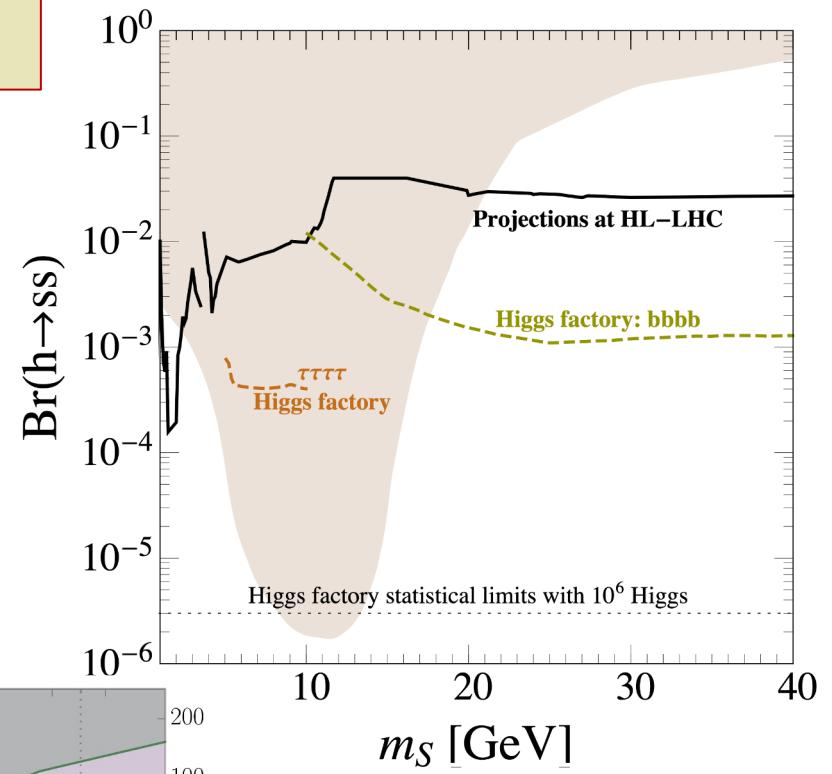
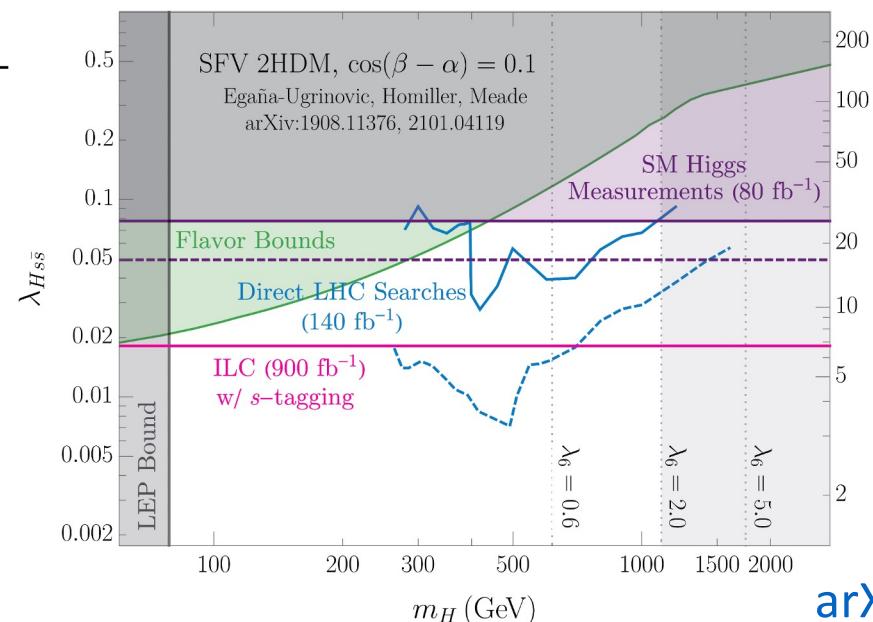
Examples to illustrate the **different patterns of Higgs coupling deviations from different BSM models**

BSM explorations: extended Higgs sectors



[arXiv:2203.07261](https://arxiv.org/abs/2203.07261)

Extended Higgs sectors:
2HDM, extra singlets, ...



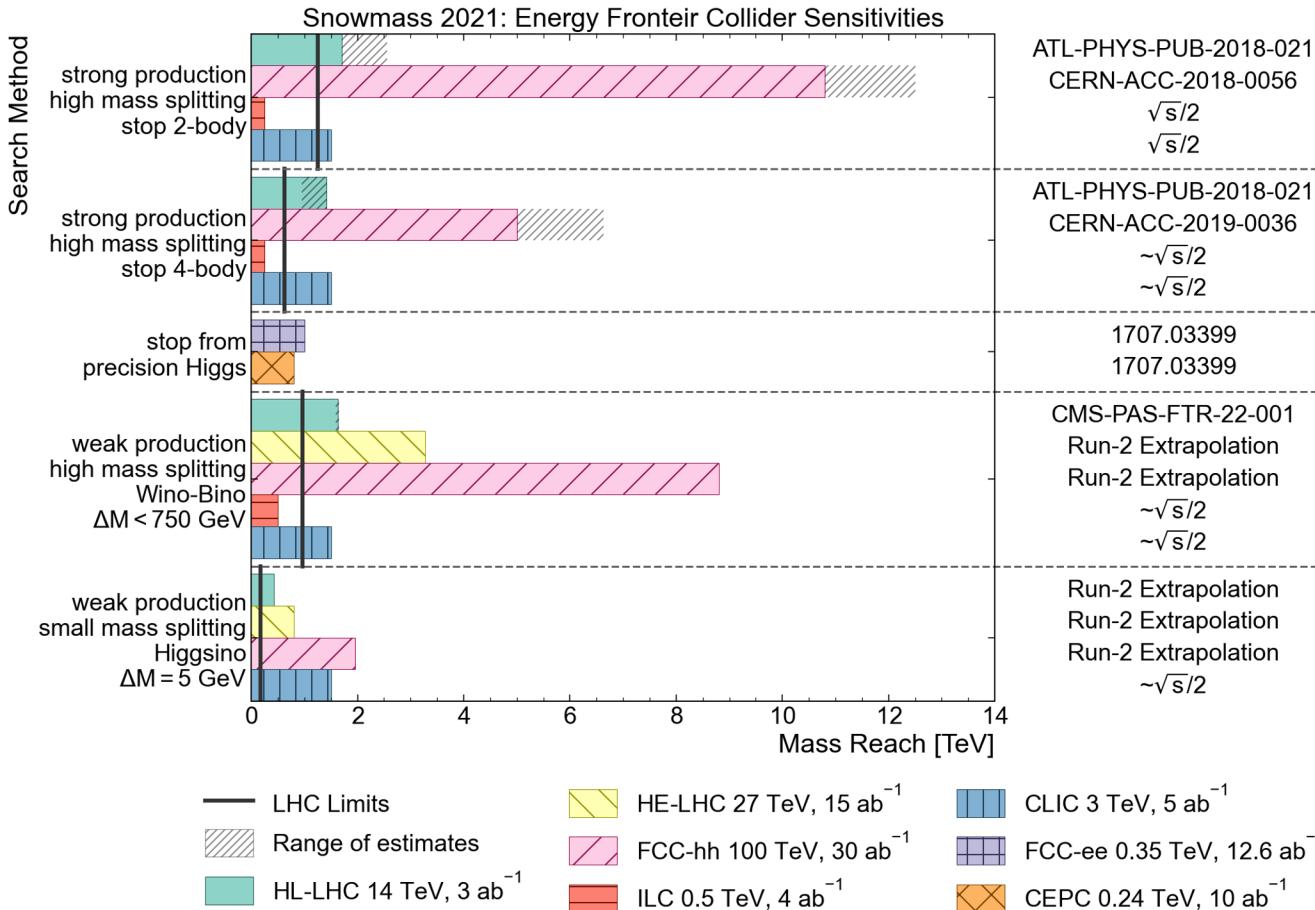
[arXiv:2203.08206](https://arxiv.org/abs/2203.08206)

Higgs and flavor:
probing anomalous
Hss coupling

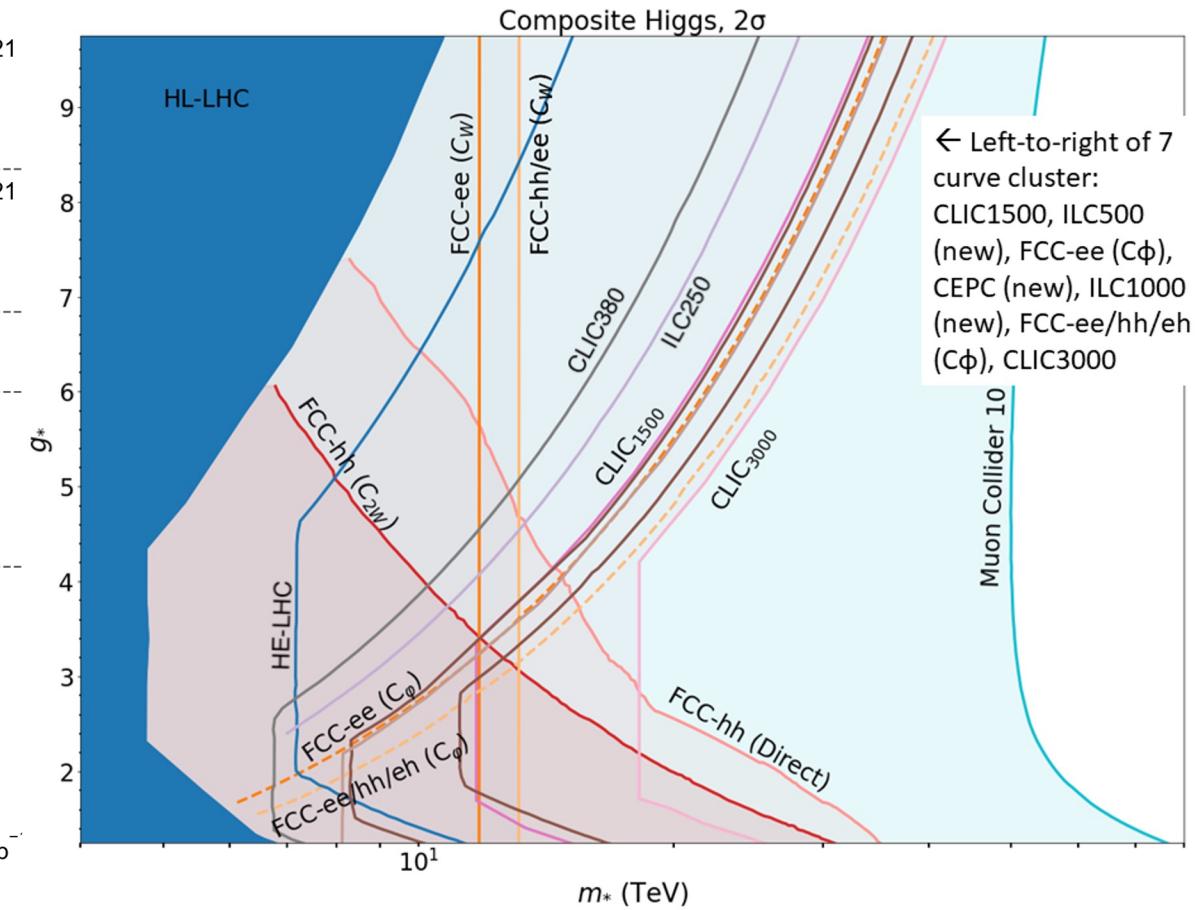
[arXiv:2203.07535](https://arxiv.org/abs/2203.07535)

Examples of BSM model specific explorations

SUSY models



Composite Higgs models



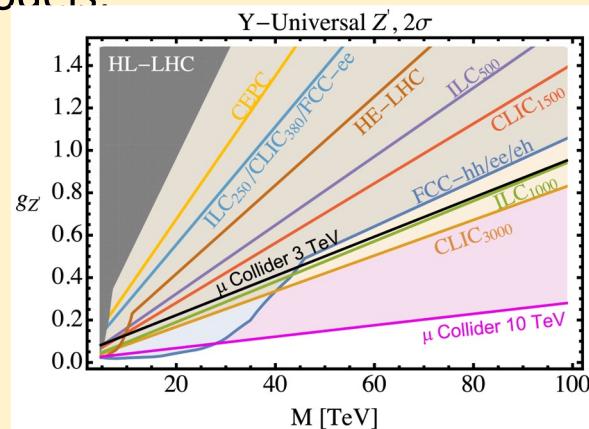
Examples of BSM general explorations

Identify important benchmarks, explore new collider options, focus on the physics messages

Heavy Bosons

Identified simplified models:

- Dilepton
- Dijets
- Diboson (VV, Vh, etc)
- Decays including Heavy Neutrinos



Layout the basic reach of future collider programs **comprehensively** in these simplified modes.

Resonance search and EFT searches are both needed.

[arXiv:1910.11775](https://arxiv.org/abs/1910.11775)

[arXiv:2203.07256](https://arxiv.org/abs/2203.07256)

| Machine | Type | \sqrt{s} (TeV) | $\int L dt$ (ab $^{-1}$) | Source | Z' Model | 5σ (TeV) | 95% CL (TeV) |
|-------------------------|---------------|------------------|---------------------------|--------|-------------------------------------|-----------------|--------------|
| HL-LHC | $p p$ | 14 | 3 | R.H. | $Z'_{SSM} \rightarrow \text{dijet}$ | 4.2 | 5.2 |
| | | | | ATLAS | $Z'_{SSM} \rightarrow l^+ l^-$ | 6.4 | 6.5 |
| | | | | CMS | $Z'_{SSM} \rightarrow l^+ l^-$ | -- | 6.8 |
| | | | | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 6 |
| ILC250/ CLIC380/ FCC-ee | $e^+ e^-$ | 0.25 | 2 | ILC | $Z'_{SSM} \rightarrow f^+ f^-$ | 4.9 | 7.7 |
| | | | | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 7 |
| HE-LHC/ FNAL-SF | $p p$ | 27 | 15 | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 11 |
| | | | | ATLAS | $Z'_{SSM} \rightarrow e^+ e^-$ | 12.8 | 12.8 |
| ILC | $e^+ e^-$ | 0.5 | 4 | ILC | $Z'_{SSM} \rightarrow f^+ f^-$ | 8.3 | 13 |
| | | | | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 13 |
| CLIC | $e^+ e^-$ | 1.5 | 2.5 | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 19 |
| Muon Collider | $\mu^+ \mu^-$ | 3 | 1 | IMCC | $Z'_{Univ}(g_z'=0.2)$ | 10 | 20 |
| ILC | $e^+ e^-$ | 1 | 8 | ILC | $Z'_{SSM} \rightarrow f^+ f^-$ | 14 | 22 |
| | | | | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 21 |
| CLIC | $e^+ e^-$ | 3 | 5 | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 24 |
| FCC-hh | $p p$ | 100 | 30 | R.H. | $Z'_{SSM} \rightarrow \text{dijet}$ | 25 | 32 |
| | | | | EPPSU* | $Z'_{Univ}(g_z'=0.2)$ | -- | 35 |
| | | | | EPPSU | $Z'_{SSM} \rightarrow l^+ l^-$ | 43 | 43 |
| Muon Collider | $\mu^+ \mu^-$ | 10 | 10 | IMCC | $Z'_{Univ}(g_z'=0.2)$ | 42 | 70 |
| VLHC | $p p$ | 300 | 100 | R.H. | $Z'_{SSM} \rightarrow \text{dijet}$ | 67 | 87 |
| Coll. In the Sea | $p p$ | 500 | 100 | R.H. | $Z'_{SSM} \rightarrow \text{dijet}$ | 96 | 130 |

Increasing Z' Sensitivity

Dark matter at colliders

Complementing observation in astrophysics experiments

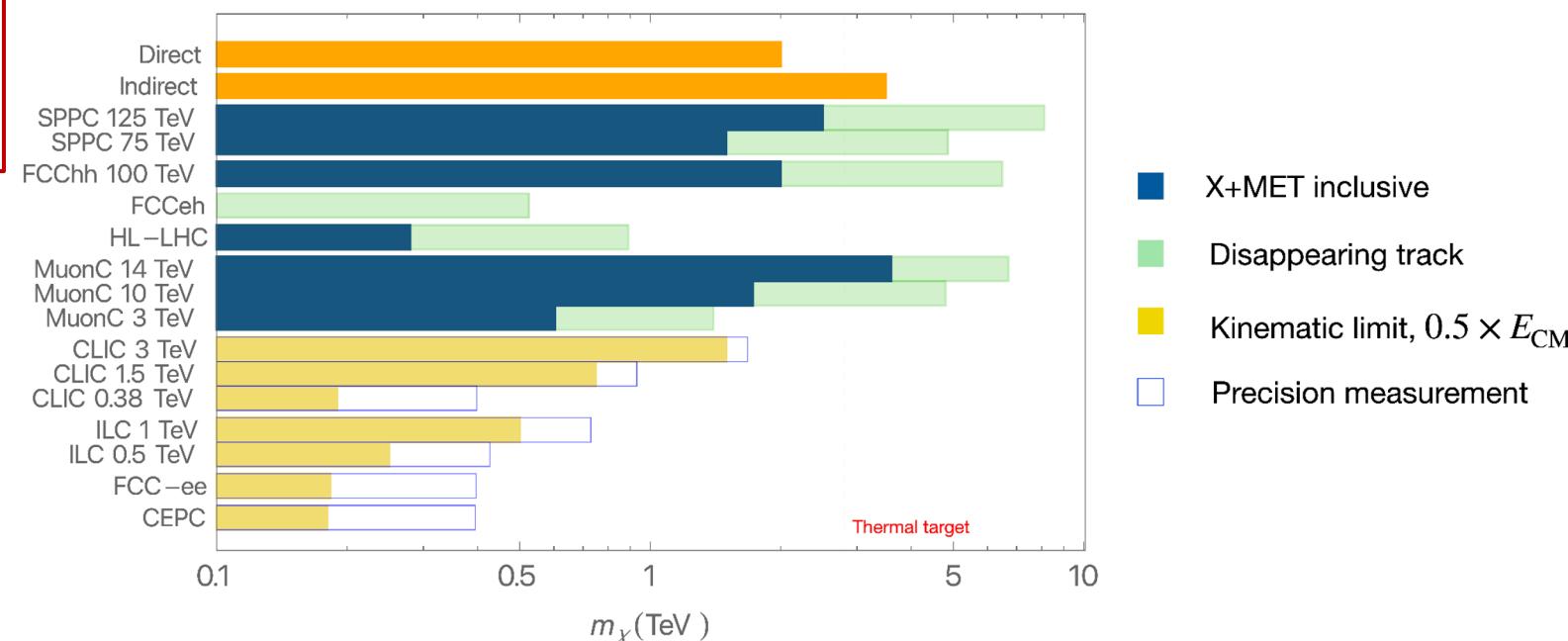
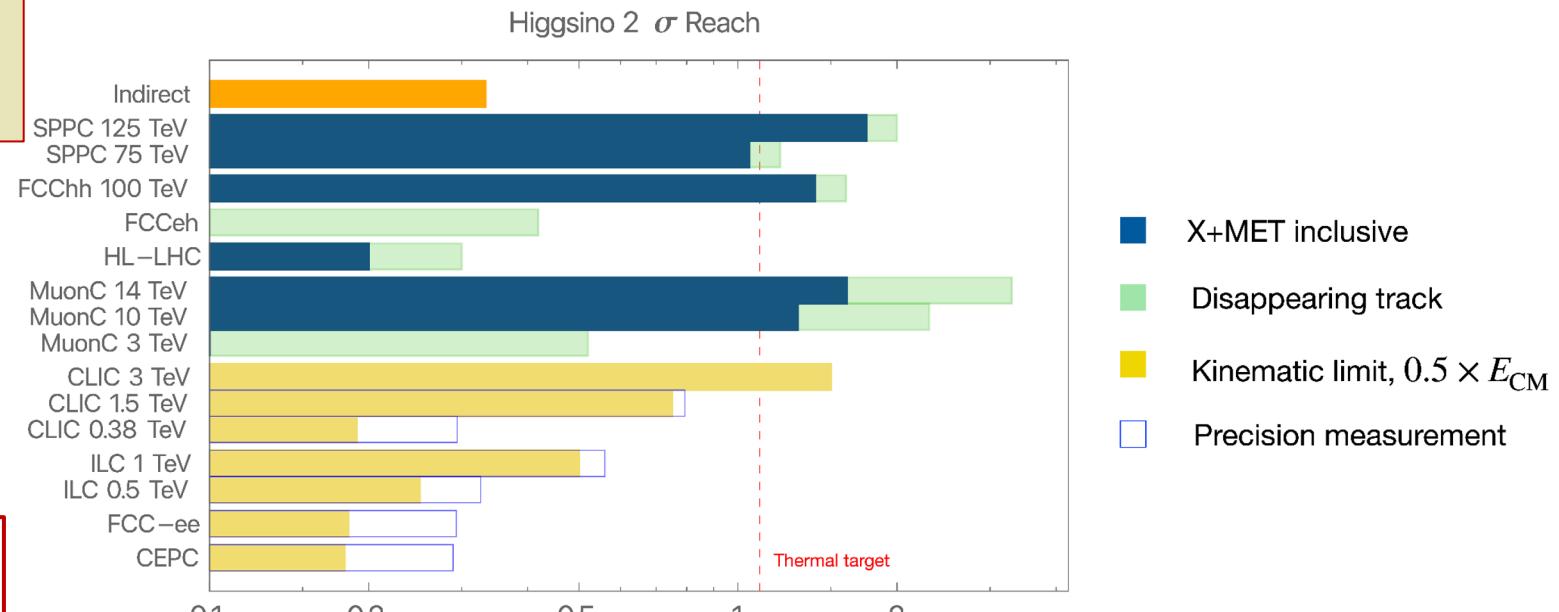
Probing interaction of DM with

SM particles

Discriminating between different models

Example of WIMP DM reach

[arXiv:2210.01770](https://arxiv.org/abs/2210.01770)



Future of Perturbative QCD calculations

| process | known | desired |
|-------------------------------------|---|---|
| $pp \rightarrow H$ | N^3LO_{HTL} , $N^2LO_{QCD}^{(t)}$, $N^{(1,1)}LO_{QCD \otimes EW}^{(HTL)}$ | N^4LO_{HTL} (incl.), $N^2LO_{QCD}^{(b,c)}$ |
| $pp \rightarrow H + j$ | N^2LO_{HTL} , NLO_{QCD} , $N^{(1,1)}LO_{QCD \otimes EW}$ | $N^2LO_{HTL} \otimes NLO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow H + 2j$ | $NLO_{HTL} \otimes LO_{QCD}$ $N^3LO_{QCD}^{(VBF*)}$ (incl.), $N^2LO_{QCD}^{(VBF*)}$, $NLO_{EW}^{(VBF)}$ | $N^2LO_{HTL} \otimes NLO_{QCD} + NLO_{EW}$, $N^2LO_{QCD}^{(VBF)}$ |
| $pp \rightarrow H + 3j$ | NLO_{HTL} , $NLO_{QCD}^{(VBF)}$ | $NLO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow VH$ | $N^2LO_{QCD} + NLO_{EW}$, $NLO_{gg \rightarrow HZ}^{(t,b)}$ | |
| $pp \rightarrow VH + j$ | N^2LO_{QCD} | $N^2LO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow HH$ | $N^3LO_{HTL} \otimes NLO_{QCD}$ | NLO_{EW} |
| $pp \rightarrow H + t\bar{t}$ | $NLO_{QCD} + NLO_{EW}$, N^2LO_{QCD} (off-diag.) | N^2LO_{QCD} |
| $pp \rightarrow H + t/\bar{t}$ | NLO_{QCD} | N^2LO_{QCD} , $NLO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow V$ | N^3LO_{QCD} , $N^{(1,1)}LO_{QCD \otimes EW}$, NLO_{EW} | $N^3LO_{QCD} + N^{(1,1)}LO_{QCD \otimes EW}$, N^2LO_{EW} |
| $pp \rightarrow VV'$ | $N^2LO_{QCD} + NLO_{EW}$, $+ NLO_{QCD} (gg)$ | $NLO_{QCD} (gg, massive loops)$ |
| $pp \rightarrow V + j$ | $N^2LO_{QCD} + NLO_{EW}$ | hadronic decays |
| $pp \rightarrow V + 2j$ | $NLO_{QCD} + NLO_{EW}$, NLO_{EW} | N^2LO_{QCD} |
| $pp \rightarrow V + b\bar{b}$ | NLO_{QCD} | $N^2LO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow VV' + 1j$ | $NLO_{QCD} + NLO_{EW}$ | N^2LO_{QCD} |
| $pp \rightarrow VV' + 2j$ | NLO_{QCD} (QCD), $NLO_{QCD} + NLO_{EW}$ (EW) | Full $NLO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow W^+W^+ + 2j$ | Full $NLO_{QCD} + NLO_{EW}$ | |
| $pp \rightarrow W^+W^- + 2j$ | $NLO_{QCD} + NLO_{EW}$ (EW component) | |
| $pp \rightarrow W^+Z + 2j$ | $NLO_{QCD} + NLO_{EW}$ (EW component) | |
| $pp \rightarrow ZZ + 2j$ | Full $NLO_{QCD} + NLO_{EW}$ | |
| $pp \rightarrow VV'V''$ | NLO_{QCD} , NLO_{EW} (w/o decays) | $NLO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow W^\pm W^+ W^-$ | $NLO_{QCD} + NLO_{EW}$ | |
| $pp \rightarrow \gamma\gamma$ | $N^2LO_{QCD} + NLO_{EW}$ | N^3LO_{QCD} |
| $pp \rightarrow \gamma + j$ | $N^2LO_{QCD} + NLO_{EW}$ | N^3LO_{QCD} |
| $pp \rightarrow \gamma\gamma + j$ | $N^2LO_{QCD} + NLO_{EW}$, $+ NLO_{QCD} (gg \text{ channel})$ | |
| $pp \rightarrow \gamma\gamma\gamma$ | N^2LO_{QCD} | $N^2LO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow 2 \text{ jets}$ | N^2LO_{QCD} , $NLO_{QCD} + NLO_{EW}$ | $N^3LO_{QCD} + NLO_{EW}$ |
| $pp \rightarrow 3 \text{ jets}$ | $N^2LO_{QCD} + NLO_{EW}$ | |
| $pp \rightarrow t\bar{t}$ | N^2LO_{QCD} (w/ decays) + NLO_{EW} (w/o decays) $NLO_{QCD} + NLO_{EW}$ (w/ decays, off-shell effects) N^2LO_{QCD} | N^3LO_{QCD} |
| $pp \rightarrow t\bar{t} + j$ | NLO_{QCD} (w/ decays, off-shell effects) NLO_{EW} (w/o decays) | $N^2LO_{QCD} + NLO_{EW}$ (w/ decays) |
| $pp \rightarrow t\bar{t} + 2j$ | NLO_{QCD} (w/o decays) | $NLO_{QCD} + NLO_{EW}$ (w/ decays) |
| $pp \rightarrow t\bar{t} + Z$ | $NLO_{QCD} + NLO_{EW}$ (w/o decays) NLO_{QCD} (w/ decays, off-shell effects) | $N^2LO_{QCD} + NLO_{EW}$ (w/ decays) |
| $pp \rightarrow t\bar{t} + W$ | $NLO_{QCD} + NLO_{EW}$ (w/ decays, off-shell effects) | $N^2LO_{QCD} + NLO_{EW}$ (w/ decays) |
| $pp \rightarrow t/\bar{t}$ | $N^2LO_{QCD}^{*}$ (w/ decays) NLO_{EW} (w/o decays) | $N^2LO_{QCD} + NLO_{EW}$ (w/ decays) |
| $pp \rightarrow tZj$ | $NLO_{QCD} + NLO_{EW}$ (w/ decays) | $N^2LO_{QCD} + NLO_{EW}$ (w/o decays) |

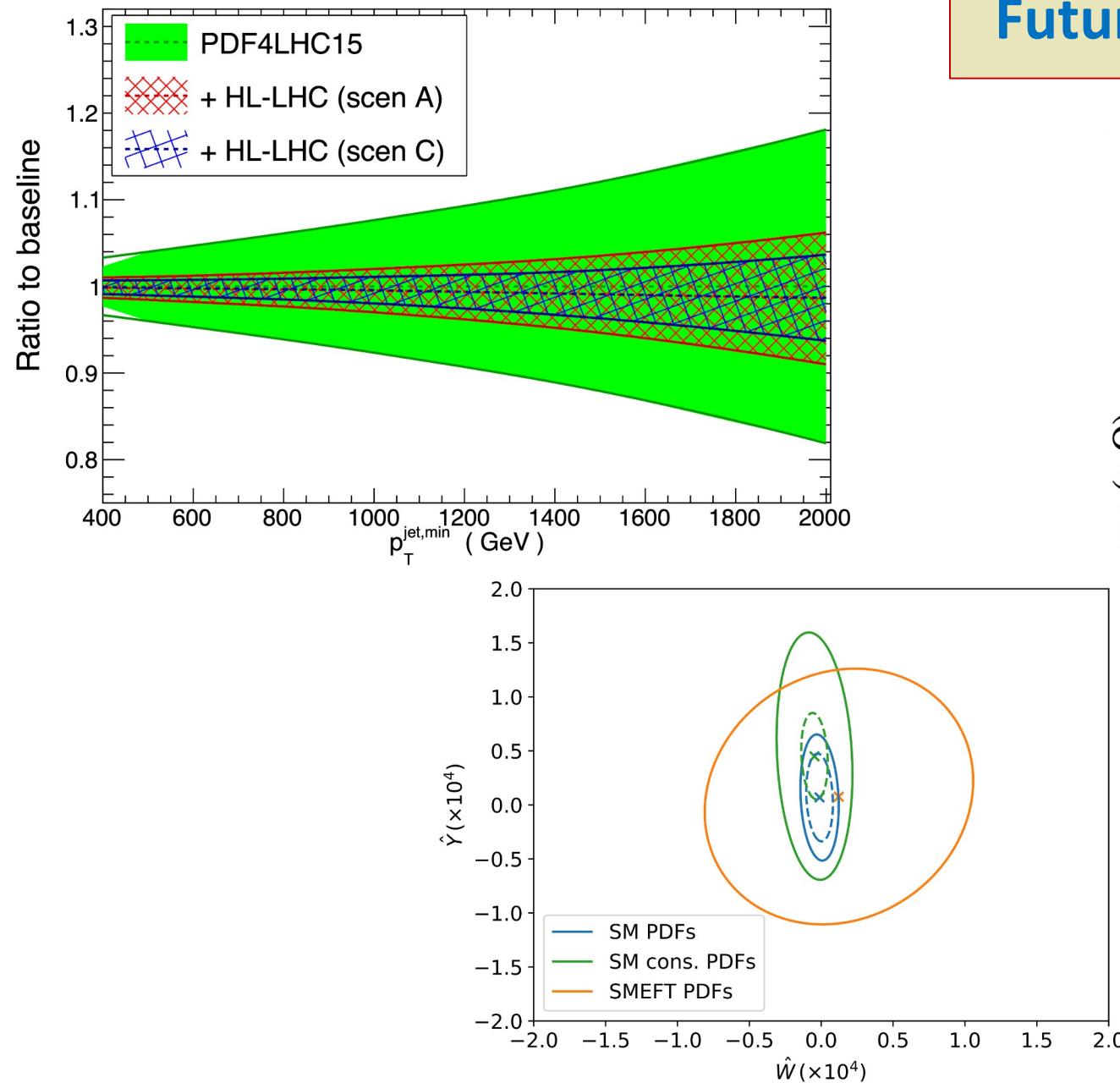
Les Houches wish-list

- α_s uncertainty is a limiting factor in many measurements, e.g. Higgs couplings, at the HL-LHC

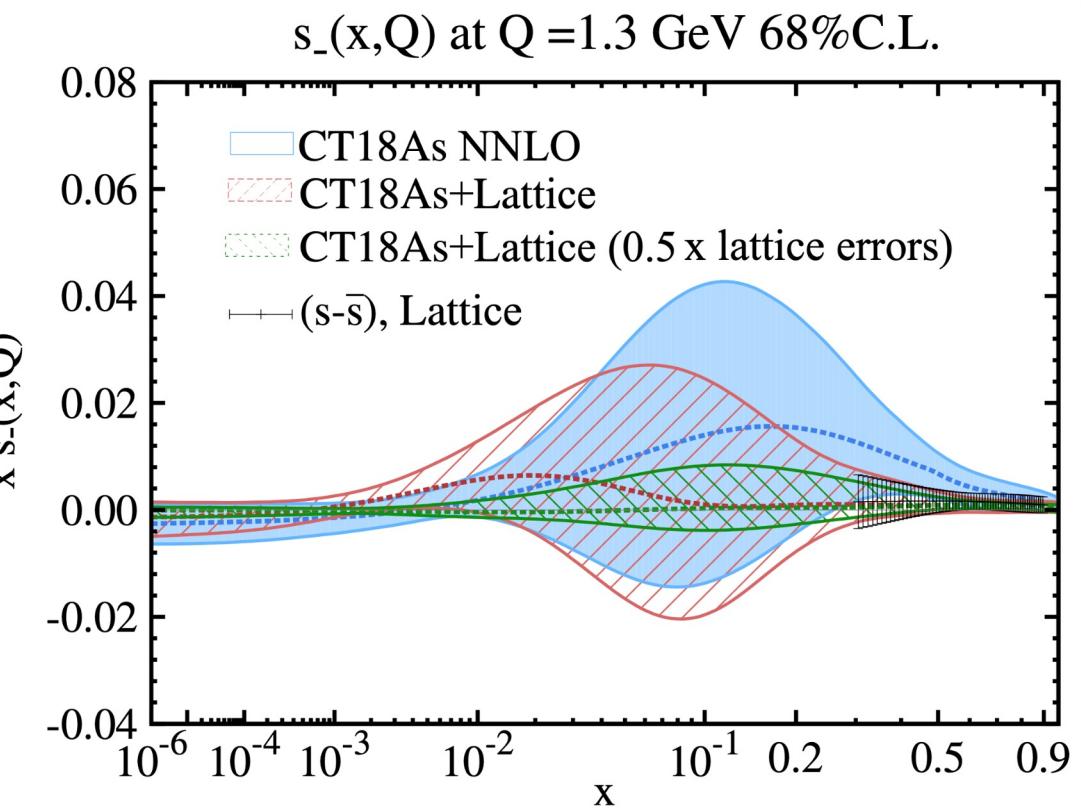
| Method | Relative $\alpha_s(m_Z)$ uncertainty | |
|--------------------------------|--------------------------------------|-------------------------|
| | Current | Near (long-term) future |
| (1) Lattice | 0.7% | $\approx 0.3\% (0.1\%)$ |
| (2) τ decays | 1.6% | $< 1\%$ |
| (3) $Q\bar{Q}$ bound states | 3.3% | $\approx 1.5\%$ |
| (4) DIS & PDF fits | 1.7% | $\approx 1\% (0.2\%)$ |
| (5) e^+e^- jets & evt shapes | 2.6% | $\approx 1.5\% (< 1\%)$ |
| (6) Electroweak fits | 2.3% | $(\approx 0.1\%)$ |
| World average | 0.8% | $\approx 0.4\% (0.1\%)$ |

- **FCC-ee:** $3 \times 10^{12} Z \rightarrow qq$ at the Z pole, and \sqrt{s} calibration 10's keV provides unparalleled α_s precision → searches for small deviations from SM predictions that could signal BSM
- **Jet substructure techniques:**
 - Identification of q/g-initiated jets in $l^+l^- \rightarrow H [\rightarrow gg] Z [\rightarrow ll]$
 - Identification of weak-strahlung emission, and $g \rightarrow tt$ in jets
 - Track functions in jet substructure

Higgs production in gluon fusion @ LHC $\sqrt{s}=14$ TeV



Future of PDF determination



[arXiv:2204.07944](https://arxiv.org/abs/2204.07944)

Setting priorities

In conclusion

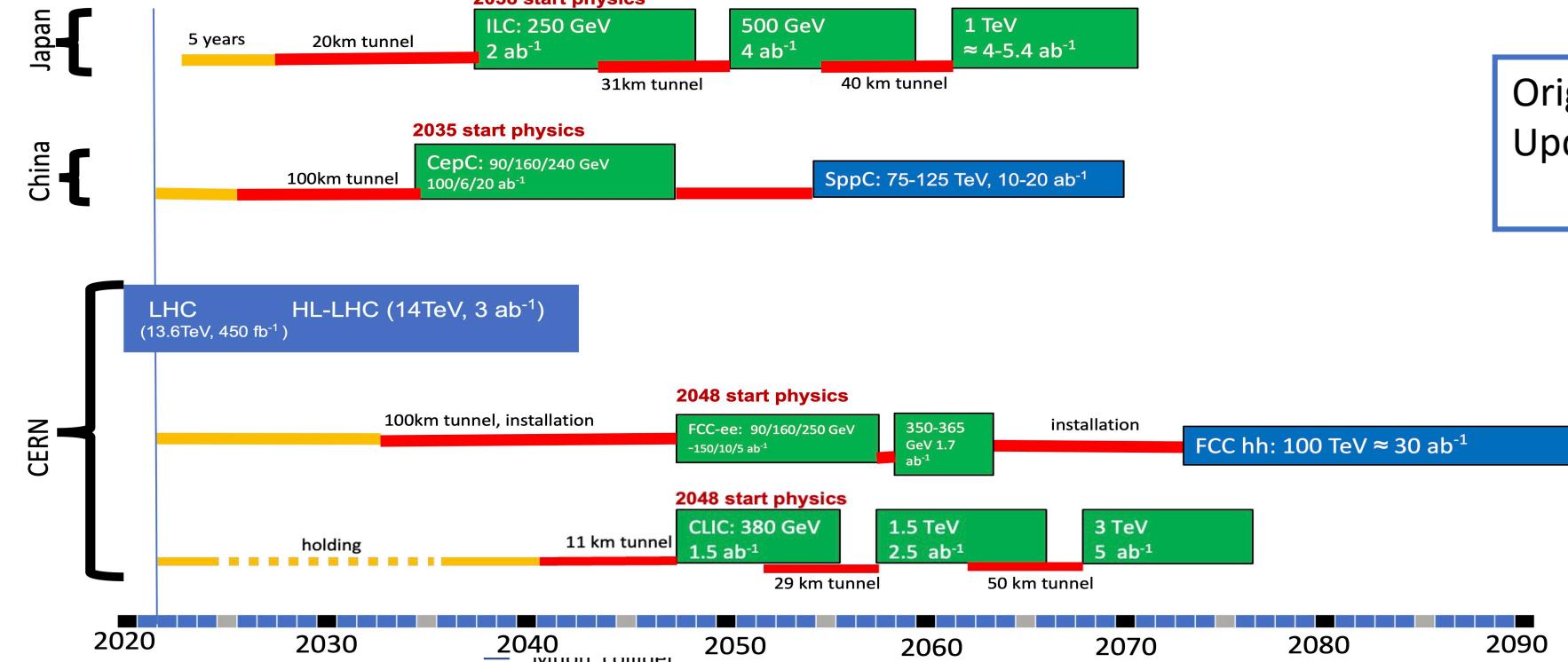
- **The Higgs discovery has been fundamental in opening new avenues to explore physics beyond the SM** and the Higgs-physics program ahead of us promises to start answering some of the remaining fundamental questions in particle physics.
- **Collider physics** remains as a **unique and necessary test of any BSM hypothesis**.
- **Many new directions have been explored during the Snowmass 2021 exercise**, building on previous studies (ESG), and have indicated the need to explore the TeV scale beyond LHC reach by pushing both **precision (Higgs factories) and energy (multi-TeV colliders)**.
- **Increasing the accuracy on SM observables** (Higgs, top, EW) could allow to **test higher scales**: a factor of 10 in precision could allow to test scale in the 10 TeV range and beyond.
- The **possibility of reaching c.o.m. energies above 500 GeV in e+e- collisions is crucial** to improve the full spectrum of HL-LHC measurements, including **top-Higgs and Higgs self-coupling**, as well as **probing extended Higgs sectors and new physics that can elude the LHC**.

The Energy Frontier vision in a nutshell

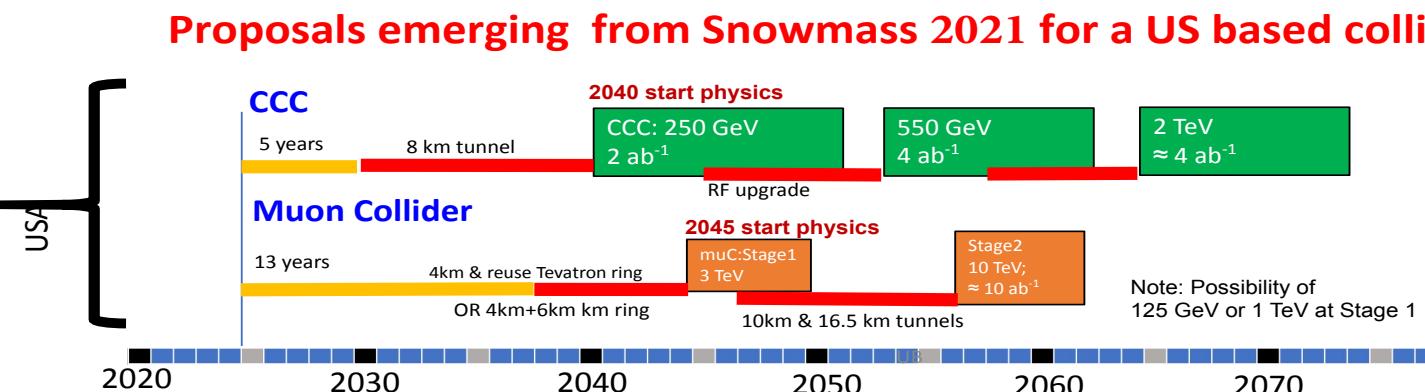
It is essential to

- Complete the HL-LHC program,
- Start now a targeted program for detector R&D for Higgs Factories
- Support a fast start of the construction of a Higgs factory
- Ensure the long-term viability of the field by developing a multi-TeV energy frontier facility such as a *muon collider* or a *hadron collider*.

Timelines



Original timeline from ESG
Updated during Snowmass 2021
(see EF Report)



Renewed interest in lepton colliders:
need supporting R&D in near future